Enumeration of planar maps: from the quadratic method to quantum gravity

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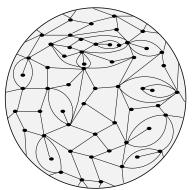


Tutte Centenary Conference Cambridge, 11 July 2017

Outline

- 1 The beginning: Tutte's recursive approach
- ② A detour in physics: matrix models and 2D quantum gravity
- 3 The bijective approach: labeled trees and the Brownian map
- Beyond the Brownian map

Planar map: definition



A planar map is a connected (multi)graph embedded in the sphere, considered up to continuous deformation. It is made of vertices, edges and faces.



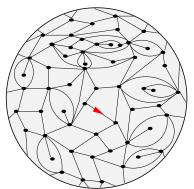








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A rooted map has a distinguished oriented edge. A pointed map has a distinguished vertex.









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Tutte's "Census" papers (1962-63)

A CENSUS OF PLANAR TRIANGULATIONS

W T THITTE

1. Triangulations. Let P be a closed region in the plane bounded by a simple closed curve, and let S be a simplicial dissection of P. We may say that S is a dissection of P into a finite number α of triangles so that no vertex of any one triangle is an interior point of an edge of another. The triangles are "topological" triangles and their edges are closed arcs which need not be straight segments. No two distinct edges of the dissection ion the same two vertices, and no two triangles have more than two vertices in common.

There are $k \ge 3$ vertices of S in the boundary of P, and they subdivide this boundary into k edges of S. We call these edges external and the remaining edges of S, if any, internal. If r is the number of internal edges we have

$$3\alpha = 2r + k$$

$$r \equiv k \pmod{3}$$
.

Let us call S a triangulation of P if it satisfies the following condition: no internal edge of S has both its ends in the boundary of P. We note that in the case k = 3 every simplicial dissection is a triangulation.

A CENSUS OF SLICINGS

W. T. TUTTE

1. Introduction. A band is a closed connected set in the 2-sphere. bounded by one or more disjoint simple closed curves.

Consider a band B with bounding curves J_1, J_2, \dots, J_k . On each curve J_i let there be chosen $m_i \geqslant 0$ points to be called vertices, with the restriction that the sum of the k integers m_i is to be even. Write

(1)
$$\sum_{i=1}^{k} m_i = 2n$$
.

Next consider a set of n disjoint open arcs in the interior of B which join the 2n vertices in pairs and partition the remainder of the interior of B into simply connected domains. We call the resulting dissection of B a slicing with respect to the given set of vertices. The arcs are the internal edges of the slicing and the simply connected domains are its internal faces, or slices.

A CENSUS OF HAMILTONIAN POLYGONS

W T THTTE

Summary. In this paper we deal with trivalent planar maps in which the boundary of each country (or "face") is a simple closed curve. One vertex is distinguished as the root and its three incident edges are distinguished as the first, second, and third major edges. We determine the average number of Hamiltonian polygons, passing through the first and second major edges, in such a "rooted map" of 2n vertices. Next we consider the corresponding problem for 3-connected rooted maps. In this case we obtain a functional equation from which the average can be computed for small values of n.

1. Rooted maps. For the purposes of this paper a blanar map M is a representation of the 2-sphere (or closed plane) as a union of a finite number of disjoint point-sets called cells. The cells are of three kinds, vertices, edges, and faces, said to have dimension 0, 1, and 2 respectively. Each vertex consists of a single point. Each edge is an open arc whose ends are distinct vertices. Each face is a simply connected domain whose boundary is a simple closed curve made up of edges and vertices. We denote the numbers of cells, vertices, edges, and faces of M by C(M), V(M), E(M), and F(M) respectively.

A CENSUS OF PLANAR MAPS

W. T. TUTTE

1. Introduction. In the series of "Census" papers, of which this is the fourth, we attempt to lay the groundwork of an enumerative theory of planar maps (12, 13, 14). The maps concerned are rooted in the sense that some edge is fixed as the root, and a positive sense of description and right and left sides are specified for it. This device simplifies the theory by ruling out the possibility of a map being symmetrical.

In this paper formulae are obtained for the number of rooted maps (with n edges), the number of non-separable rooted maps, and the number of 3connected rooted maps without multiple joins (called c-nets). Some similar enumerations, supplementing the results of earlier papers, are given for triangulations and bicubic maps.

Let M_n be the number of rooted planar maps with n edges.

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Let $M_{n,k}$ be the number of rooted planar maps with n edges and outer degree k. Now we can do something.

