

1. LÉVY'S CRITERION OF CONVERGENCE IN LAW

- (1) For every $t \in \mathbb{R}$, $\omega \in \Omega \mapsto e^{itX(\omega)}$ is measurable and bounded in module by 1, whence integrable, so $\phi(t) = \mathbb{E}[e^{itX(\omega)}]$ is well-defined. Then, for every t_0 fixed in \mathbb{R} ,

$$\lim_{t \rightarrow t_0} e^{itX} = e^{it_0X}$$

almost surely, and all these random variables have real and imaginary part uniformly bounded by 1. By Lebesgue dominated convergence,

$$\lim_{t \rightarrow t_0} \phi(t) = \lim_{t \rightarrow t_0} \mathbb{E}[e^{itX}] = \mathbb{E}[e^{it_0X}] = \phi(t_0),$$

so ϕ is continuous on \mathbb{R} .

- (2) All the functions considered are measurable on $\mathbb{R} \times [-\varepsilon, \varepsilon]$ and uniformly bounded, and integrated against a finite measure, namely, $\mu \otimes dt$. So, one can use freely Fubini's theorem to compute:

$$\begin{aligned} I_\varepsilon &= \int_{-\varepsilon}^{\varepsilon} (1 - \phi(t)) dt = \iint_{\mathbb{R} \times [-\varepsilon, \varepsilon]} (1 - e^{itx}) \mu(dx) dt \\ &= \int_{\mathbb{R}} \left(2\varepsilon - \frac{2 \sin(\varepsilon x)}{x} \right) \mu(dx) = 2\varepsilon \left(\int_{\mathbb{R}} \left(1 - \frac{\sin(\varepsilon x)}{\varepsilon x} \right) \mu(dx) \right). \end{aligned}$$

Set $y = \varepsilon x$; for $|y| \geq 2$, one has indeed

$$\frac{\sin y}{y} \leq \frac{|\sin y|}{2} \leq \frac{1}{2} \quad ; \quad 1 - \frac{\sin(\varepsilon x)}{\varepsilon x} \geq \frac{1}{2}.$$

Therefore,

$$\begin{aligned} \frac{I_\varepsilon}{2\varepsilon} &\geq \int_{\mathbb{R}} \left(1 - \frac{\sin(\varepsilon x)}{\varepsilon x} \right) \mathbf{1}_{|\varepsilon x| \geq 2} \mu(dx) \geq \frac{1}{2} \int_{\mathbb{R}} \mathbf{1}_{|\varepsilon x| \geq 2} \mu(dx) \\ \frac{I_\varepsilon}{\varepsilon} &\geq \mu(\{x \mid |\varepsilon x| \geq 2\}). \end{aligned}$$

- (3) Let $\eta > 0$. By continuity at 0 of ϕ , $|1 - \phi(t)| = |\phi(0) - \phi(t)| \leq \frac{\eta}{3}$ for t small enough, say smaller than some $\varepsilon > 0$. Then,

$$\frac{I_\varepsilon}{\varepsilon} \leq \frac{1}{\varepsilon} \int_{-\varepsilon}^{\varepsilon} \frac{\eta}{3} dt = \frac{2\eta}{3}.$$

By Lebesgue dominated convergence,

$$\lim_{n \rightarrow \infty} \frac{I_{\varepsilon, n}}{\varepsilon} = \frac{I_\varepsilon}{\varepsilon},$$

so for N big enough and all $n \geq N$,

$$\frac{I_{\varepsilon, n}}{\varepsilon} \leq \frac{3I_\varepsilon}{2\varepsilon} \leq \eta.$$

By Question (2), this implies that for all $n \geq N$,

$$\mu_n \left((-2\varepsilon^{-1}, 2\varepsilon^{-1})^c \right) \leq \eta.$$

By decreasing the value of ε , one can also assume that this is true for the n 's smaller than N (there is only a finite number of measures to consider). Thus, for every $\eta > 0$, there is a relatively compact set (bounded interval) in \mathbb{R} with complementary of measure smaller than η for all measure in $(\mu_n)_{n \in \mathbb{N}}$; *i.e.*, there is tightness.

