Observation of entanglement between collective excitation in a quantum fluid: when Faraday waves grow from vacuum fluctuations

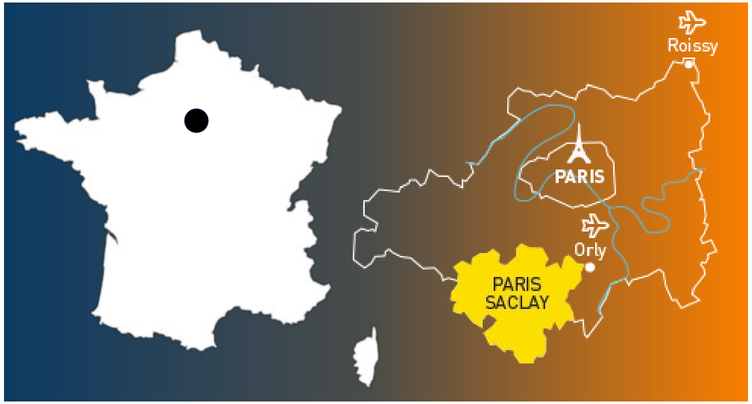
Victor Gondret,

20th of January, 2026

Pontificia Universidad católica de Chile, Santiago

Slides available at
www.normalesup.org/~gondret/talk.pdf

Paris-Saclay University



Paris-Saclay University

Do not be fooled, it is more Saclay, than Paris!

Institut d'Optique Graduate School

Is part of the university

LABORATOIRE CHARLES FABRY

The unique lab of the school

43 researchers/profs
75 PhDs/post-docs

Research @LCF

Adaptative optics

Nanophotonics

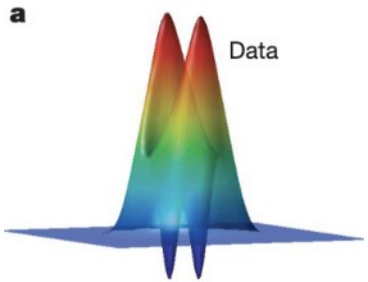
Biophotonics

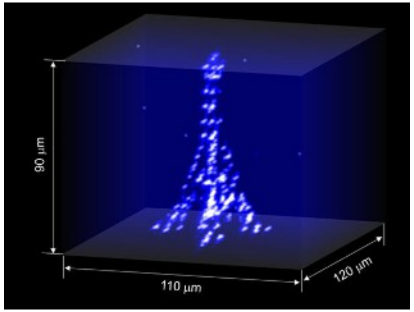
Laser

Quantum optics

Photons

Atoms



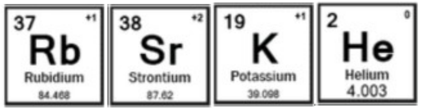


Quantum gases

Quantum gases at LCF

8 researchers/professors

6 experiments (from 1D to 3D)



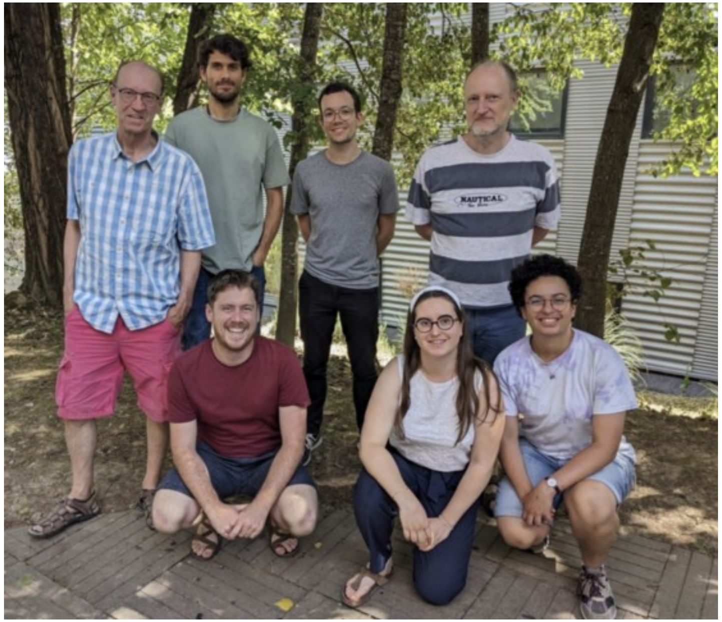
Optical lattice

3D Anderson localization

1D gases

Ultracold Fermions

Quantum microscope

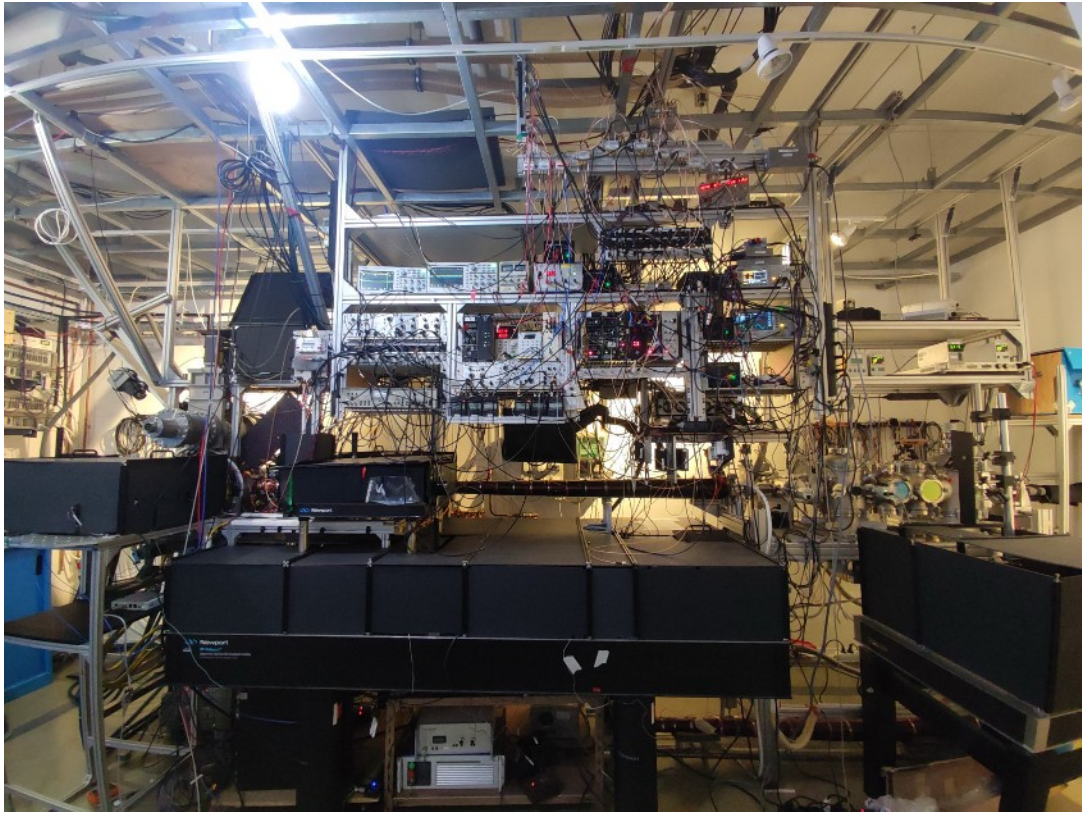


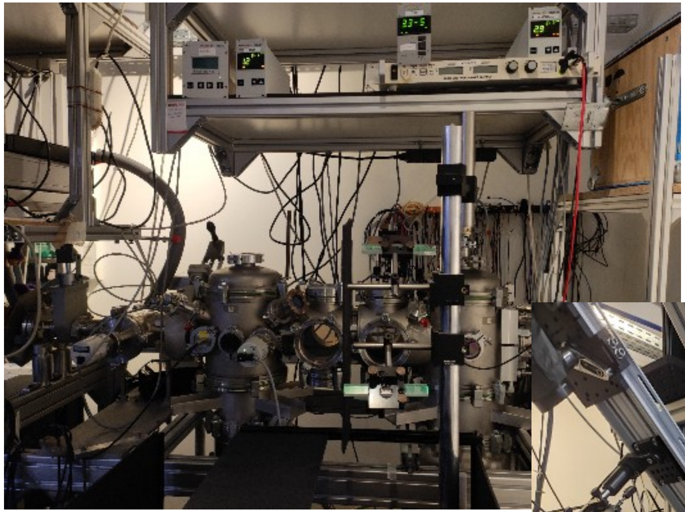
Exp: Chris Westbrook, Rui Dias, Charlie Leprince, Denis Boiron, Victor Gondret, Clothilde Lamirault, Léa Camier

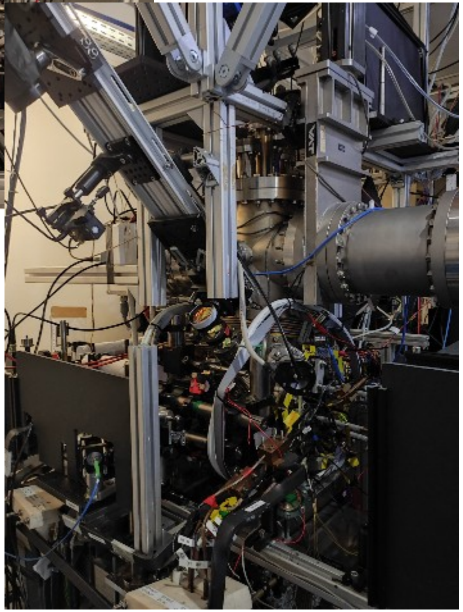
 

Th: Amaury Micheli & Scott Robertson

Experimental setup







Experiment started in 1994: oldest BEC experiment in France (metastable Helium).