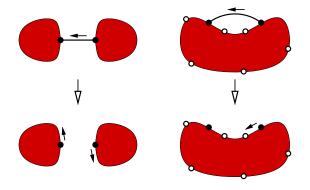
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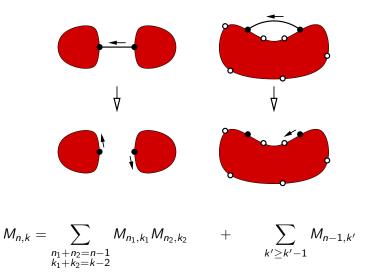
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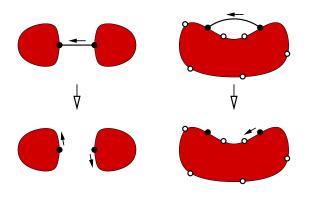
We will find and solve an equation satisfied by the generating function

$$M(x,y) := \sum_{n,k>0} M_{n,k} x^n y^k$$

though we are ultimately interested only in M(x,1) (y is a "catalytic" variable).

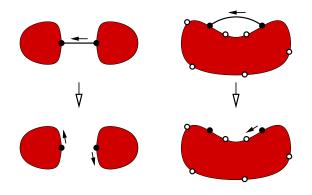






$$M_{n,k} = \sum_{\substack{n_1 + n_2 = n - 1 \\ k_1 + k_2 = k - 2}} M_{n_1,k_1} M_{n_2,k_2} + \sum_{k' \ge k - 1} M_{n-1,k'}$$

(holds for $n \ge 1$, with $M_{0,0} = 1$, $M_{0,k} = 0$ for $k \ge 1$)



$$M(x,y) = 1 + xy^{2}M(x,y)^{2} + xy\frac{yM(x,y) - M(x,1)}{y - 1}$$

$$M(x, y)$$
 satisfies

$$\Phi(M(x,y), M(x,1), x, y) = 0$$
 (1)

where

$$\Phi(M, m, x, y) = xy^2(1 - y)M^2 + (y - 1 - xy^2)M + 1 - y + xym.$$

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By differentiating (1) wrt y, we get

$$\underbrace{\frac{\partial \Phi}{\partial M}(M(x,y),M(x,1),x,y)}_{\text{(2)}}\underbrace{\frac{\partial M}{\partial y}(x,y)} + \underbrace{\frac{\partial \Phi}{\partial y}(M(x,y),M(x,1),x,y)}_{\text{(3)}} = 0.$$

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Let Y(x) be a series such that (2) vanishes when we substitute y = Y(x) (it exists and is unique!).

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Let Y(x) be a series such that (2) vanishes when we substitute y = Y(x) (it exists and is unique!). Then (3) vanishes too. Combined with (1), we get 3 equations for 4 parameters, hence by elimination we may deduce an algebraic equation relating M(x,1) and x.

In practice, we even get a rational parametrization

$$x = \frac{(Y-1)(3-2Y)}{Y^2}, \qquad M(x,1) = \frac{Y(4-3Y)}{(3-2Y)^2}$$

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Theorem (Census of planar maps, 1963)

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Some remarks:

- ullet for many classes of maps, we get a similar quadratic equation \Longrightarrow "quadratic method" (Brown, 1965)
- nowadays there is a general method to solve polynomial equations with one catalytic variable (Bousquet-Mélou and Jehanne, 2006)
- asymptotically $M_n \sim C12^n n^{-5/2}$, the -5/2 exponent is universal.

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G. 't Hooft



É. Brézin, C. Itzykson, G. Parisi, J.-B. Zuber

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Planar Diagrams

E. Brézin, C. Itzykson, G. Parisi*, and J. B. Zuber Service de Physique Théorique, Centre d'Études Nucléaires de Saclay, F-91190 Gif-sur-Yvette, France

Abstract. We investigate the planar approximation to field theory through the limit of a large internal symmetry group. This yields an alternative and powerful method to count planar diagrams. Results are presented for cubic and quartic vertices, some of which appear to be new. Quantum mechanics treated in this approximation is shown to be equivalent to a free Fermi gas system.

In 1978, BIPZ considered the special case of zero spatial dimensions, where the gauge field reduces to a simple $N \times N$ matrix.

The problem is to estimate the large N asymptotics of an integral over $N \times N$ hermitian matrices.