- (4) The sequence of measures $(\mu_n)_{n \in \mathbb{N}}$ is tight, hence relatively compact; let $(\mu_{\psi(n)})_{n \in \mathbb{N}}$ be a convergent subsequence, and μ its limit. For every $t, x \mapsto e^{itx}$ is a continuous bounded function,

$$\mu(e^{itx}) = \int_{\mathbb{R}} e^{itx} \mu(dx) = \lim_{n \rightarrow \infty} \mu_{\psi(n)}(e^{itx}) = \lim_{n \rightarrow \infty} \phi_{\psi(n)}(t) = \phi(t).$$

Hence, ϕ is indeed the characteristic function of a probability measure on \mathbb{R} .

- (5) Again, $\mu \otimes \nu$ is a finite measure on $\mathbb{R} \times \mathbb{R}$, and the functions under consideration are bounded, so one can use Fubini's theorem.

$$\begin{aligned} \int_{\mathbb{R}} e^{-itx} \phi_{\mu}(x) \nu(dx) &= \iint_{\mathbb{R} \times \mathbb{R}} e^{-itx} e^{iyx} \mu(dy) \nu(dx) \\ &= \int_{\mathbb{R}} \left(\int_{\mathbb{R}} e^{i(y-t)x} \nu(dx) \right) \mu(dy) \\ &= \int_{\mathbb{R}} \phi_{\nu}(y-t) \mu(dy). \end{aligned}$$

- (6) Let f be a bounded non-negative continuous function on \mathbb{R} . One has

$$\begin{aligned} \int_{\mathbb{R}} f(t) D(\varepsilon, t) dt &= \iint_{\mathbb{R}^2} f(t) e^{-\frac{\varepsilon(t-y)^2}{2}} \mu(dy) dt \\ &= \iint_{\mathbb{R}^2} f(x+y) e^{-\frac{\varepsilon x^2}{2}} \mu(dy) dx \\ &= \sqrt{2\pi\varepsilon^{-1}} \iint_{\mathbb{R}^2} f(x+y) \mu(dy) \nu_{\varepsilon^{-1}}(dx) \\ &= \sqrt{2\pi\varepsilon^{-1}} \mathbb{E}[f(X + Y_{\varepsilon^{-1}})]. \end{aligned}$$

- (7) For every interval $(-\eta, \eta)$ around 0,

$$\mathbb{P}[Y_{\varepsilon^{-1}} \in (-\eta, \eta)^c] \leq \frac{\mathbb{E}[(Y_{\varepsilon^{-1}})^2]}{\eta^2} = \frac{1}{(\varepsilon\eta)^2} \rightarrow 0,$$

so one has indeed $Y_{\varepsilon^{-1}} \rightarrow_{\mathbb{P}} 0$. As a consequence, the pair $(X, Y_{\varepsilon^{-1}})$ converge in probability to $(X, 0)$, and also in law. Since convergence in law is compatible with composition by continuous functions,

$$X + Y_{\varepsilon^{-1}} \rightarrow X + 0 = X.$$

- (8) Suppose ϕ_{μ} known. Then, for every $\varepsilon > 0$ and every bounded continuous function f ,

$$\mathbb{E}[f(X + Y_{\varepsilon^{-1}})] = \frac{1}{\sqrt{2\pi\varepsilon^{-1}}} \int_{\mathbb{R}} f(t) D(\varepsilon, t) dt = \frac{1}{\sqrt{2\pi\varepsilon^{-1}}} \iint_{\mathbb{R}^2} e^{-itx} \phi_{\mu}(x) \nu_{\varepsilon}(dx) dt$$

is also known, since it can be computed from f and ϕ_{μ} . Taking the limit as ε goes to infinity, one gets back $\mathbb{E}[f(X)]$ by the previous Question, that is to say $\mu(f)$. So, ϕ_{μ} determines $\mu(f)$ for every $f \in \mathcal{C}^b(\mathbb{R})$, and therefore μ by classical arguments of measure theory.