OUTLINE

I. Introduction

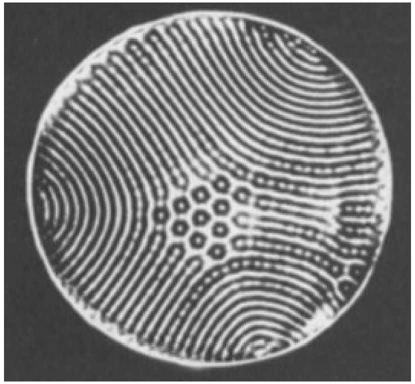
II. Model and setup

III. Growth and decay of quasiparticles

IV. Assessing entanglement of two-mode Gaussian states with many-body correlation function

V. Observation of entanglement

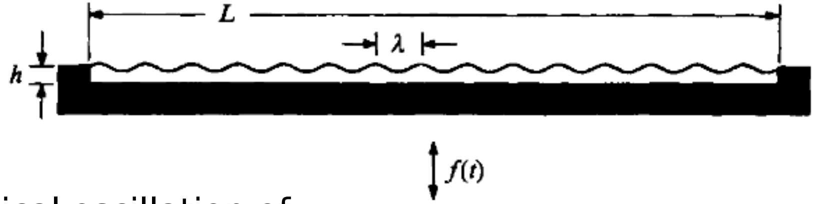
I. Introduction



Faraday waves

Broughton Suspension Bridge collapsed in 1831

Guan *et al.* PR Fluids (2023), Edwards & Fauve J. Fluid Mech. (1994)



Vertical oscillation of the tank at Ω

$$\omega_k = \sqrt{\tanh(hk)[gk + \gamma k]} = \Omega/2$$

Modulation of the effective gravity

- g gravity
- γ surface tension

Parametric oscillation ≠ forced oscillation

Ω/2	Ω
Variation of an internal parameter	External force

Parametric excitation & forced oscillation

Parametric oscillation ≠ forced oscillation

Parametric oscillation	Forced oscillation
$\Omega/2$	Ω
Variation of an internal parameter	External force

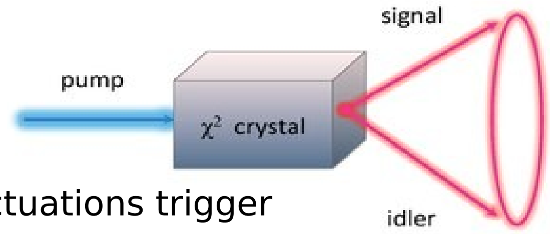
Parametric excitation & forced oscillation

Parametric oscillation ≠ forced oscillation

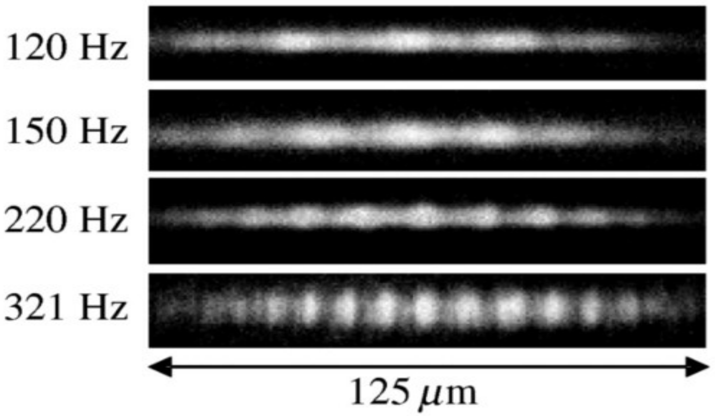
$\Omega/2$	Ω
Variation of an internal parameter	External force
Growth triggered by fluctuations	Growth initialized by the force

- Experimental imperfection
- Thermal fluctuations,
- Quantum fluctuation

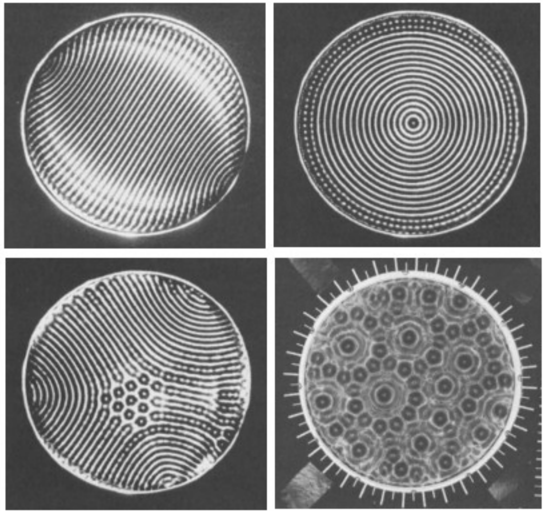
Photons



Vacuum fluctuations trigger amplification which leads to entanglement.



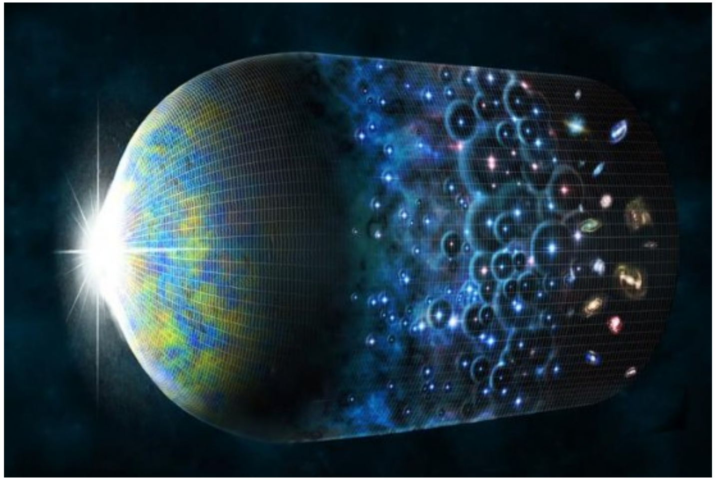
Bose-Einstein condensates



Fluid

Edwards & Fauve J. Fluid Mech. **278**, 123 (1994).

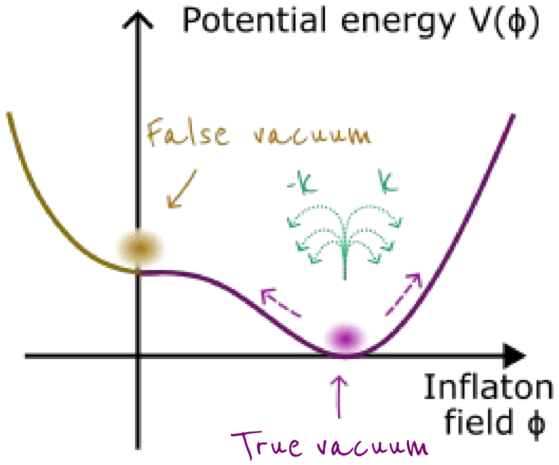
Early universe



Preheating in the early universe

The **inflaton** goes from its initial false vacuum state. Its almost constant potential energy **drives the inflation**.
A. Linde, Phys. Lett. **129B**, 177 (1983).

It starts to oscillate around its minimum and, coupled to matter fields, it creates particles through broad **parametric resonance**.
L. Kofman, A. Linde & A. Starobinsky, Phys. Rev. D 56, (1997).

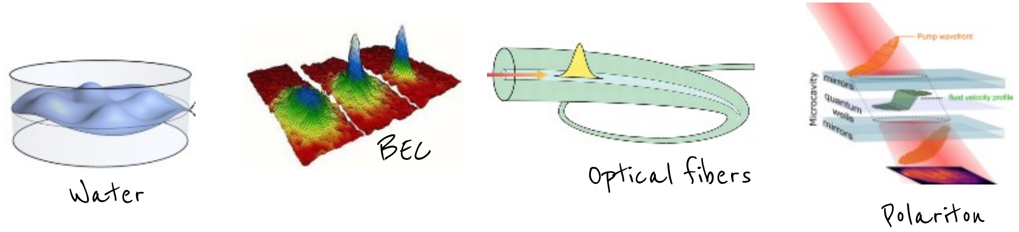


Particles are created in **pairs** with **opposite momenta from vacuum** in a highly entangled two modes squeezed state. Interactions lead to decoherence and thermalization.
D. Campo & R. Parentani, Phys. Rev. D **74**, 025001 (2006).

BUT NOT OBSERVABLE

Analog gravity

Use the tools of quantum field theory formalism to describe a condensed matter system,

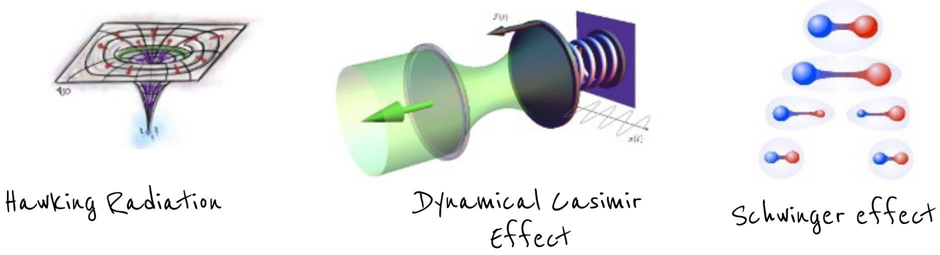


Analog gravity

In the presence of a strong coherent background, the excitations of a fluid, or *quasi*particles, can be treated using the same formalism as particles in a curved spacetime.

