

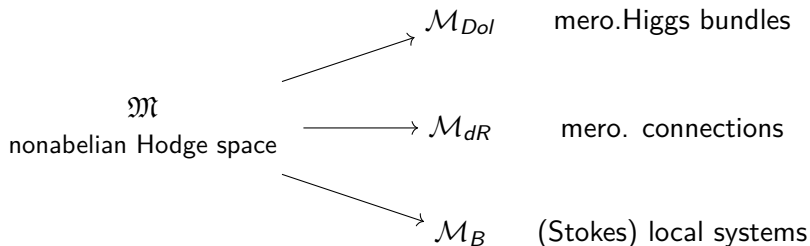
Diagrams and the classification of wild character varieties

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(irregular) nonabelian Hodge spaces



- $\mathcal{M}_{Dol} \simeq \mathcal{M}_{dR}$: (wild) nonabelian Hodge correspondence
- $\mathcal{M}_{dR} \simeq \mathcal{M}_B$: (irregular) Riemann-Hilbert correspondence
- $\mathcal{M}_B$ is a (wild) character variety.

Hyperkähler manifold, 3 different algebraic structures.

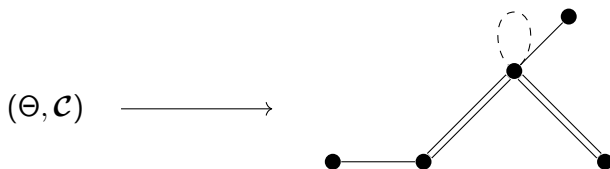
Outline

The moduli space depends on the choice of a "wild Riemann surface", i.e. a Riemann surface Σ together with singularity data $(\Theta, \mathcal{C})$.

Several different singularity data can give rise to isomorphic moduli spaces: different "representations" of the same space

Question: can we classify wild character varieties? Is there a systematic way to find all wild Riemann surfaces giving rise to the same moduli space?

Idea: associate a diagram to singularity data



such that, at least in some cases, wild Riemann surfaces corresponding to isomorphic moduli spaces have the same diagram.

- 1 Wild character varieties
- 2 Construction of diagrams
- 3 Applications to Painlevé equations
- 4 Further questions

Meromorphic connections

- Let Σ be a smooth complex algebraic curve. Here $\Sigma = \mathbb{P}^1 = \mathbb{A}^1 \cup \{\infty\}$.
- Consider (E, ∇) vector bundle with algebraic connection on $\Sigma \setminus \{a_1, \dots, a_m\}$.
- In a trivialization and with a choice of coordinate:

$$\nabla = d - A(z)dz,$$

with A having poles at singular points.

→ corresponds to a system of linear differential equations.

Formal normal form

Turritin-Levelt theorem

After passing to a finite cover $t = z^{1/r}$, any connection on the formal punctured disk is formally isomorphic (via formal gauge transformations) to a connection of the form

$$d - (dQ + \Lambda \frac{dz}{z})$$

where $Q = \text{diag}(q_1, \dots, q_k)$ diagonal matrix with coefficients in $z^{-1/r}\mathbb{C}[z^{-1/r}]$, Λ constant block diagonal.

The q_i are the *exponential factors* of the connection: indeed formal solutions $e^Q z^\Lambda$ are linear combinations of terms with $e^{q_i(z)}$.

Λ : formal monodromy

Untwisted case: when no ramification, i.e. $r = 1$.

The exponential local system

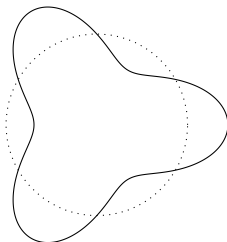
- Let ∂ be the circle of directions around $z = 0$.
- Exponential local system $\mathcal{I} \rightarrow \partial$: sections on sectors are germs of functions of the form

$$q(z) = \sum_{i=1}^s b_i z^{-i/r},$$

- Connected components are *circles* $\langle q \rangle$. $\langle q \rangle \rightarrow \partial$ cover of ∂ of order r .
 $\mathcal{I}$ is a collection of circles
- $r = \text{ram}(q)$ ramification index,
 $s = \text{lrr}(q)$ irregularity.
 s/r is the slope of q .
- (local) singularity data: choice of active circles $\langle q_i \rangle$, with multiplicities n_i , and conjugacy classes $\mathcal{C}_i \subset GL_{n_i}(\mathbb{C})$
 - ▶ Irregular class $\Theta = n_1 \langle q_1 \rangle + \dots n_k \langle q_k \rangle$.
 - ▶ $\mathcal{C} = (\mathcal{C}_1, \dots, \mathcal{C}_k)$ conjugacy classes of blocks of Λ .