- (9) We have seen that $(\mu_n)_{n \in \mathbb{N}}$ was tight under the hypotheses of Lévy's criterion, and the limit of a convergent subsequence $(\mu_{\psi(n)})_{n \in \mathbb{N}}$ is by the previous Question the unique probability measure μ such that $\phi(t) = \mu(e^{itx})$. Therefore, $\mu_n \rightarrow \mu$.

- (10) This is clear since $(t, x) \mapsto e^{itx}$ is a continuous bounded function.

- (11) The characteristic function of $Y_k = B_k - \frac{1}{k}$ is

$$\mathbb{E}[e^{itY_k}] = e^{-\frac{it}{k}} \mathbb{E}[e^{itB_k}] = \left(1 + \frac{e^{it} - 1}{k} \right) e^{-\frac{it}{k}}.$$

Therefore, since the Y_k 's are independent,

$$\mathbb{E}[e^{it\tilde{X}_n}] = \prod_{k=1}^n \left(1 + \frac{e^{\frac{it}{\log n}} - 1}{k} \right) e^{-\frac{it}{k \sqrt{\log n}}}$$

The Taylor expansion of each term $T_k(t)$ is

$$\begin{aligned} &\left(1 + \frac{1}{k} \left(\frac{it}{\sqrt{\log n}} - \frac{t^2}{2 \log n} + O\left(\frac{t^3}{(\log n)^{3/2}} \right) \right) \right) \left(1 - \frac{it}{k \sqrt{\log n}} + O\left(\frac{t^2}{k^2 \log n} \right) \right) \\ &= 1 - \frac{t^2}{2k \log n} + O\left(\frac{t^2}{k^2 \log n} + \frac{t^3}{k(\log n)^{3/2}} \right), \end{aligned}$$

with a O that is uniform because it comes only from the Taylor expansion around 0 of the exponential. So,

$$\begin{aligned}\log \mathbb{E}[e^{it\tilde{X}_n}] &= \sum_{k=1}^n \log \left(1 - \frac{t^2}{2k \log n} + O\left(\frac{t^2}{k^2 \log n} + \frac{t^3}{k(\log n)^{3/2}}\right) \right) \\ &= \sum_{k=1}^n -\frac{t^2}{2k \log n} + O\left(\frac{t^4}{k^4(\log n)^2} + \frac{t^2}{k^2 \log n} + \frac{t^3}{k(\log n)^{3/2}}\right) \\ &= -\frac{t^2}{2} + O\left(\frac{t^2}{\log n} + \frac{t^4}{(\log n)^2} + \frac{t^3}{(\log n)^{1/2}}\right) = -\frac{t^2}{2} + o(1)\end{aligned}$$

since $\sum_{k=1}^n \frac{1}{k^2}$ and $\sum_{k=1}^n \frac{1}{k^4}$ are convergent. So, $\mathbb{E}[e^{it\tilde{X}_n}]$ converges pointwise to $e^{-\frac{t^2}{2}}$, and since $X_n - \tilde{X}_n$ is a deterministic $O((\log n)^{-1/2})$, the same is true for $\mathbb{E}[e^{itX_n}]$. As $e^{-\frac{t^2}{2}}$ is the characteristic function of a standard Gaussian, the central limit theorem for X_n is proven.

2. MOD-CONVERGENCE AND BERRY-ESSEEN ESTIMATES

(1) By hypothesis, for t fixed,

$$\begin{aligned}\mathbb{E}[e^{itX_n}] &= \phi_n\left(\frac{t}{\sqrt{\lambda_n}}\right) = e^{-\frac{t^2}{2}} \psi_n\left(\frac{t}{\sqrt{\lambda_n}}\right) \\ &= e^{-\frac{t^2}{2}} \left(\psi\left(\frac{t}{\sqrt{\lambda_n}}\right) + o\left(\frac{1}{\sqrt{\lambda_n}}\right) \right) \\ &= e^{-\frac{t^2}{2}} \left(\psi(0) + o\left(\frac{1}{\sqrt{\lambda_n}}\right) \right) \rightarrow e^{-\frac{t^2}{2}}.\end{aligned}$$

By Lévy criterion of convergence in law, $X_n \rightharpoonup \mathcal{N}_{(0,1)}$.

