## An elementary proof of the contraction theorem

Quang-Cuong Pham LPPA, Collège de France Paris, France cuong.pham@normalesup.org

April 10, 2007

## Abstract

Nonlinear contraction theory was introduced by Lohmiller and Slotine in [1]. The development of this theory relies heavily on differential notions, especially on the so-called "infinitesimal displacement at fixed time"  $\delta \mathbf{x}$ . However, these notions were never explicitly defined, and as a consequence, some of the proofs remained rather obscure. In these notes, we aim at providing rigorous proofs of the basic theorems of nonlinear contraction theory, using only elementary differential calculus.

Consider a smooth nonlinear dynamical equation

$$\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}, t) \tag{1}$$

Consider two trajectories starting respectively at **a** and **b**. At time t = 0, consider a smooth path  $\gamma(0)$  that connects **a** and **b**. The intermediate values along that path are parameterized by  $u \in [0, 1]$ , with  $\mathbf{x}(0, 0) = \mathbf{a}$  and  $\mathbf{x}(0, 1) = \mathbf{b}$ . For all t > 0 define now the path  $\gamma(t) = {\mathbf{x}(t, u) | u \in [0, 1]}$ , where  $\mathbf{x}(t, u)$  is the position at time t of a particle that starts at  $\mathbf{x}(0, u)$  at t = 0. Standard results in differential calculus guarantee that for all  $t, \gamma(t)$  is a smooth path whose **M**-length is given by

$$L_{\mathbf{M}}(\gamma(t)) = \int_0^1 \sqrt{\left(\frac{\partial \mathbf{x}}{\partial u}(t, u)\right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left(\frac{\partial \mathbf{x}}{\partial u}(t, u)\right)} du$$

Define now

$$\phi(t) = \int_0^1 \left(\frac{\partial \mathbf{x}}{\partial u}(t, u)\right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left(\frac{\partial \mathbf{x}}{\partial u}(t, u)\right) du$$

By Schwarz's inequality, one has  $L^2_{\mathbf{M}}(\gamma(t)) \leq \phi(t)$ . The time derivative of  $\phi$  is next given by

$$\begin{split} \frac{d\phi}{dt} &= \frac{d}{dt} \int_0^1 \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) du \\ &= \int_0^1 \frac{\partial}{\partial t} \left( \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) \right) du \\ &= \int_0^1 \left( \frac{\partial}{\partial t} \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) \right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) \\ &+ \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right)^T \left( \frac{\partial}{\partial t} \mathbf{M}(\mathbf{x}(t, u), t) \right) \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) \\ &+ \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right)^T \mathbf{M}(\mathbf{x}(t, u), t) \left( \frac{\partial}{\partial t} \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) \right) du \end{split}$$

Using successively Schwarz's theorem, equation (1) and the chain rule yields

$$\frac{\partial}{\partial t} \left( \frac{\partial \mathbf{x}}{\partial u}(t, u) \right) = \frac{\partial}{\partial u} \left( \frac{\partial \mathbf{x}}{\partial t}(t, u) \right) = \frac{\partial}{\partial u} \left( \mathbf{f}(\mathbf{x}(t, u), t)) = \mathbf{J}(\mathbf{x}(t, u), t) \cdot \frac{\partial \mathbf{x}}{\partial u}(t, u) \right)$$

where  $\mathbf{J}(\mathbf{x}(t, u), t)$  denotes the Jacobian matrix of  $\mathbf{f}$  computed at  $(\mathbf{x}(t, u), t)$ . Thus

$$\begin{array}{rcl} \frac{d\phi}{dt} &=& \int_0^1 \left(\frac{\partial \mathbf{x}}{\partial u}(t,u)\right)^T \left(\mathbf{J}^T \mathbf{M} + \dot{\mathbf{M}} + \mathbf{M} \mathbf{J}\right) \left(\frac{\partial \mathbf{x}}{\partial u}(t,u)\right) du \\ &\leq& -2\lambda \int_0^1 \left(\frac{\partial \mathbf{x}}{\partial u}(t,u)\right)^T \mathbf{M} \left(\frac{\partial \mathbf{x}}{\partial u}(t,u)\right) du \\ &=& -2\lambda \phi(t) \end{array}$$

where  $\lambda$  verifies

$$\forall \mathbf{x}, t \quad \mathbf{J}^T \mathbf{M} + \dot{\mathbf{M}} + \mathbf{M} \mathbf{J} \le -2\lambda \mathbf{M}$$

Applying Gronwall's lemma then yields

$$\forall t \ge 0 \quad \phi(t) \le \phi(0) e^{-2\lambda t}$$

Thus  $\phi(t)$  tends exponentially to 0 with rate  $2\lambda$ , which implies that  $L_{\mathbf{M}}(\gamma(t))$  tends exponentially to 0 with rate  $\lambda$ .

## References

 W. Lohmiller, J.-J. Slotine. On Contraction Analysis for Nonlinear Systems. Automatica, 34 (6):671–682, 1998.