Stokes phenomenon

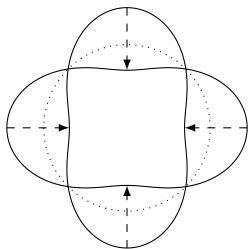
Main idea: growth rates of e^{q_i} change when going around the singularity



Stokes diagram: draw $e^{\operatorname{Re}(q(\epsilon i \theta))}$, $\epsilon \ll 1$, here for $q = z^3$.

Stokes matrices

Change of dominance between the exponential factors e^{q_i} depending on the direction.



Singular directions: $d \in \mathbb{A}$ where $q_i - q_j \in \mathbb{R}_{<0}$ for some (i, j) .

Stokes arrow $\langle q_i \rangle \rightarrow \langle q_j \rangle$

Associated block in Stokes matrix $S_d = \begin{pmatrix} 1 & 0 & * \\ 0 & \ddots & 0 \\ 0 & 0 & 1 \end{pmatrix}$ in position (i, j) .

Explicit presentation of wild character varieties

$$\mathcal{M}_B \cong \{(C_1^{-1} h_1 S_{k_1}^{(1)} \dots S_1^{(1)} C_1) \dots (C_m^{-1} h_m S_{k_m}^{(m)} \dots S_1^{(m)} C_m) = 1\} / G \times \mathbf{H},$$

This corresponds to a quasi-Hamiltonian description

- Fission spaces: $\mathcal{A}(V_{a_i}^0) \ni (C_i, h_i, S_1^{(i)}, \dots, S_{k_i}^{(i)})$
- For several singularities, fusion of fission spaces (here $G = GL_n(\mathbb{C})$)

$$\mathrm{Hom}_{\mathbb{S}}(\Theta) \simeq \mathcal{A}(V_{a_1}^0) \otimes \dots \otimes \mathcal{A}(V_{a_m}^0) // G.$$

- $\mathcal{M}_B$ is now obtained by taking the multiplicative symplectic reduction at the conjugacy classes of formal monodromies.

$$\mathcal{M}_B(\Theta, \mathcal{C}) = \mathrm{Hom}_{\mathbb{S}}(\Theta) //_{\mathcal{C}} \mathbf{H}.$$

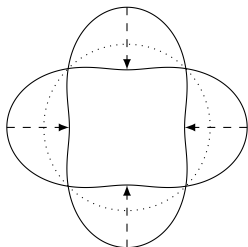
where $\mathbf{H} = H_1 \times \dots \times H_m$, H_i is a subgroup of G of block diagonal matrices

Untwisted example

- Singularity at infinity, two exponential factors $q_1 = z^2$, $q_2 = -z^2$, 4 singular directions, 4 Stokes matrices.
- Moduli space

$$\mathcal{M}_B = \{hS_4S_3S_2S_1 = 1\} // H$$

$$\text{with } S_{2i+1} = \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix}, S_{2i} = \begin{pmatrix} 1 & 0 \\ * & 1 \end{pmatrix}, h = \begin{pmatrix} * & 0 \\ 0 & * \end{pmatrix}$$

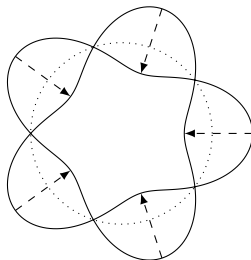


Twisted example

- Singularity at infinity, one exponential factor $z^{5/2}$, 5 singular directions, 5 Stokes matrices.
- Moduli space (of dimension 2)

$$\mathcal{M}_B = \{hS_5S_4S_3S_2S_1 = 1\} // H$$

$$\text{with } S_{2i+1} = \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix}, S_{2i} = \begin{pmatrix} 1 & 0 \\ * & 1 \end{pmatrix}, h = \begin{pmatrix} 0 & * \\ * & 0 \end{pmatrix}$$