3. Combinatorics of Quartic Vertices

1) Vacuum Diagrams

Setting each diagram equal to unity, apart from the overall weight, is equivalent to treat a field theory in zero dimension, in which space-time is reduced to one or to a finite number of points. It means that

$$\exp -N^2 E^{(0)}(g) = \lim_{N \to \infty} \int d^{N^2} M \exp \left[-\int_{1}^{\infty} \text{tr } M^2 + \frac{g}{N} \text{tr } M^4 \right]. \tag{3}$$

The integration measure on hermitian matrices is

$$d^{N^2}M \equiv \prod_i dM_{ii} \prod_{i < j} d(\operatorname{Re} M_{ij}) d(\operatorname{Im} M_{ij})$$
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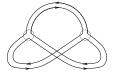
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Such an integral may be viewed as a perturbation of a Gaussian integral (GUE partition function). By a formal expansion of the non Gaussian term, we get a perturbative series which is standardly represented via Feynman diagrams.

In the context of $N \times N$ matrices, the Feynman diagrams can be represented as "fatgraphs" or "ribbon graphs" (because a matrix carries two indices). For instance, the diagrams contributing to the expansion of $\langle ({\rm Tr} M^3)^2 \rangle$ are:





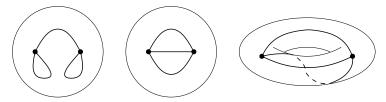


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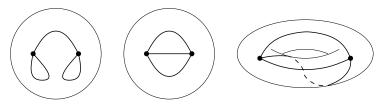
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When *N* gets large, the dominant diagrams correspond to planar maps!

Table 1. Counting rules for the vacuum amplitude $E^{(0)}(g)$ in the planar limit, up to order three

2g	⊖ 2g²	∞ 16 g ²	$\bigotimes_{\frac{32}{3}g^3}$
64 g ³	0000 128g ³	256 3 g³	

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38 E. Brézin et al.

Since the integrand (3) depends only on the eigenvalues λ_i , this allows us to integrate over U and, up of to a g-independent normalizing factor we obtain

$$\exp -N^2 E^{(0)}(g) = \lim_{N \to \infty} \int \prod_{i} d\lambda_i \prod_{i < j} (\lambda_i - \lambda_j)^2 \exp \left[\frac{1}{2} \sum_{i} \lambda_i^2 + \frac{g}{N} \sum_{i} \lambda_i^4 \right]. \tag{6}$$

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The steepest descent method consists in understanding the large N-limit of the empirical distribution of eigenvalues $\mu_N := \frac{1}{N} \sum_{i=1}^N \delta_{\lambda_i}$, by assuming (or showing) that it converges to a continuous distribution μ solving a minimization problem.

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The k-th moment of the limiting empirical distribution μ corresponds to the generating function of rooted planar maps with outer degree k.

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- the *m*-th Catalan number $c_m = \frac{(2m)!}{m!(m+1)!}$,
- i.e. the number of rooted plane trees with m edges,
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To count general (non unicellular) maps we need to study the non Gaussian $(g \neq 0)$ case. Nowadays, this is routinely done using the so-called loop or Schwinger-Dyson equations for the resolvent

$$F(z) := \int_{\text{supp}\mu} \frac{d\mu(x)}{x - z} = \sum_{k=0}^{\infty} \frac{1}{z^{k+1}} \int x^k d\mu(x)$$

which are basically equivalent to Tutte's recursive equations (z is the catalytic variable).

"There are methods and formulae in science, which serve as master-keys to many apparently different problems. The resources of such things have to be refilled from time to time. In my opinion at the present time we have to develop an art of handling sums over random surfaces."

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His paper introduces Liouville quantum gravity:

$$Z = \int D\varphi(\xi) \exp\left(-\frac{26-D}{48\pi} \int_{\xi} \left[\frac{1}{2} (\partial_{\mu}\varphi)^2 + \mu^2 e^{\varphi}\right]\right). \tag{25}$$

Here $D\varphi(\xi)$ denotes the (ill-defined!) functional integration over a field $\varphi: \Sigma \to \mathbb{R}$, with Σ a Riemann surface.

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In 1984-1985, several papers suggested another approach: discretize quantum gravity using random triangulations.

Liouville quantum gravity should emerge for a very large number of vertices and a very small mesh size. The computations require being able to enumerate maps with various constraints, and matrix models were used as an efficient tool.



Picture by N. Curien



V. Kazakov



F. David



J. Ambjørn, B. Durhuus, J. Fröhlich

A highlight of this approach is the "solution" of the 2D Ising model on random triangulations by Kazakov in 1986, using a two-matrix model.

Combinatorially, it corresponds to enumerating planar triangulations endowed with a (non necessarily proper) 2-coloring, fixing the number of vertices and the number of monochromatic edges. (For proper colorings, we recover Tutte's "chromatic sum" at $\lambda=2$.)



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Quite surprisingly, the resulting critical exponents were in complete agreement with predictions from Liouville quantum gravity. The Potts model (q-colorings for general q) was similarly "solved" via matrix models a few years later (Daul, Zinn-Justin, Bonnet, Eynard...), still consistently with LQG.