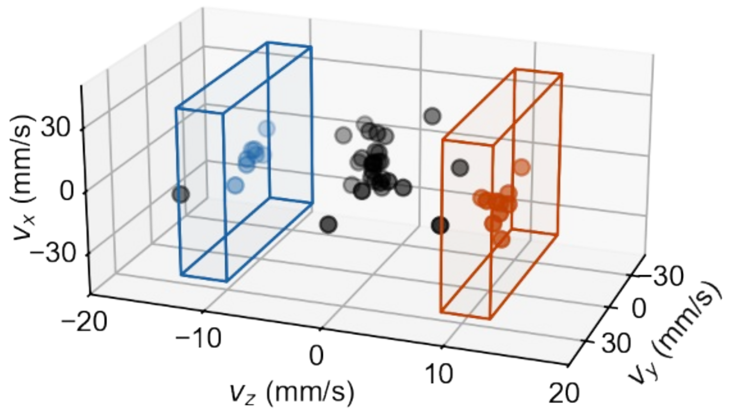
Unruh, *Experimental black-hole evaporation?* Phys. Rev. Lett. **46**, 1351 (1981)

Shape the fluid to mimic famous effect of QFT.



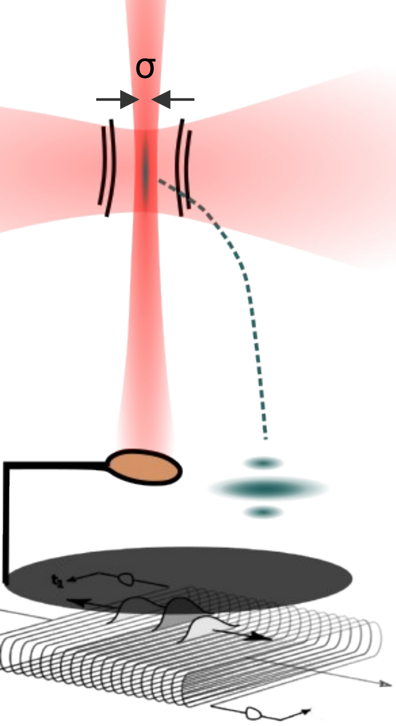
Goal: witness vacuum fluctuation amplification

II. Model and setup



Collective excitations in a BEC

- g effective interaction strength
- n density
- m atomic mass
- $\hbar$ reduced Planck cte



- BEC of He* in 10s with few thousands atoms at 25(5) nK.
- 1 kHz & 50 Hz: effective 1D dynamics

Bose gas with contact interaction

$$\hat{H} = \sum_k \frac{\hbar^2 k^2}{2m} \hat{a}_k^\dagger \hat{a}_k + \frac{g}{2V} \sum_{k_1,k_2,q} \hat{a}_{k_1+q}^\dagger \hat{a}_{k_2-q}^\dagger \hat{a}_{k_2} \hat{a}_{k_1}$$

with $\hat{a}_k$ the atomic annihilation operator.

How to change gn ?

- g with a Feshbach resonance (Chicago, Rice, Heidelberg)
- n transverse with trap modulation (Mexico, NIST, Palaiseau, Trento, Utrecht)

Bogoliubov description:

We treat the BEC as a coherent state and quantized other modes k. Introduce the quasiparticle modes b_k which diagonalize the Hamiltonian.

$$\hat{a}_k = u_k \hat{b}_k + v_k \hat{b}_{-k}^\dagger$$

with u_k and v_k the Bogoliubov coefficients and

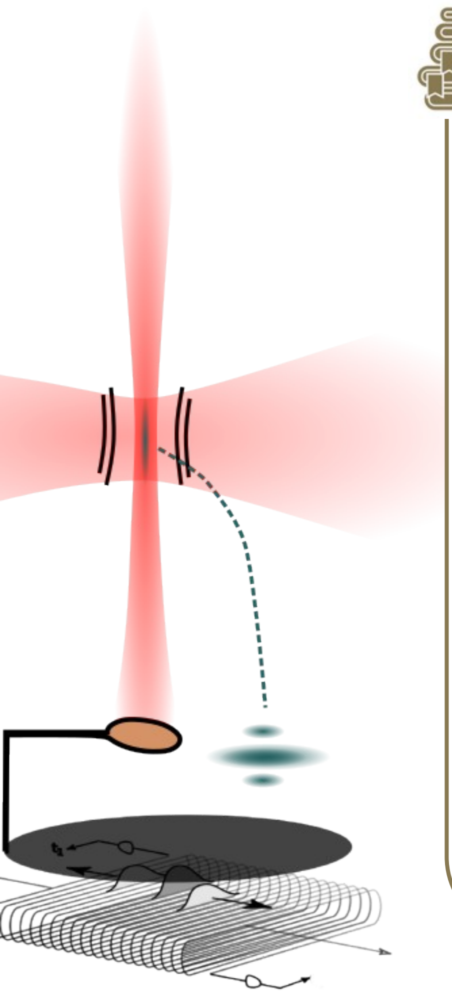
$$\omega_k = \sqrt{\frac{gn}{m} k^2 + \left(\frac{\hbar k^2}{2m}\right)^2}$$

Quasiparticle evolution

$$\partial_t \hat{b}_k = -i\omega_k \hat{b}_k + \frac{\dot{\omega}_k}{2\omega_k} \hat{b}_{-k}^\dagger$$

Busch *et al.* Phys. Rev. A **89** (2014)

Theoretical details



Theoretical cheatsheet

If non-zero temperature, both thermal and vacuum fluctuations trigger the exponential growth.

Amplification of vacuum fluctuation is witnessed by two-mode entanglement.

A large temperature prevent the appearance of entanglement.

Beyond Bogoliubov numerics: quasiparticle interactions further destroy entanglement.

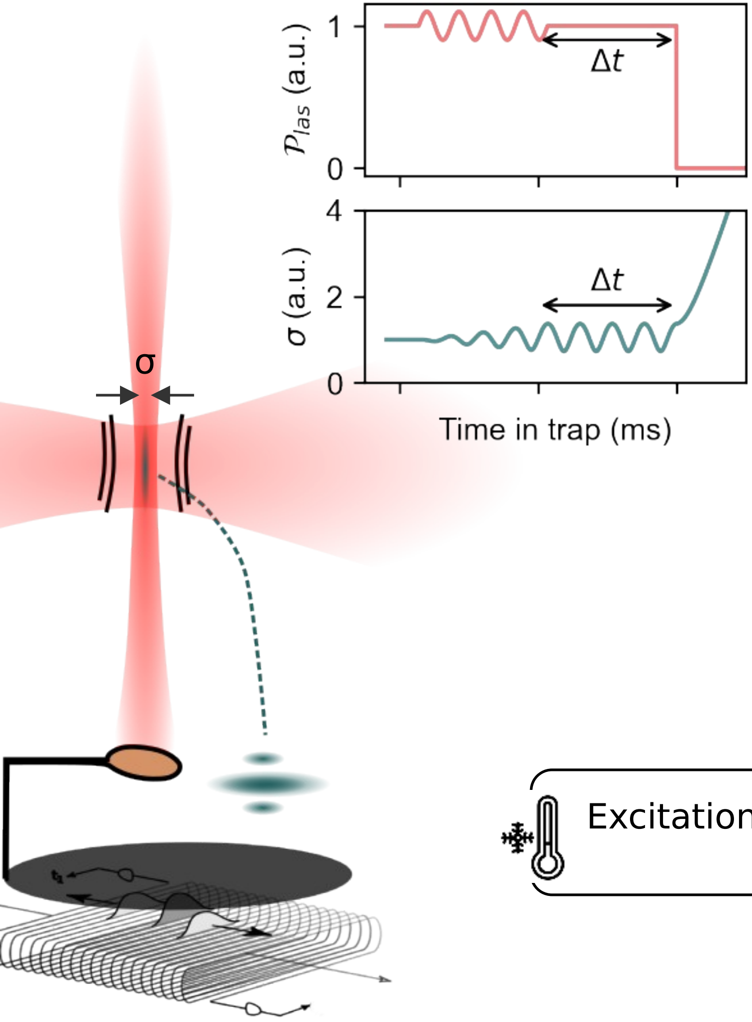
Busch *et al.* Phys. Rev. A **89** (2014),
Robertson *et al.*, Phys. Rev. D **95**, 065020 (2017),
Robertson *et al.*, Phys. Rev. D **98**, 056003 (2018).
Micheli & Robertson.Comptes Rendus. Phys., (2024),

Two-mode squeezing model.

Damped two-mode squeezing model.

Not analytics

Faraday waves with quantum fluids



Protocol

(1) Excite the transverse breathing mode of the BEC at Ω for 4 periods,

(2) Let it breath for Δt: longitudinal collective excitations with $\omega_k = \Omega/2$ are parametrically excited

$$\omega_k = \sqrt{\frac{gn}{m}k^2 + \left(\frac{\hbar k^2}{2m}\right)^2}$$

Modulation of interactions at Ω

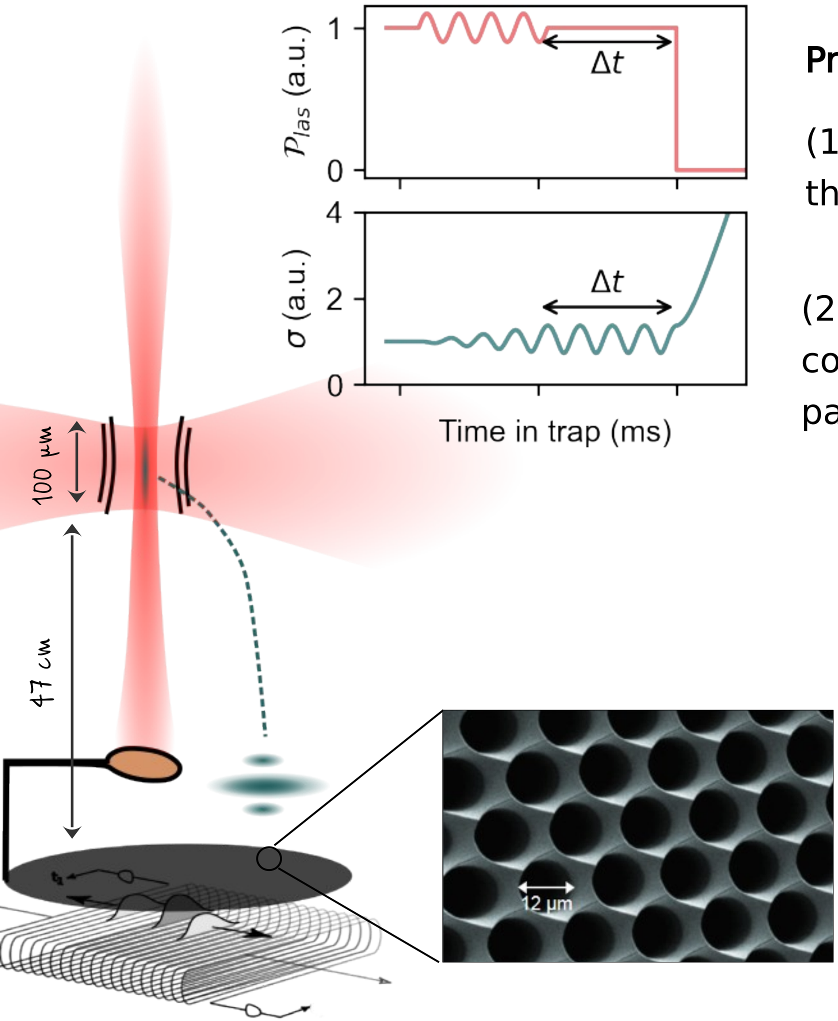
The pattern appears as Δt increases!