(2) Suppose $\mu_n \rightharpoonup \mu$. Since μ is supposed absolutely continuous w.r.t. Lebesgue measure, each interval $I = (-\infty, x)$ is a continuity set for μ :

$$\mu(I^0) = \mu(I) \quad \text{as } I^0 = I \quad ; \quad \mu(\bar{I}) = \mu(I) + \mu(\{x\}) = \mu(I) \quad \text{since } \mu(\{x\}) = 0.$$

Therefore, by Portmanteau's theorem, $\lim_{n \rightarrow \infty} \mu_n(-\infty, x) = \mu(-\infty, x)$, so F_{μ_n} converges pointwise to F_μ . By Dini's theorem, since the cumulative functions are non-decreasing functions bounded from above by 1 and from below by 0, this is in fact a uniform convergence, so

$$d(\mu_n, \mu) = \sup_{x \in \mathbb{R}} |F_{\mu_n}(x) - F_\mu(x)| \rightarrow 0.$$

Conversely, suppose $d(\mu_n, \mu) \rightarrow 0$, and consider an open set $U \subset \mathbb{R}$. It can be written as a countable union of disjoint open intervals: $U = \bigsqcup_{k \in \mathbb{N}} (a_k, b_k)$. Therefore,

$$\mu(U) = \sum_{k \in \mathbb{N}} \mu(a_k, b_k) = \sum_{k=0}^{\infty} F_\mu(b_k) - F_\mu(a_k),$$

using for the second equality the fact that $\mu(\{a_k\}) = 0$ for all k . Fix $\varepsilon > 0$, and K such that

$$\sum_{k=0}^K F_\mu(b_k) - F_\mu(a_k) \geq \mu(U) - \varepsilon.$$

By continuity of F_μ , one can then choose $\eta > 0$ such that

$$\sum_{k=0}^K F_\mu(b_k) - F_\mu(a_k + \eta) \geq \mu(U) - 2\varepsilon;$$

notice that $\bigsqcup_{k=0}^K [a_k + \eta, b_k) \subset U$. Since F_{μ_n} converges to F_μ , one has then

$$\begin{aligned}\liminf_{n \rightarrow \infty} \mu_n(U) &\geq \liminf_{n \rightarrow \infty} \left(\sum_{k=0}^K \mu_n([a_k + \eta, b_k)) \right) = \liminf_{n \rightarrow \infty} \left(\sum_{k=0}^K F_{\mu_n}(b_k) - F_{\mu_n}(a_k + \eta) \right) \\ &\geq \mu(U) - 2\varepsilon.\end{aligned}$$

It is true for every $\varepsilon > 0$, so $\liminf_{n \rightarrow \infty} \mu_n(U) \geq \mu(U)$, and by Portmanteau's theorem $\mu_n \rightarrow \mu$.

- (3) The positivity of Δ_T is obvious since $1 \geq \cos Tx$. Notice then that the function $x \mapsto \frac{1 - \cos Tx}{\pi T x^2}$ is in $\mathcal{L}^1(\mathbb{R}) \cap \mathcal{L}^2(\mathbb{R})$ (it is continuous and bounded by $\frac{1}{x^2}$ at infinity). Therefore, it is the inverse Fourier transform of its Fourier transform. So, it is equivalent to prove

$$\widehat{\Delta_T}(t) = \mathbf{1}_{|t| \leq T} \left(1 - \frac{|t|}{T} \right)$$

and to prove that

$$\Delta_T(x) = \frac{1}{2\pi} \int_{-T}^T \left(1 - \frac{|t|}{T} \right) e^{-itx} dt.$$