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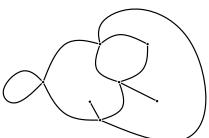
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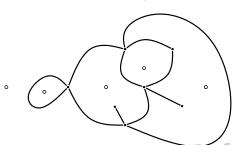
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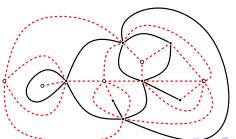
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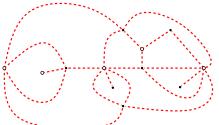
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Let us also observe that a planar quadrangulation with n edges has n+2 vertices, therefore

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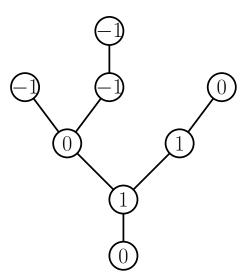
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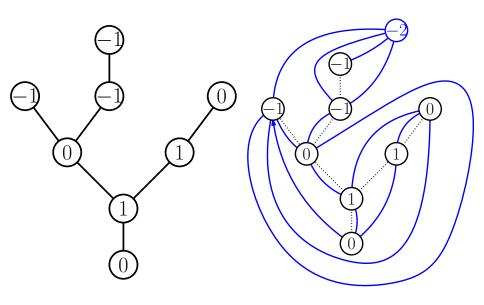
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Idea: let us consider a labeling ℓ of the vertices of the tree by integers, such that

- $\ell(\text{root}) = 0$,
- $|\ell(u) \ell(v)| \neq 1$ if u, v are neighbours.





The CVS bijection has the following interesting property: let v^* be the added vertex, then for any vertex v of the tree we have

$$d_Q(v,v^*) = \ell(v) - \min \ell + 1$$

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It is not difficult to see that, for a uniform random quadrangulation with n faces, the distances should be of order $n^{1/4}$ (the branches of the trees have length $\propto n^{1/2}$, on each branch the labels form a lazy random walk on \mathbb{Z}).

This was made more precise in several works. Chassaing and Schaeffer proved in 2002 that

$$\frac{1}{n^{1/4}}\max d_{Q_n}(\cdot, v^*) \xrightarrow[n \to \infty]{(d)} w_{\mathrm{ISE}}$$

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Of course, v^* plays no special role, it is just a uniformly chosen vertex in the quadrangulation.

With Di Francesco and Guitter, I obtained in 2003 an exact expression for the generating function $R_{\ell}(t)$ of planar quadrangulations with two marked vertices at distance $\leq \ell$. It reads

$$R_{\ell} = R \frac{(1 - x^{\ell})(1 - x^{\ell+3})}{(1 - x^{\ell+1})(1 - x^{\ell+2})}$$

where the series R = R(t) and x = x(t) satisfy

$$R = 1 + 3tR^2$$
, $x + \frac{1}{x} + 1 = \frac{1}{tR^2}$.

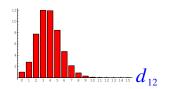
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 $a \rightarrow a \rightarrow n = 50$

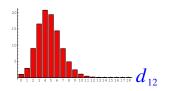
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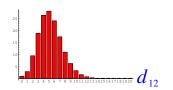
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Consequences of the CVS bijection

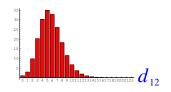
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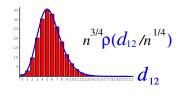
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This encodes the exact law for the distance between two uniform points in a random quadrangulation of size n.

The limit $n \to \infty$ can be analyzed by standard analytical methods, and we confirm a prediction of Ambjørn and Watabiki (1996).



Beyond the CVS bijection

The CVS bijection encodes the distances from only one specific vertex. Can we go beyond this limitation?

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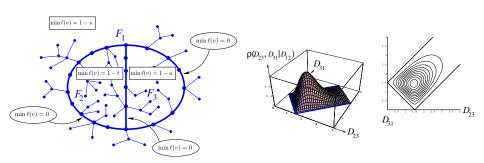
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In parallel, there has been an ongoing effort by probabilists to define a random continuous object which is to random quadrangulations what Brownian motion is to random walks, i.e. a scaling limit.

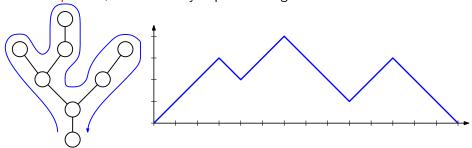
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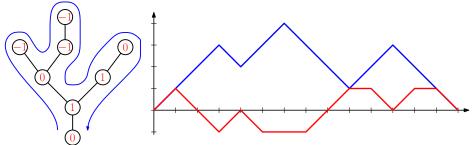
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Recall that a rooted plane tree with n edges can be described by its contour process, which is a Dyck path of length 2n.