Excitation procedure does not heat the cloud.

We can hope to get an entangled state!

Chevy et al., Phys. Rev. Lett. **88**, 250402 (2002)
Jackson and Zaremba Phys. Rev. Lett. **89**, 150402 (2002)

Faraday waves with quantum fluids



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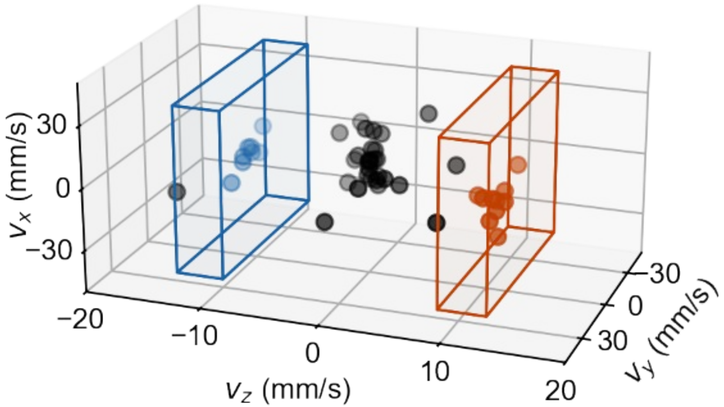
$$\omega_k = \sqrt{\frac{gn}{m}k^2 + \left(\frac{\hbar k^2}{2m}\right)^2}$$

(3) Switch off the trap: cloud expansion

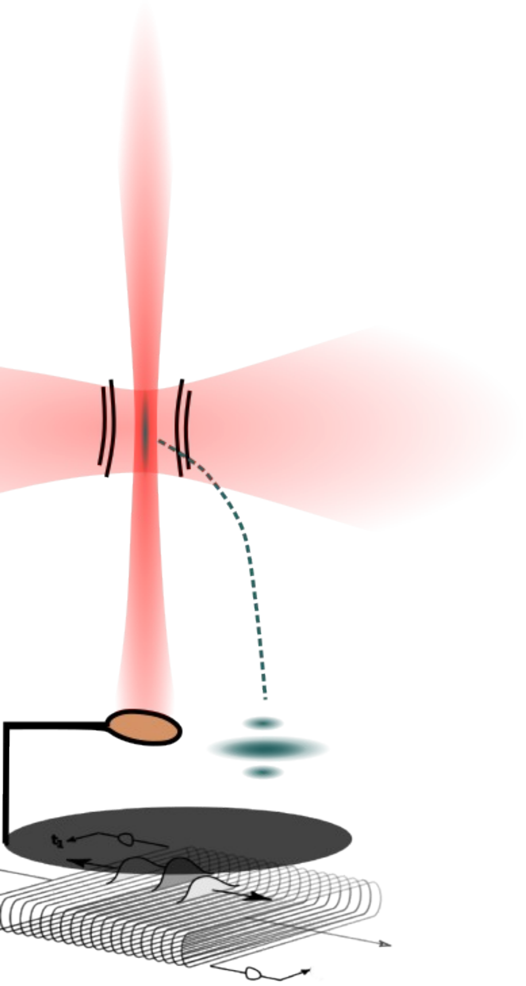
(4) Single particle detection after time of flight

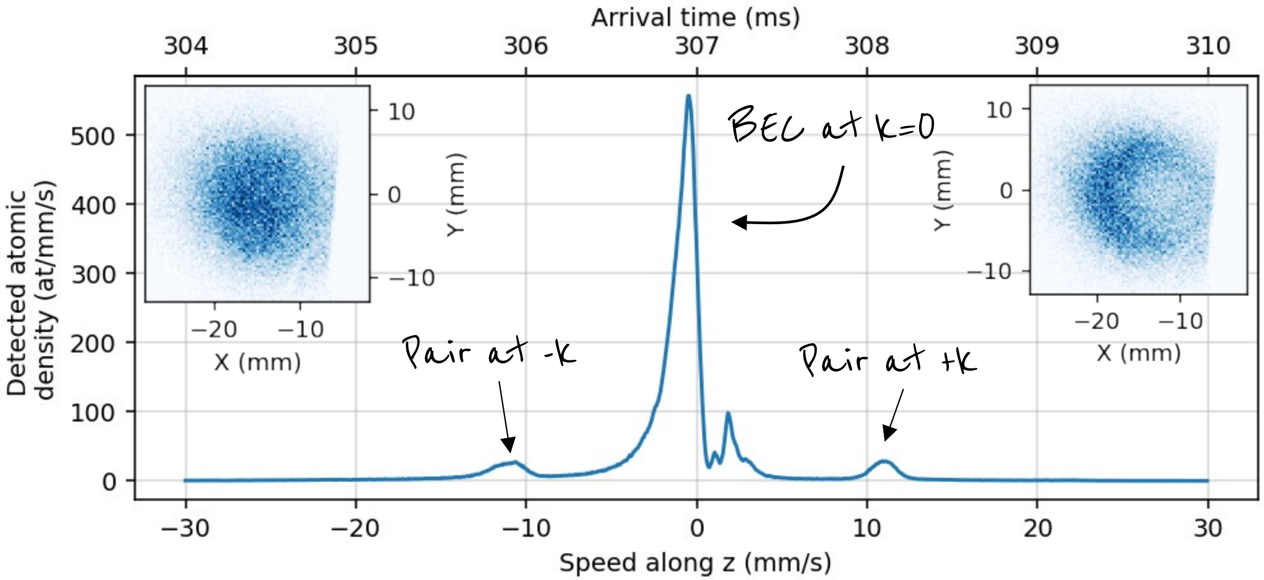
$(t,x,y) \leftrightarrow (v_Z, v_X, v_Y)$



Single shot "image", each dot is an atom. We count atoms in voxels which define the modes → measure the full particle number probability distribution

Saturation of the detector



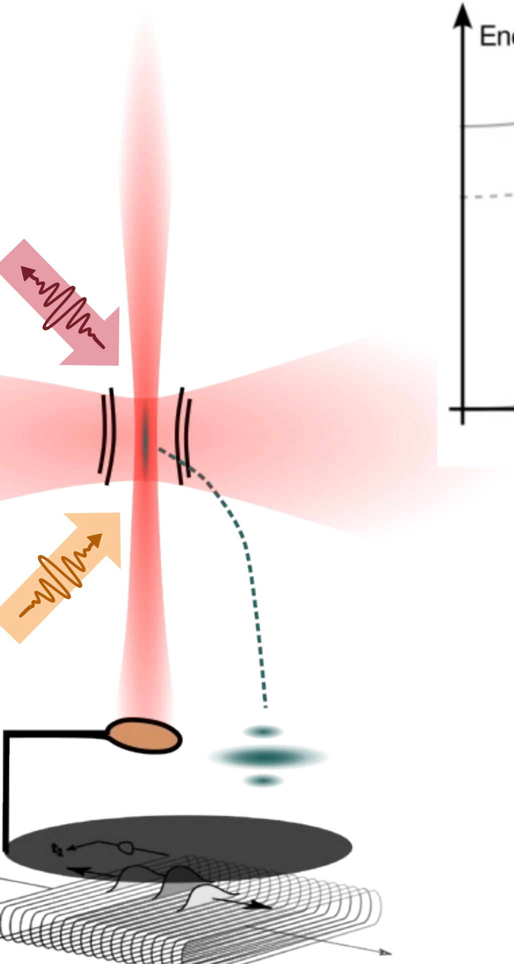


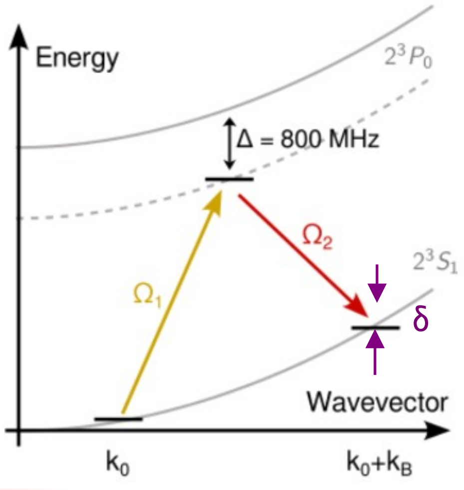
The BEC saturation affects the 2nd pair detectivity….

Use a velocity selective two-photon process to deflect only the BEC.

LABORATOIRE CHARLES FABRY

How to kick off atoms? With light!