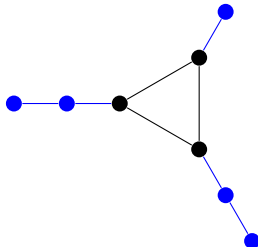


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Diagram associated to a connection

General structure:

- Core diagram encoding the irregular class Θ
- "Legs" (linear quivers) encoding the conjugacy classes of formal monodromies



This relies on the fact that a conjugacy class $\mathcal{C} \subset GL(N)$ can be characterized by a linear quiver

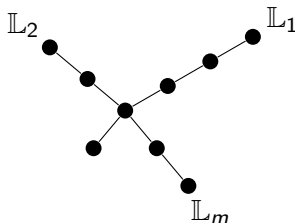
Regular connections

- Singularity data: conjugacy classes $\mathcal{C}$ of formal monodromies.
- Quivers from conjugacy classes [\[Kraft-Procesi 82, Crawley-Boevey 01-06\]](#). Let $\mathcal{C} \subset GL(N)$ conjugacy class. Choose *marking* i.e. polynomial $P = (X - \xi_1) \dots (X - \xi_k)$ s.t $P(A) = 0$ for any $A \in \mathcal{C}$. Define



with $d_i = \text{rank}(A - \xi_1) \dots (A - \xi_{i-1})$.

- Glue those legs $\rightarrow$ star-shaped quiver associated to the connection.



One irregular singularity: core diagram [Boalch-Yamakawa 20]

Setting: only one singularity at infinity, exponential factors q_i with ramification orders r_i .

The vertices correspond to the exponential factors q_i .

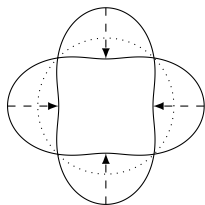
To find the number of edges/loops between q_i and q_j (one edge = one arrow in each direction)

- Count the number of Stokes arrows between the corresponding circles in the Stokes diagrams
- subtract $r_i r_j$ if $i \neq j$, or $r_i^2 - 1$ if $i = j$ (corresponds to relations in the quasi-Hamiltonian presentation)

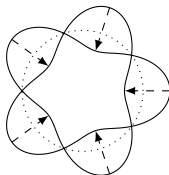
Notice there can be edges/loops with negative multiplicities

Simple examples

Untwisted example: $2 - 1 = 1$ edges between $\langle q_1 \rangle$ and $\langle q_2 \rangle$



Twisted example: $5 - 2^2 + 1 = 2$ oriented loops at $\langle q \rangle$



Properties of the quivers

Dimension of the wild character variety: $\dim \mathcal{M}_B = 2 - (\mathbf{d}, \mathbf{d})$,
with $\mathbf{d}$ dimension vector, $(\cdot, \cdot)$ bilinear form defined by the Cartan matrix
of the diagram.

In the case of regular singularities, or in the "simply-laced" case (one
untwisted irregular singularity of order ≤ 3 + simple poles), we have more:

- Quiver modularity theorem: The additive moduli space $\mathcal{M}^* \subset \mathcal{M}_B$ is
isomorphic to the Nakajima quiver variety of Γ

$$\mathcal{M}^* \simeq N(\Gamma, \mathbf{d}, \lambda).$$

- $\mathcal{M}_B$ can be seen as a multiplicative quiver variety. The several
readings correspond to isomorphisms between the wild character
varieties.
- Weyl group action: basically corresponds to exchanging eigenvalues of
formal monodromies.

Extension to several irregular singularities

Idea: use transformations to reduce to the case with one singularity at infinity.

Action of $SL_2(\mathbb{C})$ on singularity data, including Fourier-Laplace transform.

Theorem (D.)

There is a well-defined way to associate to any connection (E, ∇) on a Zariski open subset of the affine line, with modified formal data $(\check{\Theta}, \check{\mathcal{C}})$ a core diagram $\Gamma(\check{\Theta}, \check{\mathcal{C}})$ such that

- ① *If $\check{\Theta}$ has support at infinity, then $\Gamma(\check{\Theta}, \check{\mathcal{C}})$ is the diagram of [BY].*
- ② *$\Gamma(\check{\Theta}, \check{\mathcal{C}}) = \Gamma(A \cdot (\check{\Theta}, \check{\mathcal{C}}))$ for any $A \in SL_2(\mathbb{C})$.*

The diagram comes together with a dimension vector $\mathbf{d} \in \mathbb{Z}^I$, and a vector of labels $\mathbf{q} \in (\mathbb{C}^*)^I$ (where I is the set of vertices).

