

Wess-Zumino model

$$S = \int d^4x d^2\theta d^2\bar{\theta} \Phi \phi + \left(\int d^4x d^2\theta \left(\frac{1}{2} m \phi^2 + \frac{1}{3} \lambda \phi^3 \right) + \text{c.c.} \right)$$

$$\phi(y, \theta) = \varphi(y) + \theta \psi(y) + \theta^2 F(y)$$

$$[Q_\alpha, Y(x, \theta, \bar{\theta})] = -(-i \partial_\alpha - \sigma^\mu_{\alpha\dot{\alpha}} \bar{\theta}^{\dot{\alpha}} \partial_\mu) Y(x, \theta, \bar{\theta})$$

$$[\bar{Q}_{\dot{\alpha}}, Y(x, \theta, \bar{\theta})] = -(i \partial_{\dot{\alpha}} + \theta^\mu \sigma^\mu_{\dot{\alpha}\alpha} \partial_\mu) Y(x, \theta, \bar{\theta})$$

Action on a chiral superfield.

$$\partial_\alpha = \partial_\alpha + \frac{\partial y^\nu}{\partial \theta^\alpha} \frac{\partial}{\partial y^\nu}$$

$$y^\nu = x^\nu + i \theta \sigma^\nu \bar{\theta}$$

$$= \partial_\alpha + i(\sigma^\nu \bar{\theta})_\alpha \partial_{y^\nu}$$

$$\Rightarrow -i \partial_\alpha - (\sigma^\mu \bar{\theta})_\alpha \partial_\mu = -i \partial_\alpha + (\sigma^\mu \bar{\theta})_\alpha \partial_{y^\mu} - (\sigma^\mu \bar{\theta})_\alpha \partial_{y^\mu} = -i \partial_\alpha$$

$$\Rightarrow [Q_\alpha, \phi(y, \theta)] = +i \partial_\alpha \phi$$

$$\partial_{\dot{\alpha}} = \partial_{\dot{\alpha}} + \frac{\partial y^\nu}{\partial \bar{\theta}^{\dot{\alpha}}} \frac{\partial}{\partial y^\nu} = \partial_{\dot{\alpha}} - i(\theta \sigma^\nu)_{\dot{\alpha}} \partial_\nu$$

$$\Rightarrow i \partial_{\dot{\alpha}} + (\theta \sigma^\nu)_{\dot{\alpha}} \partial_\nu = i \partial_{\dot{\alpha}} + 2(\theta \sigma^\mu)_{\dot{\alpha}} \partial_\mu$$

$$\Rightarrow [\bar{Q}_{\dot{\alpha}}, \phi(y, \theta)] = -2(\theta \sigma^\mu)_{\dot{\alpha}} \partial_\mu \phi(y, \theta)$$

$$\partial_\alpha \phi = \psi_\alpha + 2\theta_\alpha F$$

$$[Q_\alpha, \phi] = [Q_\alpha, \varphi] + [Q_\alpha, \theta \psi] + [Q_\alpha, \theta^2 F] \\ = i(\psi_\alpha + 2\theta_\alpha F)$$

$$\Rightarrow [Q_\alpha, \varphi] = i\psi_\alpha$$

$$[Q_\alpha, \theta \psi] = [Q_\alpha, \theta^\beta \psi_\beta] = -\theta^\beta \{Q_\alpha, \psi_\beta\} = 2i\theta_\alpha F \\ = +\varepsilon_{\alpha\beta} \theta^\beta 2iF$$

$$\Rightarrow \{Q_\alpha, \psi_\beta\} = -2i \varepsilon_{\alpha\beta} F$$

Finally $[Q_\alpha, F] = 0$

The commutators involving $\bar{Q}$ are

$$[\bar{Q}_{\dot{\alpha}}, \phi(\gamma, \theta)] = [\bar{Q}_{\dot{\alpha}}, \phi] + [\bar{Q}_{\dot{\alpha}}, \theta \psi] + [\bar{Q}_{\dot{\alpha}}, \theta^2 F]$$

$$= -2(\theta \sigma^\mu)_{\dot{\alpha}} \partial_\mu \phi$$

$$\Rightarrow \boxed{[\bar{Q}_{\dot{\alpha}}, \phi] = 0}$$

$$[\bar{Q}_{\dot{\alpha}}, \theta \psi] = -2(\theta \sigma^\mu)_{\dot{\alpha}} \partial_\mu \psi$$

$$\Rightarrow \boxed{\{\bar{Q}_{\dot{\alpha}}, \psi_\alpha\} = + 2 \sigma_{\alpha \dot{\alpha}}^\mu \partial_\mu \psi}$$

$$[\bar{Q}_{\dot{\alpha}}, \theta^2 F] = -2(\theta \sigma^\mu)_{\dot{\alpha}} \partial_\mu (\theta \psi) = -2 \theta^\alpha \sigma_{\alpha \dot{\alpha}}^\mu \theta^\beta \partial_\mu \psi_\beta$$

$$= -2 \left(-\frac{1}{2} \varepsilon^{\mu\rho} \theta^2 \right) \sigma_{\alpha \dot{\alpha}}^\mu \partial_\mu \psi_\beta$$

$$= + (\psi \sigma^\mu)_{\dot{\alpha}} \theta^2$$

$$\Rightarrow \boxed{[\bar{Q}_{\dot{\alpha}}, F] = + (\psi \sigma^\mu)_{\dot{\alpha}} \partial_\mu}$$