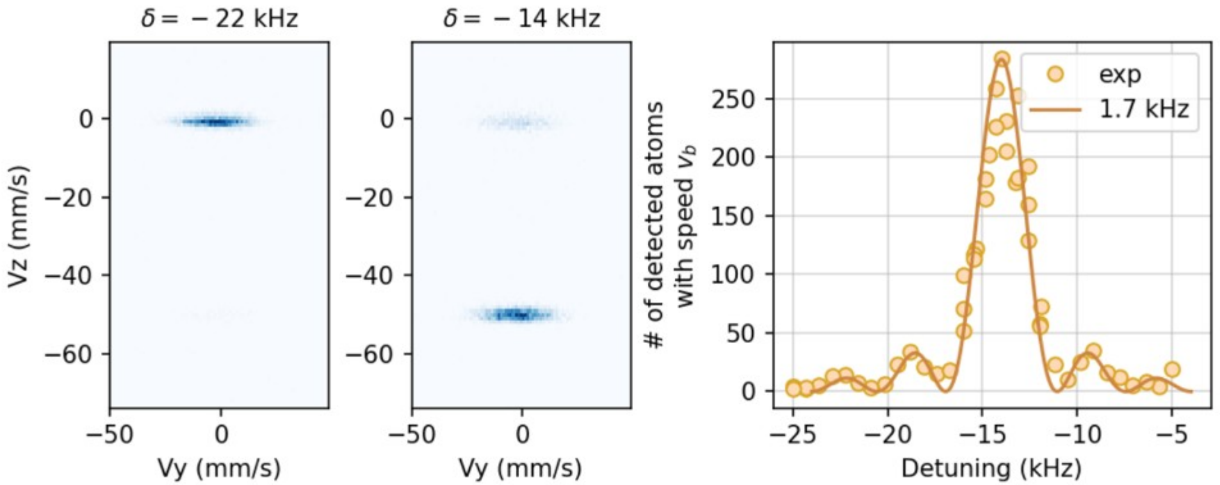
Use a velocity selective two-photon process to deflect only the BEC.

The two-photons Rabi frequency:

$$\Omega_R = \frac{\Omega_1 \Omega_2^*}{2\Delta}$$

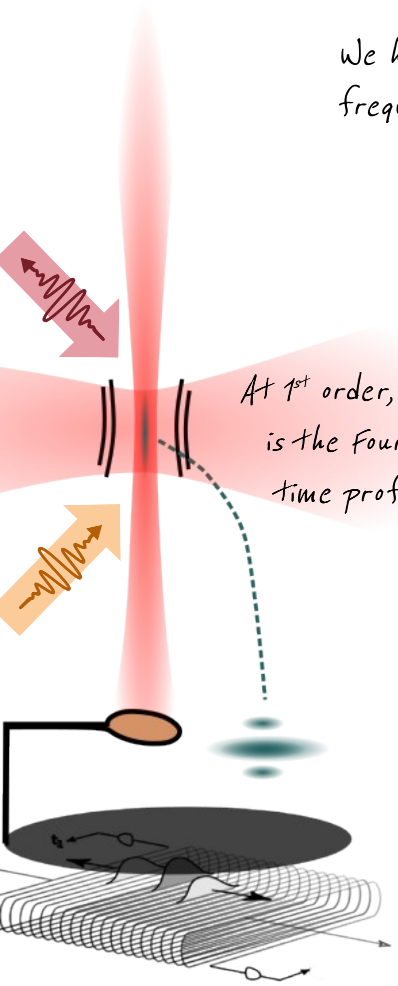
→ By changing the detuning δ between the two lasers, different velocity speeds can be addressed: δ ↔ v_{res}.

The velocity kick is fixed by the angle between the two lasers



Leprince *et al* Phys. Rev. A **111**, 063304 (2025)

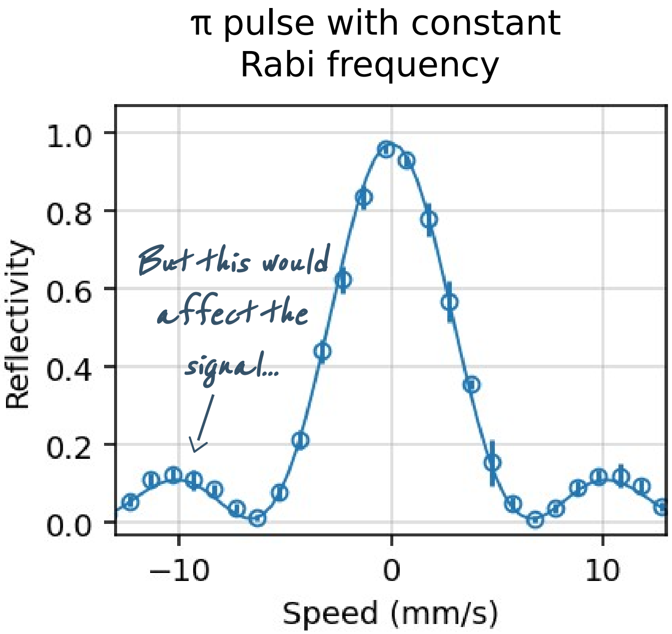
Bragg deflection of the BEC



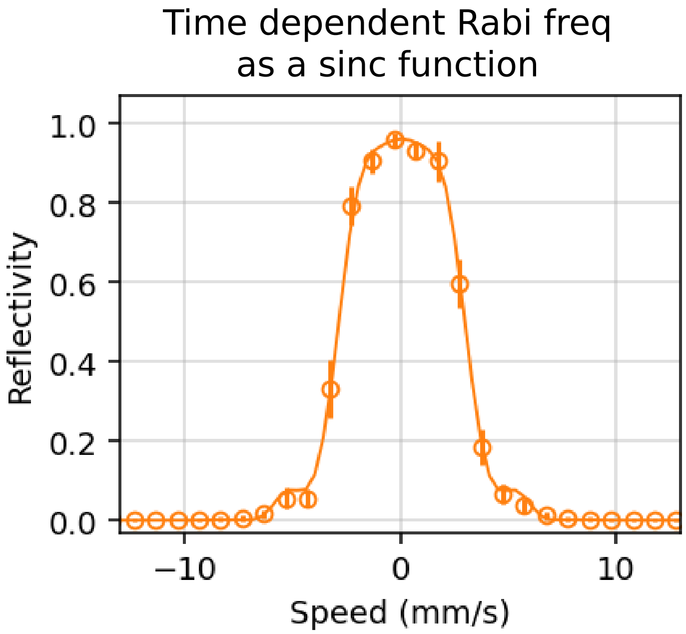
We have an equivalence
frequency ⇔ momentum

At 1st order, the reflectivity profile is the Fourier transform of the time profile of the light pulse

Use a velocity selective two-photon process to deflect only the BEC.



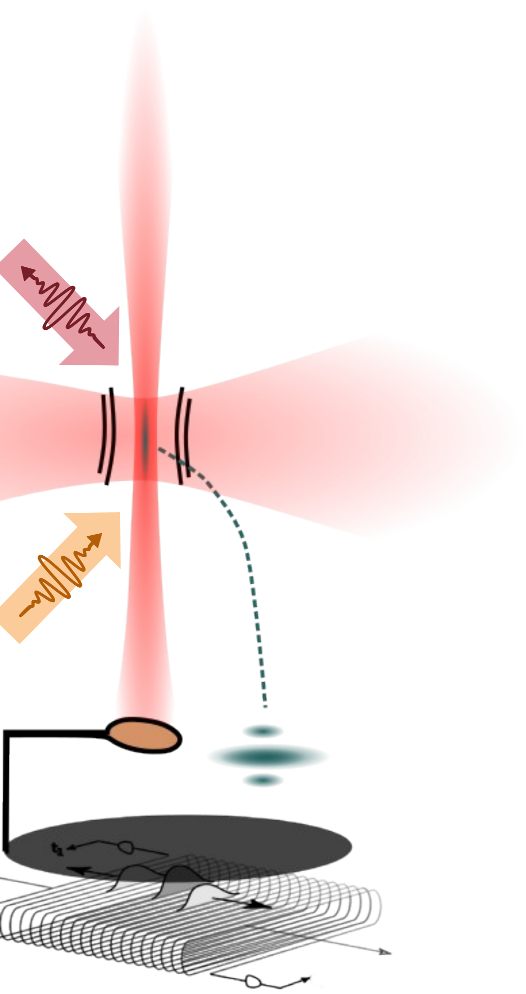
(it looks like a |sinc| function)



(it looks more like a square)

Leprince *et al* Phys. Rev. A **111**, 063304 (2025)