Explicit description

- As before, the diagram consists of a core to which we attach legs
- The vertices of the core correspond to the active circles
- The multiplicity of edges/loops in the core is given by

Definition

Suppose i, j are active circles at a_i, a_j , with $i = \langle q_i \rangle$, $j = \langle q_j \rangle$ are circles at a_i, a_j respectively. Let $\alpha_i = \text{lrr}(q_i)$, $\beta_i = \text{ram}(q_i)$ and similarly for j . If $a_i = a_j$ let us denote by B_{ij}^∞ the number of edges between i and j in the diagram of [BY].

- 1 If $a_i = a_j = \infty$ then $B_{ij} = B_{ij}^\infty$.
- 2 If $a_i = \infty$ and $a_j \neq \infty$ then $B_{ij} = B_{ji} = \beta_i(\alpha_j + \beta_j)$.
- 3 If $a_i \neq \infty, a_j \neq \infty$ and $a_i \neq a_j$ then $B_{ij} = B_{ji} = 0$.
- 4 If $a_i = a_j \neq \infty$ then $B_{ij} = B_{ji} = B_{ij}^\infty - \alpha_i\beta_j - \alpha_j\beta_i$.

Fourier-Laplace transformation

- Connections on the Riemann sphere closely related to modules on the Weyl algebra $A_1 = \mathbb{C}[z]\langle \partial_z \rangle$, with $[\partial_z, z] = 1$.
- Fourier transform: automorphism of the Weyl algebra:

$$\begin{cases} z & \mapsto -\partial_z \\ \partial_z & \mapsto z \end{cases}$$

- If M module over the Weyl algebra, Fourier transform $\Rightarrow \mathcal{F}M$
- More generally: we can act with any matrix $A \in SL_2(\mathbb{C})$.
- Such symplectic transformations are generated by 3 types of elementary transformations.
 - ▶ Fourier-Laplace
 - ▶ Scalings $z \mapsto z/\lambda$
 - ▶ Twists, corresponding to $\nabla \mapsto \nabla + \lambda z dz$.

The stationary phase formula [Malgrange 91, Fang 09, Sabbah 08]

It relates the irregular class of a connection and its Fourier transform.

- Solutions are linear combinations of terms of the form

$$f(z) = e^{q(z)} g(z),$$

- The Fourier transform is an integral $\widehat{f}(\xi) = \int_{\gamma} e^{q(z) - \xi z} g(z)$
- The behaviour of the integral when $\xi \rightarrow \infty$ is determined by the critical point of the exponential factor, i.e. z_0 such that

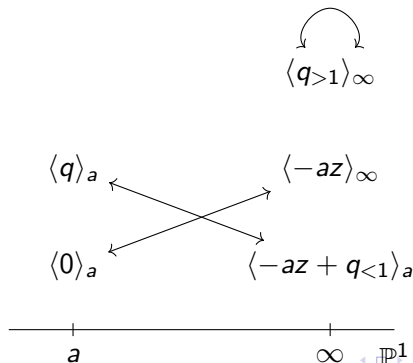
$$\frac{\partial q}{\partial z}(z_0) = \xi.$$

- New exponential factor $\widetilde{q}(\xi) = q(z_0(\xi)) - \xi z_0(\xi)$
 $\Rightarrow$ Legendre transform of q .

The stationary phase formula

Different types of circles:

- ① The *pure circles* at infinity, of the form $\langle \alpha z \rangle_\infty$, with $\alpha \in \mathbb{C}$.
- ② Other circles of slope ≤ 1 at infinity, of the form $\langle \alpha z + q \rangle_\infty$, with $\alpha \in \mathbb{C}$, and $q \neq 0$ of slope < 1 ,
- ③ Circles $\langle q \rangle_\infty$ of slope > 1 at infinity,
- ④ Irregular circles at finite distance $\langle q \rangle_a$, with $q \neq 0$, $a \in \mathbb{C}$.
- ⑤ The tame circles $\langle 0 \rangle_a$, $a \in \mathbb{C}$ at finite distance.