This integral is easily computed:

$$\begin{aligned} \frac{1}{2\pi} \int_{-T}^T \left(1 - \frac{|t|}{T} \right) e^{-itx} dt &= \frac{1}{2\pi} \left[\frac{e^{-itx}}{-ix} \right]_{-T}^T + \frac{1}{2\pi T} \int_{-T}^0 t e^{-itx} dt - \frac{1}{2\pi T} \int_0^T t e^{-itx} dt \\ &= \frac{\sin Tx}{\pi x} + \frac{1}{2\pi T} \left(\left[\frac{t e^{-itx}}{-ix} \right]_{-T}^0 - \left[\frac{t e^{-itx}}{-ix} \right]_0^T \right) \\ &\quad + \frac{1}{2\pi T} \left(\int_{-T}^0 \frac{e^{-itx}}{ix} dt - \int_0^T \frac{e^{-itx}}{ix} dt \right) \\ &= \frac{1}{2\pi T x^2} \left([e^{-itx}]_{-T}^0 - [e^{-itx}]_0^T \right) = \frac{1 - \cos Tx}{\pi T x^2}. \end{aligned}$$

We have then computed the Fourier transform of Δ_T , and in particular,

$$\int_{\mathbb{R}} \Delta_T(x) dx = \widehat{\Delta_T}(0) = 1.$$

Finally, set $h > 0$. The mass of $(-h, h)^c$ is smaller than

$$\frac{2}{\pi T} \int_h^\infty \frac{1 - \cos Tx}{x^2} dx \leq \frac{2}{\pi T} \int_h^\infty \frac{2}{x^2} dx = \frac{4}{\pi T h}.$$

- (4) If $\eta = 0$, then $F = G$, $D = F - G = 0$, and so is its convolution by Δ_T :

$$D_T(x) = \int_{\mathbb{R}} D(x-y) \Delta_T(y) dy = \int_{\mathbb{R}} 0 dy = 0 \quad \Rightarrow \quad \eta_T = 0.$$

Suppose $\eta > 0$ and fix x_0 such that $D(x_0) = \eta$. Since F is increasing and G is Lipschitz with constant m , for every $|y| \leq \frac{\eta}{2m}$,

$$\begin{aligned} D(x-y) &= F(x-y) - G(x-y) \\ &\geq F(x_0) - G(x_0) - m(x-y-x_0) = \eta - mh + my = \frac{\eta}{2} + my. \end{aligned}$$

Then,

$$\begin{aligned} \eta_T &\geq D_T(x) = \int_{\mathbb{R}} D(x-y) \Delta_T(y) dy \\ &\geq -\eta \int_{(-h,h)^c} \Delta_T(y) dy + \int_{(-h,h)} \left(\frac{\eta}{2} + my \right) \Delta_T(y) dy \\ &\geq -\eta \int_{(-h,h)^c} \Delta_T(y) dy + \frac{\eta}{2} \int_{(-h,h)} \Delta_T(y) dy \quad \text{since } y \Delta_T(y) \text{ is odd;} \\ &\geq \frac{\eta}{2} - \frac{3\eta}{2} \int_{(-h,h)^c} \Delta_T(y) dy \\ &\geq \frac{\eta}{2} - \frac{6\eta}{\pi T h} = \frac{\eta}{2} - \frac{12m}{\pi T}. \end{aligned}$$

The inequality is also trivially true when $\eta = 0$.

(5) Recall that the Fourier transform of a convolution is the product of Fourier transforms:

$$\begin{aligned} (\widehat{a * b})(t) &= \int_{\mathbb{R}} \left(\int_{\mathbb{R}} a(x-y) b(y) dy \right) e^{itx} dx \\ &= \iint_{\mathbb{R} \times \mathbb{R}} (a(x-y) e^{it(x-y)} dx) (b(y) e^{ity} dy) = \widehat{a}(t) \widehat{b}(t) \end{aligned}$$

We use this result with $a = D$ and $b = \Delta_T$, after a Fourier inversion of D_T (we have the right to do so, since a and b are both in $\mathcal{L}^1(\mathbb{R}) \cap \mathcal{L}^2(\mathbb{R})$):

$$\begin{aligned} D_T(x) &= \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{D_T}(t) e^{-itx} dt \\ &= \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{D}(t) \widehat{\Delta_T}(t) e^{-itx} dt \\ &= \frac{1}{2\pi} \int_{\mathbb{R}} (\widehat{F}(t) - \widehat{G}(t)) \widehat{\Delta_T}(t) e^{-itx} dt. \end{aligned}$$