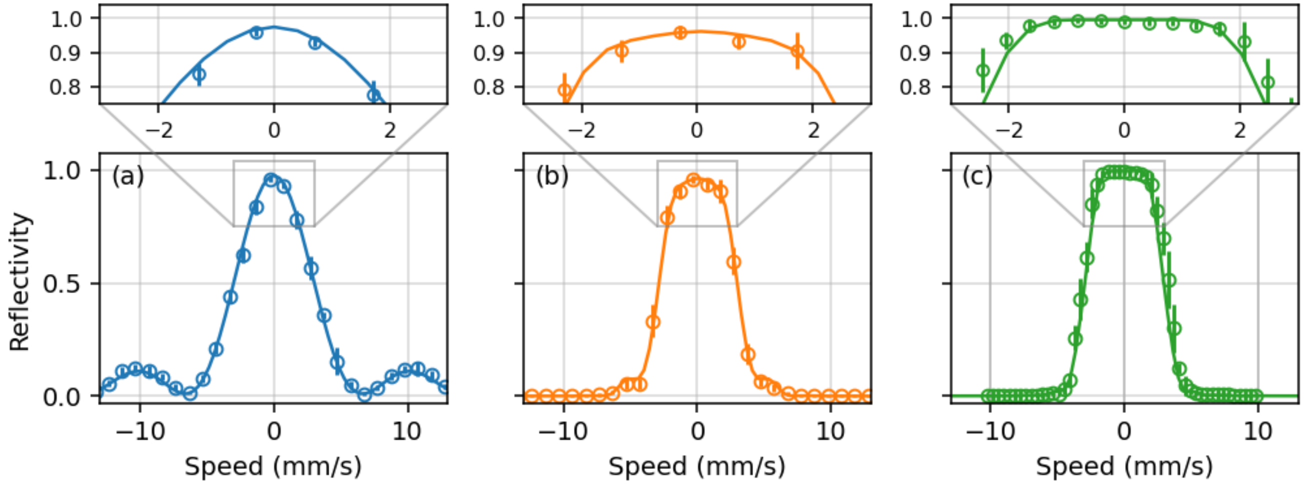
Summary



Constant pulse

Sinc pulse

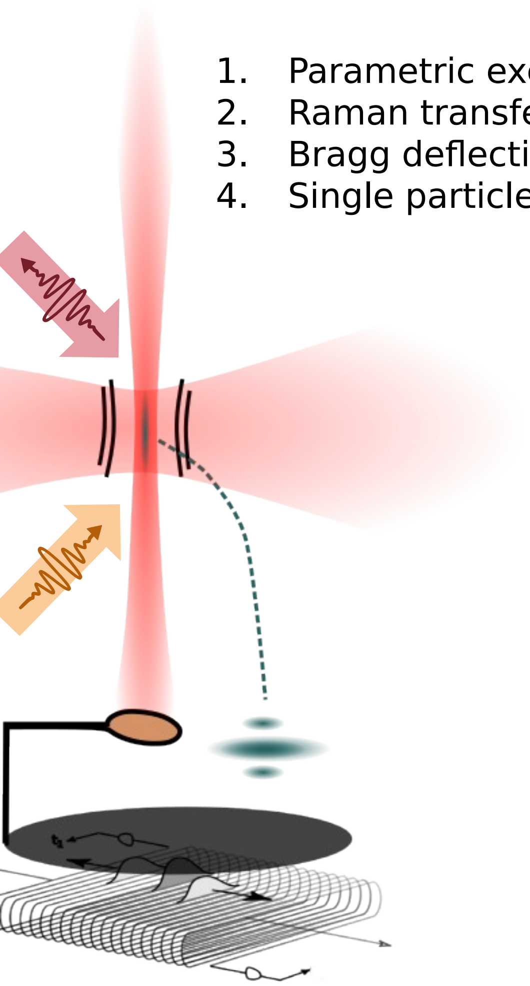
Optimized pulse
("Reburp")

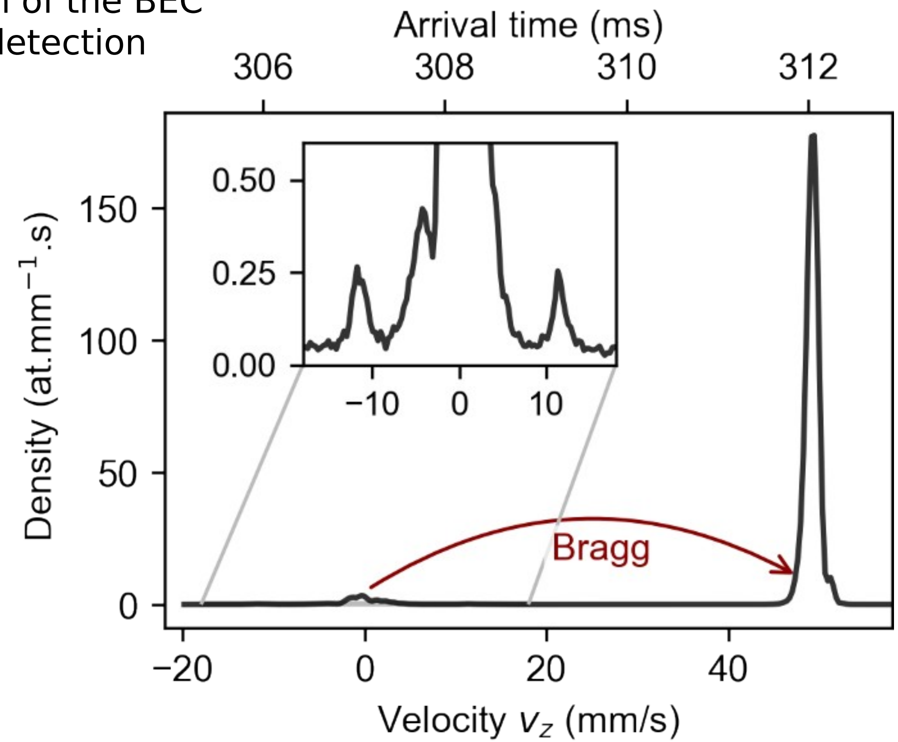


Leprince *et al* Phys. Rev. A **111**, 063304 (2025)

So does it work?

1. Parametric excitation
2. Raman transfer (+kick)
3. Bragg deflection of the BEC
4. Single particle detection



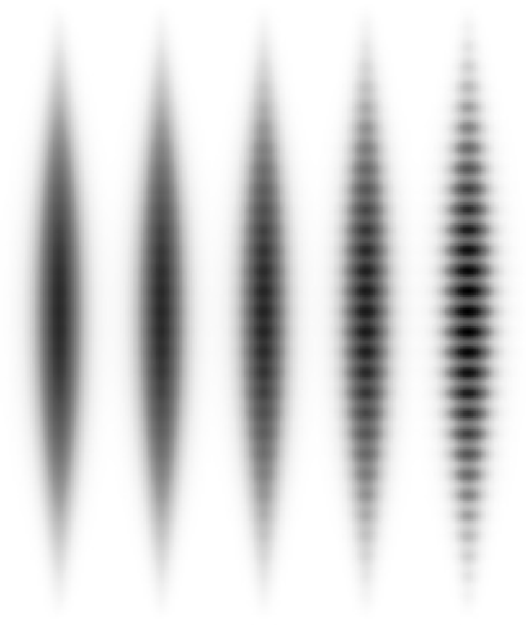


In the following, we use a pulse-shaped Bragg deflector

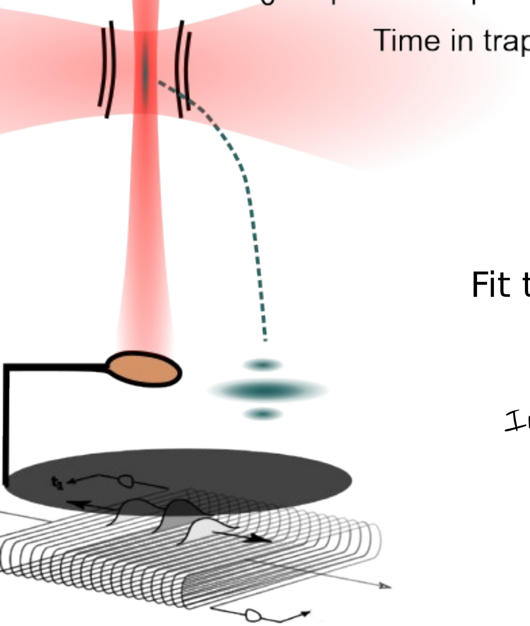
Leprince *et al* Phys. Rev. A **111**, 063304 (2025)

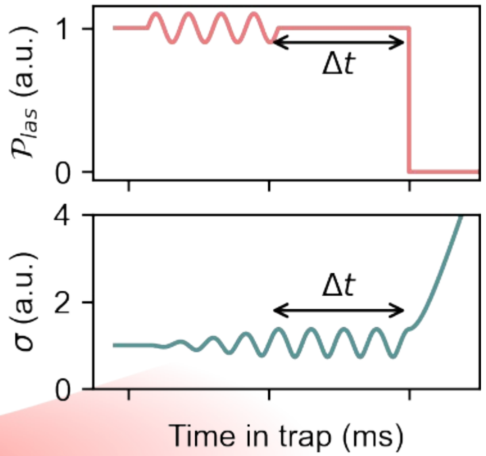
III. Growth and decay of quasiparticles

Gondret *et al.*, Parametric pair production of collective excitations in a Bose–Einstein condensate, Comptes Rendus. Physique **25**, 1 (2025).



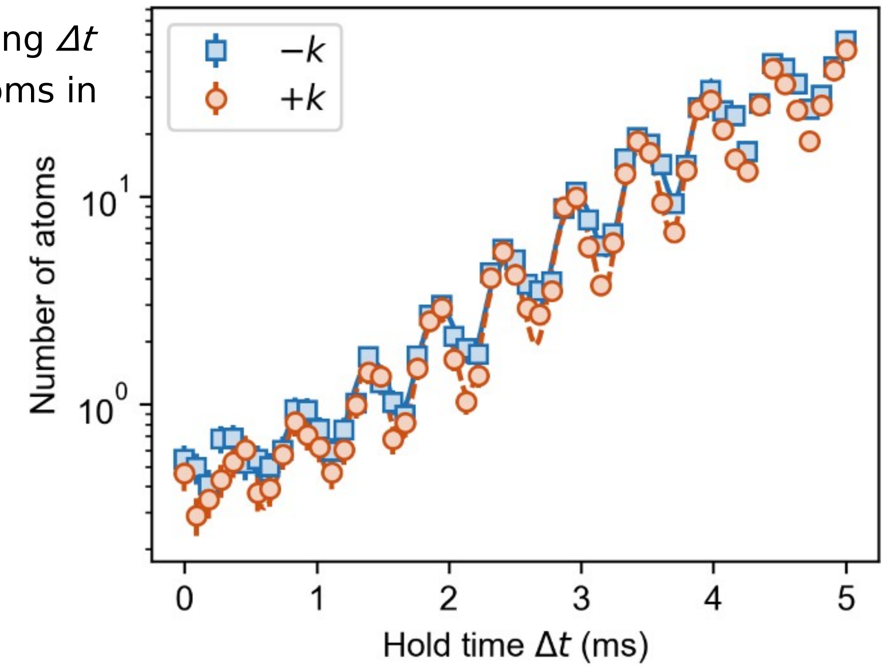
Growth of the (quasi)particle number





Repeat the experiment varying Δt and count the number of atoms in each mode

- 👍 Apparent pairwise production
- 👍 Exponential growth[1] as expected in parametric amplification
- 🤔 A large oscillation in the growth



Fit the data:

Exponential

$$n_k(\Delta t) = n_0 + \Delta n \, e^{G_k \Delta t} \times [1 + A_k \cos(2\omega_k \Delta t + \varphi_k)]\,.$$

Initial thermal population

Fluctuations that trigger the growth $\propto (n_k + n_{-k} + 1)$

Oscillation part

[1] Gondret *et al.*, Comptes Rendus. Physique **25**, 1 (2025).