Main idea

Putting everything at infinity

- Apply twist at infinity $q \mapsto q + \lambda z^2$.
Now all active circles at ∞ have slope > 1 .
- Now apply Fourier: only one singularity at infinity.

Crucial fact: this does not depend on the choices. Indeed

Theorem

If Θ is an irregular class at infinity and $A \in SL_2(\mathbb{C})$ such that $A \cdot \Theta$ is also at infinity, then

$$\Gamma_c(A \cdot \Theta) = \Gamma_c(\Theta).$$

Main ingredient of proof: the diagram is invariant under Fourier transform.

Properties of the diagram

Dimension of the wild character variety: we still have

$$\dim \mathcal{M}_B(E, \nabla) = 2 - (\mathbf{d}, \mathbf{d}).$$

Idea of proof: use the quasi-Hamiltonian description of the wild character variety

$$\mathcal{M}_B(V) = \mathrm{Hom}_{\mathbb{S}}(E, \nabla) //_{\mathcal{C}} \mathbf{H}.$$

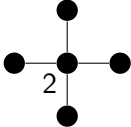
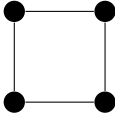
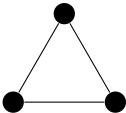
From the invariance under $SL_2(\mathbb{C})$, get several readings of the diagram

It is still possible to interpret some Weyl reflections on the diagram as coming from operations on connections

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Diagrams for Painlevé equations

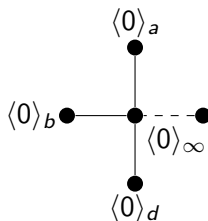
- Painlevé equations can be obtained from isomonodromic deformation equations of some meromorphic connections. Such a connection is a Lax pair for the Painlevé equation.
- They correspond to two-dimensional moduli spaces.
- Standard (rank 2) Lax pairs and diagrams:

VI	V	IV	III	II
$1+1+1+1$	$2+1+1$	$3+1$	$2+2$	4
				

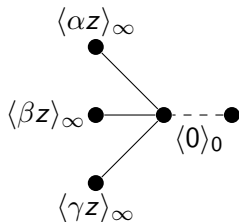
- Get affine Dynkin diagrams corresponding to the Okamoto symmetries of Painlevé equations

Painlevé VI: several Lax pairs

Standard Lax pair: rank 2, 4 simple poles (including one at infinity)



Harnad dual: rank 3, one irreg sing. at infinity, and one simple pole at finite distance



Painlevé III

From alternative Lax pair [Boalch-Yamakawa 20]: one pole of order 2, twisted, 2 simple poles, rank 2

A diagram representing the alternative Lax pair for Painlevé III. It consists of a horizontal line with three black dots. The left dot is labeled $\langle 0 \rangle_a$ and the right dot is labeled $\langle 0 \rangle_a$. The middle dot is labeled $\langle \alpha z^{1/2} \rangle_\infty$ below it. A dashed oval encircles the middle dot and the line segment between it and the right dot.

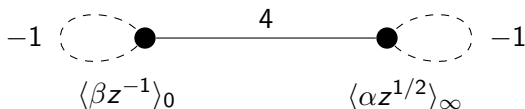
From the usual Lax pair: two irregular poles of order 2

A diagram representing the usual Lax pair for Painlevé III. It consists of a horizontal line with three black dots. The left dot is labeled $\langle \lambda_1 z \rangle_\infty$ and the right dot is labeled $\langle \lambda_2 z \rangle_\infty$. The middle dot is labeled $\langle \mu z^{-1} \rangle_0$ below it. A dashed oval encircles the middle dot and the line segment between it and the right dot.