Notice then that $\widehat{a'}(t) = -it \widehat{a}(t)$ by integration by parts:

$$\widehat{a'}(t) = \int_{\mathbb{R}} a'(x) e^{itx} dx = -it \int_{\mathbb{R}} a(x) e^{itx} dx = -it \widehat{a}(t).$$

Therefore,

$$\begin{aligned} \widehat{F}(t) - \widehat{G}(t) &= \frac{\widehat{\mu}(t) - \widehat{g}(t)}{-it} = \frac{\phi_{\mu}(t) - \phi_G(t)}{-it}; \\ D_T(x) &= \frac{1}{2\pi} \int_{\mathbb{R}} \widehat{\Delta_T}(t) \left(\frac{\phi_{\mu}(t) - \phi_G(t)}{-it} \right) e^{-itx} dt \\ &\leq \frac{1}{2\pi} \int_{\mathbb{R}} |\widehat{\Delta_T}(t)| \left| \frac{\phi_{\mu}(t) - \phi_G(t)}{t} \right| dt \\ &\leq \frac{1}{2\pi} \int_{-T}^T \left| \frac{\phi_{\mu}(t) - \phi_G(t)}{t} \right| dt. \end{aligned}$$

This is true for every x , so η_T has the same bound, and then

$$\eta \leq 2\eta_T + \frac{24m}{\pi T} \leq \frac{1}{\pi} \int_{-T}^T \left| \frac{\phi_{\mu}(t) - \phi_G(t)}{t} \right| dt + \frac{24m}{\pi T}.$$

(6) We calculate

$$\begin{aligned} \phi_{\mu_n}(t) &= \mathbb{E}[e^{itX_n}] = \mathbb{E}[e^{itY_n/\sqrt{\lambda_n}}] = e^{-\frac{t^2}{2}} \psi_n\left(\frac{t}{\sqrt{\lambda_n}}\right) \\ &= e^{-\frac{t^2}{2}} \left(\psi\left(\frac{t}{\sqrt{\lambda_n}}\right) + O\left(\frac{t}{\sqrt{\lambda_n}}\right) o\left(\frac{1}{\sqrt{\lambda_n}}\right) \right) \end{aligned}$$

since one has uniform convergence $\psi_n \rightarrow \psi$ at speed $o(\sqrt{\lambda_n}^{-1})$ if the argument $t/\sqrt{\lambda_n}$ stays bounded. The remainder is a $o\left(\frac{t}{\sqrt{\lambda_n}}\right)$; then we can expand ψ in Taylor series to get

$$\phi_{\mu_n}(t) = e^{-\frac{t^2}{2}} \left(1 + \frac{\psi'(0)t}{\sqrt{\lambda_n}} + o\left(\frac{t}{\sqrt{\lambda_n}}\right) \right).$$

Then,

$$\begin{aligned} \int_{\mathbb{R}} g_n(x) e^{itx} dx &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \left(1 - \frac{i\psi'(0)x}{\sqrt{\lambda_n}} \right) e^{itx - \frac{x^2}{2}} dx \\ &= \left(1 - \frac{\psi'(0)}{\sqrt{\lambda_n}} \frac{\partial}{\partial t} \right) \left(\frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{itx - \frac{x^2}{2}} dx \right) \\ &= \left(1 - \frac{\psi'(0)}{\sqrt{\lambda_n}} \frac{\partial}{\partial t} \right) e^{-\frac{t^2}{2}} \\ &= e^{-\frac{t^2}{2}} \left(1 + \frac{\psi'(0)t}{\sqrt{\lambda_n}} \right). \end{aligned}$$

This function is bounded, so the integral $G_n(x)$ is indeed m -Lipschitz for a certain constant m . Moreover, one has indeed $G_n(+\infty) = 1$ and $G_n(-\infty) = 0$, so Berry's lemma will apply.