Measuring the growth

$$\omega_k = \sqrt{\frac{gn}{m}k^2 + \left(\frac{\hbar k^2}{2m}\right)^2}$$

Theory: assuming a sine modulation of gn with amplitude a, the growth is analytical:

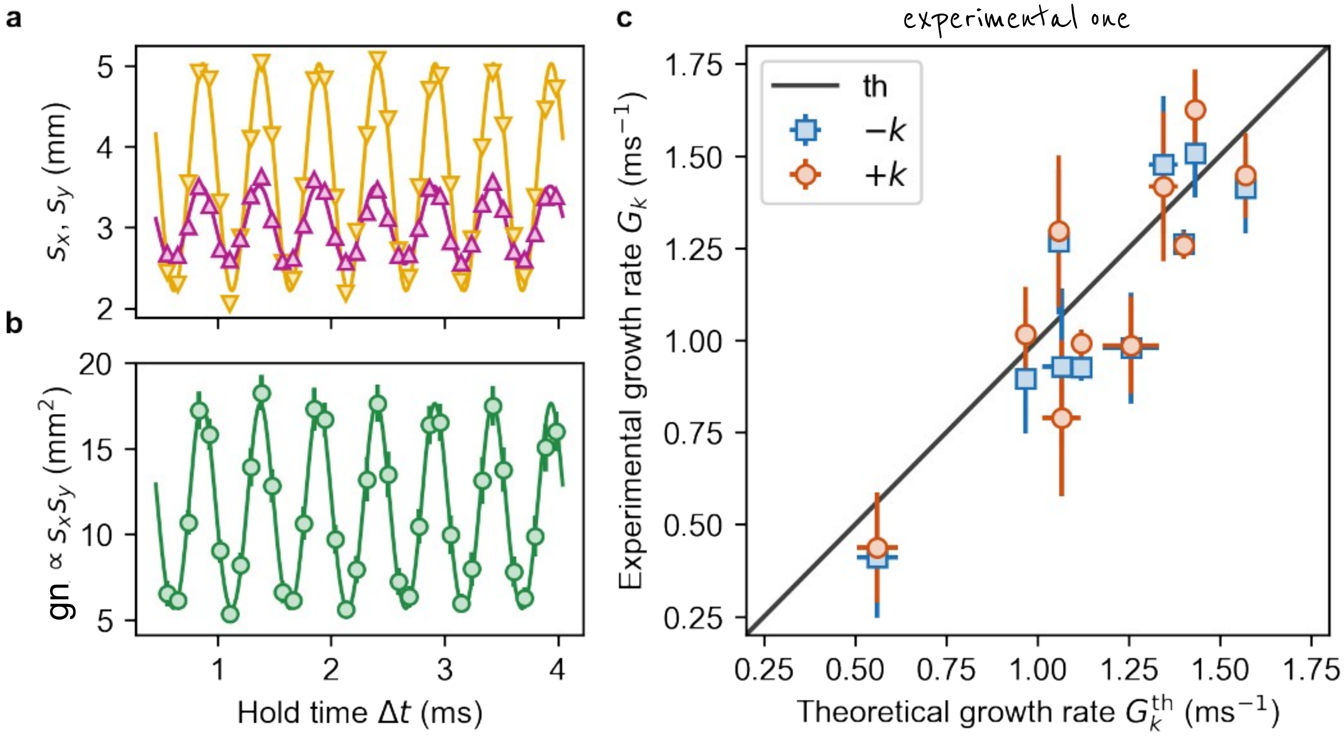
$$G_k^{\mathrm{th}} = \frac{a}{2}\frac{\omega_k}{1+k^2\xi^2/4}$$

BEC healing length

→ This model does not account for damping: the discrepancy between theory and experiment gives the value of the decay rate in the experiment.

We have access to the density modulation of the BEC

Comparison of the undamped theoretical growth rate to the experimental one

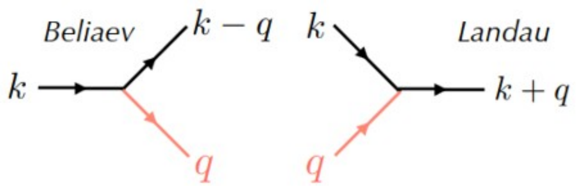


[1] Gondret *et al.*, Comptes Rendus. Physique **25**, 1 (2025).

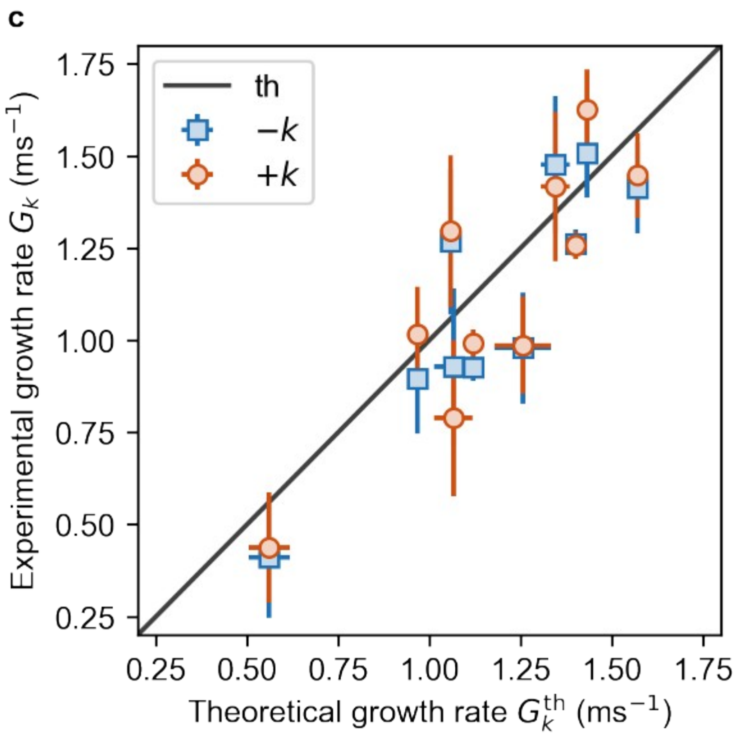
Why measuring the decay is interesting?

In a pure 1D gas, collective excitations do not decay because the system is **integrable** [1].

Recent prediction [2] derive an analytical formula for the decay of Bogoliubov quasiparticles in elongated Bose gases.



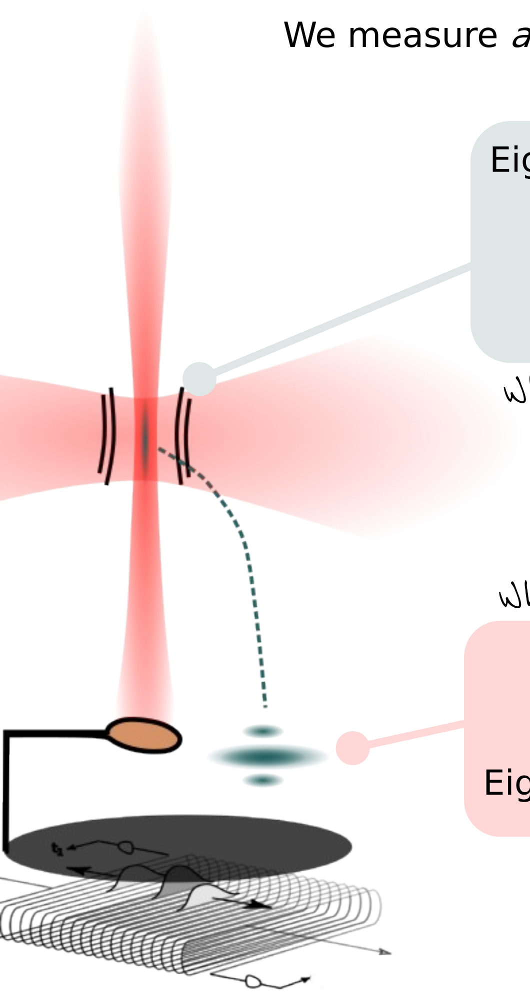
Although 3D, our cloud approach the 1D regime.
Can we check the validation of the prediction?



[1] Bouchoule *et al.,* Phys. Rev. Lett. **130**, 140401 (2023).
[2] Micheli & Robertson, Phys. Rev. B **106**, 214528 (2022).