Solution for F

The terms involving F in the action are

$$S = \int d^4x \left(\bar{F} F + (m \psi + \lambda \psi^2) F + (\bar{m} \bar{\psi} + \bar{\lambda} \bar{\psi}^2) \bar{F} \right)$$

The equations of motion for F are

$$\bar{F} = -(m \psi + \lambda \psi^2)$$

$$F = -(\bar{m} \bar{\psi} + \bar{\lambda} \bar{\psi}^2)$$

On-shell,

$$[Q_\alpha, \psi] = +i \psi_\alpha \quad [\bar{Q}_{\dot{\alpha}}, \psi] = 0$$

$$\{Q_\alpha, \psi_\beta\} = 2i \varepsilon_{\alpha\beta} (m \psi + \lambda \psi^2)$$

$$\{\bar{Q}_{\dot{\alpha}}, \psi_\alpha\} = 2 \sigma_{\alpha \dot{\alpha}}^\mu \partial_\mu \psi$$

Let us now study explicitly the supersymmetry algebra when realized on ψ, ψ fields.

$$\{ \bar{Q}_{\dot{\alpha}}, [Q_{\alpha}, \psi] \} = [\underbrace{\{ Q_{\alpha}, \bar{Q}_{\dot{\alpha}} \}}_{2\sigma_{\alpha\dot{\alpha}}^{\mu} p_{\mu}}, \psi] - [Q_{\alpha}, \underbrace{[\bar{Q}_{\dot{\alpha}}, \psi]}_0]$$

$$= \{ \bar{Q}_{\dot{\alpha}}, i\psi_{\alpha} \} = i 2\sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu} \psi$$

$$\Rightarrow [p_{\mu}, \psi] = i \partial_{\mu} \psi \quad \checkmark$$

$$[\bar{Q}_{\dot{\alpha}}, \{ Q_{\alpha}, \psi_{\beta} \}] = [\underbrace{\{ Q_{\alpha}, \bar{Q}_{\dot{\alpha}} \}}_{2\sigma_{\alpha\dot{\alpha}}^{\mu} p_{\mu}}, \psi_{\beta}] - [Q_{\alpha}, \{ \bar{Q}_{\dot{\alpha}}, \psi_{\beta} \}]$$

$$LHS = [\bar{Q}_{\dot{\alpha}}, 2i \epsilon_{\alpha\beta} (\overline{m\psi + \lambda\psi^2})]$$

$$= 2i \epsilon_{\alpha\beta} [\underbrace{Q_{\alpha}^{\dagger}, (m\psi + \lambda\psi^2)^{\dagger}}]$$

$$- [Q_{\alpha}, m\psi + \lambda\psi^2]^{\dagger} = - (m(i\psi_{\alpha}) + 2\lambda\psi(i\psi_{\alpha}))^{\dagger}$$

$$= i(\bar{m} \bar{\psi}_{\dot{\alpha}} + 2\bar{\lambda} \bar{\psi} \bar{\psi}_{\dot{\alpha}})$$

$$RHS = -2i \sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu} \psi_{\beta} - [Q_{\alpha}, 2\sigma_{\beta\dot{\alpha}}^{\mu} \partial_{\mu} \psi]$$

$$= 2i \sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu} \psi_{\beta} - 2\sigma_{\beta\dot{\alpha}}^{\mu} \partial_{\mu} (i\psi_{\alpha})$$

Contract LHS and RHS by $\epsilon^{\alpha\beta}$

$$\Rightarrow 4 (\bar{m} \bar{\psi}_{\dot{\alpha}} + 2\bar{\lambda} \bar{\psi} \bar{\psi}_{\dot{\alpha}}) = 4i \sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu} \psi^{\alpha}$$

So we need to check $i\mathbb{I}$

$$\boxed{\sigma_{\alpha\dot{\alpha}}^{\mu} \partial_{\mu} \psi^{\alpha} = -i(\bar{m} \bar{\psi}_{\dot{\alpha}} + 2\bar{\lambda} \bar{\psi} \bar{\psi}_{\dot{\alpha}})}$$