(7) Berry's lemma ensures that

$$\begin{aligned} \eta_n &= \sup_{x \in \mathbb{R}} |F_{\mu_n}(x) - G_n(x)| \\ &\leq \frac{1}{\pi} \int_{-K\sqrt{\lambda_n}}^{K\sqrt{\lambda_n}} \left| \frac{\phi_{\mu_n}(t) - \widehat{g}_n(t)}{t} \right| dt + \frac{24m}{\pi K \sqrt{\lambda_n}} \\ &\leq \frac{1}{\pi \sqrt{\lambda_n}} \left(\int_{(-K\sqrt{\lambda_n}, K\sqrt{\lambda_n}) \setminus (-\varepsilon\sqrt{\lambda_n}, \varepsilon\sqrt{\lambda_n})} + \int_{(-\varepsilon\sqrt{\lambda_n}, \varepsilon\sqrt{\lambda_n})} \right) \left| e^{-\frac{t^2}{2}} v\left(\frac{t}{\sqrt{\lambda_n}}\right) \right| dt + \frac{24m}{\pi K \sqrt{\lambda_n}} \end{aligned}$$

where $v(\cdot)$ is a function going to zero at zero. Fix $\theta > 0$, and ε such that $|v(\cdot)|$ is smaller than θ on $(-\varepsilon, \varepsilon)$. The integral on $(-\varepsilon\sqrt{\lambda_n}, \varepsilon\sqrt{\lambda_n})$ is then smaller than

$$\frac{\theta}{\pi \sqrt{\lambda_n}} \int_{\mathbb{R}} e^{-\frac{t^2}{2}} dt \leq \frac{\theta}{\sqrt{\lambda_n}}.$$

Fix then K such that $\frac{24m}{\pi K} \leq \theta$. The third summand in the upper bound is then also smaller than $\frac{\theta}{\sqrt{\lambda_n}}$, and on the other hand, $v(\cdot)$ is bounded by some constant $C(K)$ on $(-K, K)$. This leads to

$$\begin{aligned} \eta_n &\leq \frac{2\theta}{\sqrt{\lambda_n}} + \frac{1}{\pi \sqrt{\lambda_n}} \int_{(-K\sqrt{\lambda_n}, K\sqrt{\lambda_n}) \setminus (-\varepsilon\sqrt{\lambda_n}, \varepsilon\sqrt{\lambda_n})} C(K) e^{-\frac{t^2}{2}} dt \\ &\leq \frac{2\theta}{\sqrt{\lambda_n}} + \frac{C(K)}{\pi \sqrt{\lambda_n}} \int_{\mathbb{R} \setminus (-\varepsilon\sqrt{\lambda_n}, \varepsilon\sqrt{\lambda_n})} e^{-\frac{t^2}{2}} dt \\ &\leq \frac{2\theta}{\sqrt{\lambda_n}} + \frac{C(K) e^{-\frac{\varepsilon^2 \lambda_n}{2}}}{\pi \varepsilon \lambda_n} \leq \frac{3\theta}{\sqrt{\lambda_n}} \quad \text{for } n \text{ big enough.} \end{aligned}$$

Since this is true for every θ , the estimate $o(1/\sqrt{\lambda_n})$ is shown.

(8) The previous Question proves that up to a uniform $o(1/\sqrt{\lambda_n})$,

$$F_{\mu_n}(x) \simeq F_{\mathcal{N}(0,1)}(x) + \int_{-\infty}^x \frac{-i\psi'(0)y}{\sqrt{\lambda_n}} \frac{e^{-\frac{y^2}{2}}}{\sqrt{2\pi}} dy = F_{\mathcal{N}(0,1)}(x) + \frac{i\psi'(0)}{\sqrt{2\pi\lambda_n}} e^{-\frac{x^2}{2}},$$

so

$$d(\mu_{X_n}, \mathcal{N}(0,1)) \simeq \frac{|\operatorname{Im}(\psi'(0))|}{\sqrt{2\pi\lambda_n}} \sup_{x \in \mathbb{R}} \left(e^{-\frac{x^2}{2}} \right) \simeq \frac{|\operatorname{Im}(\psi'(0))|}{\sqrt{2\pi\lambda_n}}.$$