Degenerate Painlevé equations

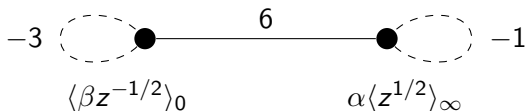
Degenerate PIII: two second order poles, one of them twisted (slope 1/2).

$$\check{\Theta} = \langle \alpha z^{1/2} \rangle_{\infty} + \langle \beta z^{-1} \rangle_0$$



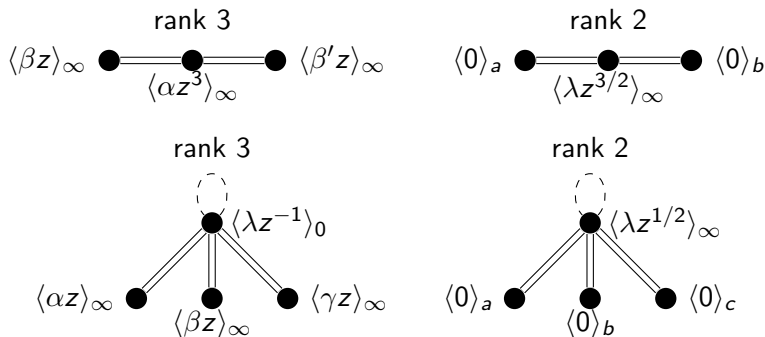
Doubly degenerate PIII: two twisted second order poles (slopes 1/2)

$$\check{\Theta} = \langle \alpha z^{1/2} \rangle_{\infty} + \langle \beta z^{-1/2} \rangle_0$$



Higher dimensional cases

Among 4-dimensional Painlevé-type equations, two examples of different known Lax pairs corresponding to the same diagram:



Proposition

All Lax pairs for a same given 4-dimensional Painlevé-type equation listed by [\[Kawakami et al. 18\]](#) correspond to different readings of the same diagram.

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Orbits under basic operations

Basic operations:

- $SL_2(\mathbb{C})$
- Twists
- Möbius transformations (change the point at infinity)

The diagram is invariant under $SL_2(\mathbb{C})$ but not under other basic operations.

Question: orbits under repeated application of basic operations? Is there a minimal diagram?

In the rigid case:

Theorem (Katz '96, Arinkin '10)

Any rigid irreducible connection on $\mathbb{P}^1$ can be brought to the trivial rank one connection by repeated application of basic operations.

Diagrams with one vertex and few loops

Consider an exponential factor $\langle q \rangle$ at infinity, let $k = B_{\langle q \rangle, \langle q \rangle} / 2$ the corresponding number of loops.

Theorem

If $k = 0, 1, 2$, $\langle q \rangle$ can be brought by repeated application of basic operations to an exponential factor which is an admissible deformation of

k	simple exponential factor
0	0
1	$z^{5/3}$
2	$z^{7/5}$

Towards explicit isomorphisms?

- One expects the different readings of the diagram correspond to isomorphisms of the moduli spaces, compatible with the symplectic structure
- Such isomorphisms should be obtained by the transformation of Stokes data under Fourier transform
- This has been studied by many people [\[Malgrange, Mochizuki, Sabbah, Hien, D'Agnolo...\]](#), but difficult to get explicit formulas beyond the simply laced case.
- A simple twisted case : connections with only 1 exponential factor $\langle z^{\alpha/\beta} \rangle_{\infty}$ and $\langle z^{\alpha/\alpha-\beta} \rangle_{\infty}$ with $\alpha > \beta$.

Isomorphism for slopes $5/2 \leftrightarrow 5/3$



- $\mathcal{M}_B = \{s_1, \dots, s_5, hS_5S_4S_3S_2S_1 = 1\}$ with
 $h = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, S_1 = \begin{pmatrix} 1 & s_1 \\ 0 & 1 \end{pmatrix}, S_2 = \begin{pmatrix} 1 & 0 \\ s_2 & 1 \end{pmatrix}, \dots$
- $\mathcal{M}'_B = \{t_1, \dots, t_{10}, kT_{10} \dots T_1 = 1\}$ with
 $k = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}, T_1 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ t_1 & 0 & 1 \end{pmatrix}, T_2 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & t_2 & 1 \end{pmatrix}, \dots$
- $\mathcal{M}_B \simeq \mathcal{M}'_B$ via
 $\Phi : (s_1, s_2, s_3, s_4, s_5) \mapsto (s_3, -s_5, -s_2, s_4, s_1, -s_3, s_5, s_2, -s_4, -s_1)$
- The isomorphism is compatible with the symplectic structures