The terms in the action involving $\psi, \bar{\psi}$

$$S = \int d^4x d^2\theta d^2\bar{\theta} (\bar{\theta} \bar{\psi}(\bar{\gamma})) (\theta \psi)(\gamma) + \left(\int d^4x d^2\theta \left(\frac{1}{2} m (\theta \psi)^2 + \lambda \psi (\theta \psi)^2 \right) + c.c \right)$$

$$\bar{\theta} \bar{\psi}(\bar{\gamma}) \theta \psi(\gamma) = (\bar{\theta} \bar{\psi}(x) - i(\theta \sigma^{\mu} \bar{\theta}) \bar{\theta} \partial_{\mu} \bar{\psi}(x)) \times$$

$$(\theta \psi(x) + i(\theta \sigma^{\mu} \bar{\theta}) \theta \partial_{\mu} \psi(x))$$

$$= \bar{\theta} \bar{\psi}(x) \theta \psi(x) + i(\bar{\theta} \bar{\psi})(\theta \sigma^{\mu} \bar{\theta}) \theta \partial_{\mu} \psi$$

$$- i(\theta \sigma^{\mu} \bar{\theta})(\bar{\theta} \partial_{\mu} \bar{\psi})(\theta \psi)$$

The only terms which survive the full superspace integral are

$$\begin{aligned}
 & i (\bar{\theta} \bar{\psi}) (\theta \sigma^\mu \bar{\theta}) (\theta \partial_\mu \psi) - i (\theta \sigma^\mu \bar{\theta}) (\bar{\theta} \partial_\mu \psi) (\theta \psi) \\
 &= 2 i (\bar{\theta} \bar{\psi}) (\theta \sigma^\mu \bar{\theta}) (\theta \partial_\mu \psi) - \underbrace{\partial_\mu (i (\bar{\theta} \bar{\psi}) (\theta \psi) (\theta \sigma^\mu \bar{\theta}))}_{\text{drops the total derivative}} \\
 &= 2 i \bar{\theta}_{\dot{\alpha}} \bar{\psi}^{\dot{\alpha}} \theta^\beta \sigma_{\beta\dot{\beta}}^\mu \bar{\theta}^{\dot{\beta}} \theta^\alpha \partial_\mu \psi_\alpha \\
 &= + 2 i \underbrace{(\bar{\theta}_{\dot{\alpha}} \bar{\theta}^{\dot{\alpha}})}_{\frac{1}{2} \delta_{\dot{\alpha}}^{\dot{\beta}} \bar{\theta}^2} (\underbrace{\theta^\beta \theta^\alpha}_{-\frac{1}{2} \varepsilon^{\beta\alpha} \theta^2}) \bar{\psi}^{\dot{\alpha}} \sigma_{\beta\dot{\beta}}^\mu \partial_\mu \psi_\alpha \\
 &= -\frac{i}{2} \theta^2 \bar{\theta}^2 \bar{\psi}^{\dot{\beta}} \sigma_{\beta\dot{\beta}}^\mu \partial_\mu \psi^\beta \\
 &= \frac{i}{2} \theta^2 \bar{\theta}^2 (\partial_\mu \psi) \sigma^\mu \bar{\psi} = -\frac{i}{2} \theta^2 \bar{\theta}^2 \bar{\psi} \sigma^\mu (\partial_\mu \psi)
 \end{aligned}$$

The superpotential contributes

$$\begin{aligned}
 & \frac{1}{2} m (\theta \psi)^2 + \lambda \varphi (\theta \psi)^2 \\
 (\theta \psi)^2 &= \theta^\alpha \psi_\alpha \theta^\beta \psi_\beta = -\theta^\alpha \theta^\beta \psi_\alpha \psi_\beta = \frac{1}{2} \varepsilon^{\alpha\beta} \theta^2 \psi_\alpha \psi_\beta \\
 &= -\frac{1}{2} \theta^2 \psi^\beta \psi_\beta = -\frac{1}{2} \theta^2 \psi^2
 \end{aligned}$$

Finally, the terms that contribute to the equation of motion for ψ are

$$-\frac{i}{2} \bar{\psi} \sigma^\mu \partial_\mu \psi + \left[\left(\frac{1}{2} m + \lambda \varphi \right) \left(-\frac{1}{2} \psi^2 \right) + \text{c.c} \right]$$

Variation wrt $\bar{\psi}_{\dot{\alpha}}$

$$\Rightarrow -\frac{i}{2} (\bar{\sigma}^\mu \partial_\mu \psi)^{\dot{\alpha}} + \overline{\left(\frac{1}{2} m + \lambda \varphi \right) \left(-\frac{1}{2} \right) (\psi^2)} = 0$$

$$\Rightarrow \boxed{(\bar{\sigma}^\mu \partial_\mu \psi)^{\dot{\alpha}} = i \left(\bar{m} + 2 \bar{\lambda} \bar{\varphi} \right) \bar{\psi}^{\dot{\alpha}}}$$

This is the same as the condition found before if we use $\bar{\sigma}_\mu^{\dot{\alpha}\alpha} = \varepsilon^{\dot{\alpha}\beta} \varepsilon^{\alpha\gamma} \sigma_{\mu,\beta\gamma}$