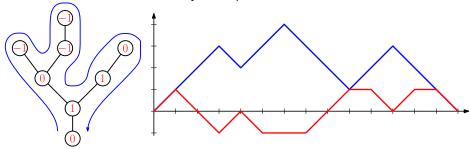


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Theorem (Chassaing-Schaeffer, 2002)

By appropriate rescalings, the contour and the label processes of a uniform random labeled tree with n edges converge jointly to a continuous process (e,Z) called the Brownian snake.

Now, how do we construct a "map" from the Brownian snake? We need to find a continuum analogue of the CVS bijection. (The analogue of the tree/Dyck path correspondance already allows to construct the Brownian Continuum Random Tree.)

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(Several properties proved in the meantime: BM is homeomorphic to the sphere, Hausdorff dimension is 4, geodesics, etc.)

Outline

- The beginning: Tutte's recursive approach
- A detour in physics: matrix models and 2D quantum gravity
- The bijective approach: labeled trees and the Brownian map
- 4 Beyond the Brownian map

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Is there nothing else to see? Since Kazakov's solution, we know that for maps endowed with a critical Ising model,

$$Z_n^{\text{Ising}}(\nu_c) \sim C \kappa^n n^{-7/3}$$
.

(For the q-state Potts model, the exponent varies with q.)

Slicings

We again find inspiration in the work of Tutte:

Theorem (Census of slicings, 1962)

The number of bipartite rooted planar maps with e edges and f faces, among which f_i have degree 2i for all i, is

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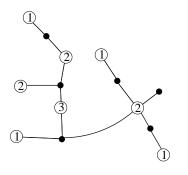
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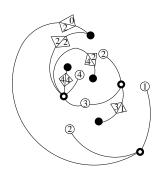
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Schaeffer gave a bijective proof of this formula in 1997, using a different sort of trees. In 2004, with Di Francesco and Guitter, we generalized the CVS bijection to this setting (and more). This is now known as the BDG bijection.

Mobiles





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To reach new universality classes, we need to consider Boltzmann random maps: to a bipartite map $\mathfrak m$ we assign a probability

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where $(q_i)_{i\geq 1}$ is a given sequence of face weights. By the BDG bijection, we obtain a certain distribution over mobiles (labeled trees), where q_i controls the proportion of vertices of degree i. The structure of such a tree is best described by its Lukasiewicz path (instead of contour), as q_i controls the probability to jump up by i-1.

By an appropriate choice of the q_i 's, we can ensure that the Lukasiewicz path makes "large jumps" in such a way that its scaling limit is not a Brownian excursion, but another Lévy process.

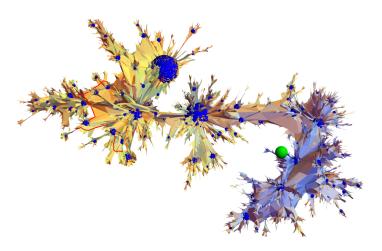
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Theorem (Le Gall-Miermont, 2011)

For any $d \in (2,4)$, there exists nongeneric sequences of face weights $(q_i)_{i\geq 1}$ for which the corresponding Boltzmann random maps of size n have distances of order $n^{1/d}$.

After rescaling, these maps converge (along subsequences) to a random compact metric space of Hausdorff dimension d.

By a continuous analogue of the BDG bijection, Le Gall and Miermont also constructed the conjecturally unique limit, the stable map.

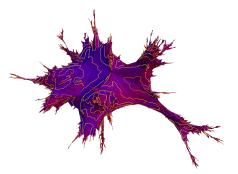


Picture by T. Budd of the dual of a stable map

Can we connect this with physics?

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Yes, Le Gall and Miermont conjectured that stable maps are connected with the so-called O(n) loop model on random maps (related to FK percolation, n=1 corresponds to percolation and the Ising model).



Picture by J. Bettinelli

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Latest developments: results about the nestings of loops, that connect with Liouville quantum gravity (Borot-B-Duplantier 2016, Chen-Curien-Maillard 2017).

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This is just my own perspective and I certainly forgot many other deep aspects, e.g.:

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- much work on the connection between random maps and Liouville quantum gravity – see the recent preprint of Gwynne-Miller-Sheffield on Tutte embedding of the mated-CRT map.

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To conclude let me quote a colleague of Tutte: "If I have seen further it is by standing on the shoulders of giants."