Why this oscillation? Mapping the quasiparticles onto the particles

We measure *atoms* and not *quasiparticles*



Eigenbasis in an interacting gas

$$\omega_k = \sqrt{\frac{gn}{m}k^2 + \left(\frac{\hbar k^2}{2m}\right)^2}$$

What we produce

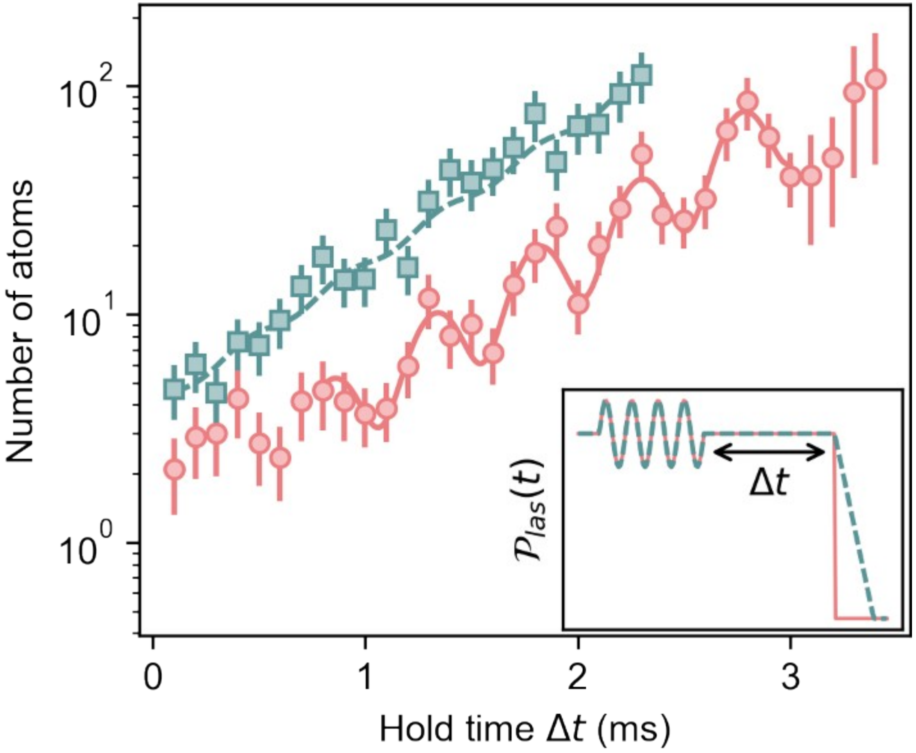
Equivalence if

$\partial_t \omega_k / \omega_k \ll \omega_k$

What we measure

$$\omega_k = \frac{\hbar k^2}{2m}$$

Eigenbasis for non-interacting atoms



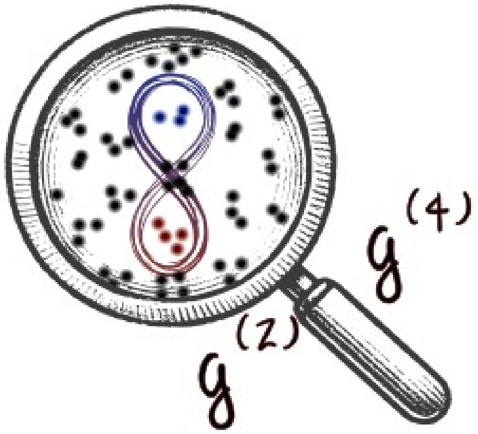
→ In the following, we slowly turn off interactions.

[1] Gondret *et al.*, Comptes Rendus. Physique **25**, 1 (2025).

We measure the full probability distribution of the quasiparticle number of a two-mode state.

→ How extract entanglement from it?

IV. Assessing entanglement of two-mode Gaussian states with many-body correlation function



Gondret *et al.*, Quantifying Two-Mode Entanglement of Bosonic Gaussian States from Their Full Counting Statistics, Phys. Rev. Lett. **135**, 100201 (2025).

What is entanglement

HOW?

Just violate a Bell inequality

Bell *Physics* (1964)
CHSH *Phys. Rev. Lett.* (1969)

Entanglement $\Leftrightarrow$ Bell inequalities
$\Leftrightarrow$ Distillability
$\Leftrightarrow$ Teleportation

EQUIVALENCE ONLY FOR PURE STATES

Gisin, *Phys. Lett. A* (1991)
Gisin & Peres, *Phys. Lett. A* (1992)
Popescu & Rohrlich, *Phys. Lett. A* (1992)

WHAT ABOUT *MIXED* STATES?

Teleportation $\not\Rightarrow$ Bell inequalities

Popescu Phys. Rev. Lett. (1994)

Define a partition 1-2 (two modes here). Any **separable** state can be written as

$$\rho = \sum_i \alpha_i \rho_{i,1} \otimes \rho_{i,2}$$

where $\alpha_i \geq 0$ are probabilities.

Other states are non-separable / entangled.

Werner *Phys. Rev. A* (1989)

How to assess entanglement?

SO HOW?

Many entanglement witnesses and criteria in the literature

PPT:

$$\hat{\rho}^{t_2} \geq 0$$

Peres, *Phys. Rev. Lett.* (1996)

$$|\langle \hat{a}_1 \hat{a}_2 \rangle|^2 \leq n_1 n_2$$

Hillery & Zubairy *Phys. Rev. Lett.* (2006)

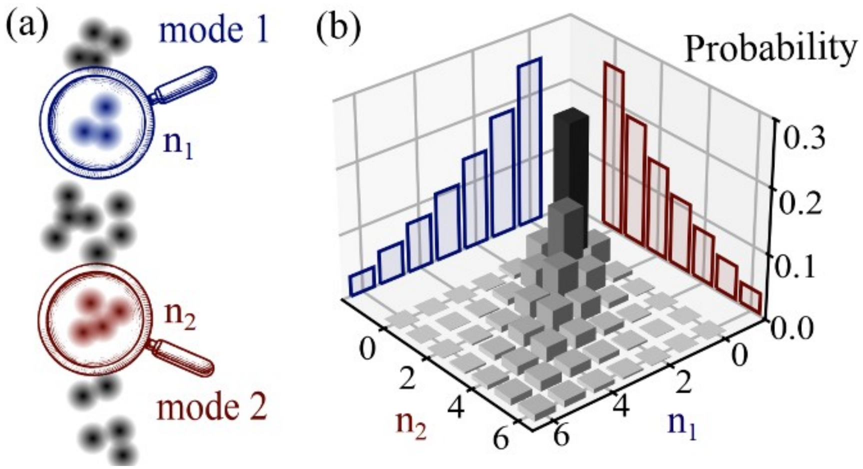
⋮

EXERIMENTAL TOOLS NEEDED

To measure the mean/variances of field operators, one needs homodyne-like detection schemes[1] or to reconstruct the state measuring non-commuting operators[2] (e.g. $\hat{x}$ and $\hat{p}$)

[1] Gross *et al.* Nature (2011)
[2] Bergschneider *et al.* Nat. Phys. (2019)

COULD WE USE THE FULL COUNTING STATISTICS?



Yields any order of *particle number* correlation function

$$G_{12}^{(m,p)} = \left\langle \left(\hat{a}_1^\dagger\right)^m \left(\hat{a}_2^\dagger\right)^p \hat{a}_1^m \hat{a}_2^p \right\rangle$$

See also Barasiński *et al* PRL (2023)

Can we assess entanglement in general from the FCS?

Entanglement definition

A (k,-k) state is entangled if it is not separable i.e. it cannot be written as

$$\hat{\rho} = \sum_i \alpha_i \hat{\rho}_{i,k} \otimes \hat{\rho}_{i,-k}$$

with $\alpha_i \geq 0$.

Werner, Phys. Rev. A **40**, 4277 (1989)

Take a two-mode squeezed state

$$\hat{\rho}_{TMS} \propto \sum_{i,n} \kappa^i \kappa^n |i\rangle\langle i|_k \otimes |n\rangle\langle n|_{-k}$$

→ Can we prove it is entangled from its full counting statistics?

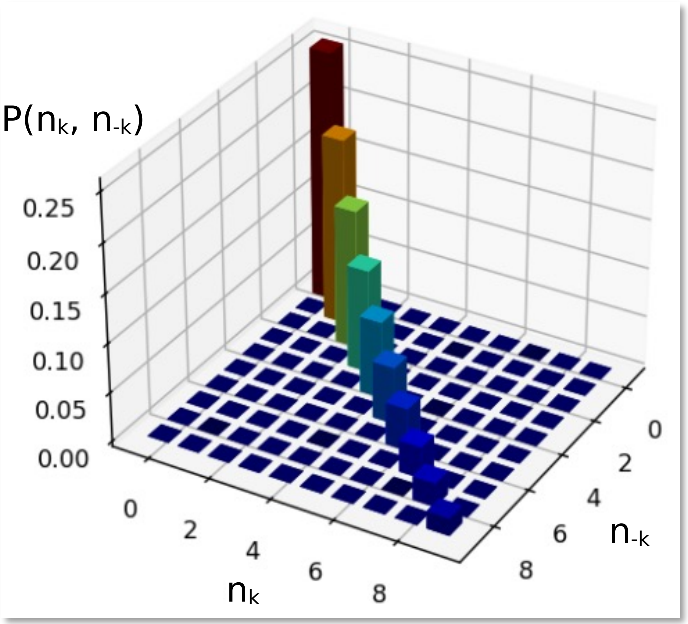


Fig: full counting statistics of a two-mode squeezed state

NO
(in general)

Ex: the following separable state has the same full counting statistics

$$\hat{\rho}_{separ} \propto \sum_i \kappa^i |i\rangle\langle i|_k \otimes |i\rangle\langle i|_{-k}$$

This state IS separable

Conclusion?

One cannot assess the entanglement of *any* quantum state from its full counting statistics.

It only measures the diagonal terms of the density matrix

So... thank you ??

Wait a minute... this is not true for *Gaussian* states!

Two-body correlation function to witness entanglement

DEFINITION

A Gaussian state is defined by its 1st and 2nd moments: it has vanishing cumulants of order > 2.

PROPERTIES

Any operator that involves more than 2 fields $\hat{a}^{(\dagger)}$ can be expressed with 1- and 2-field operators.

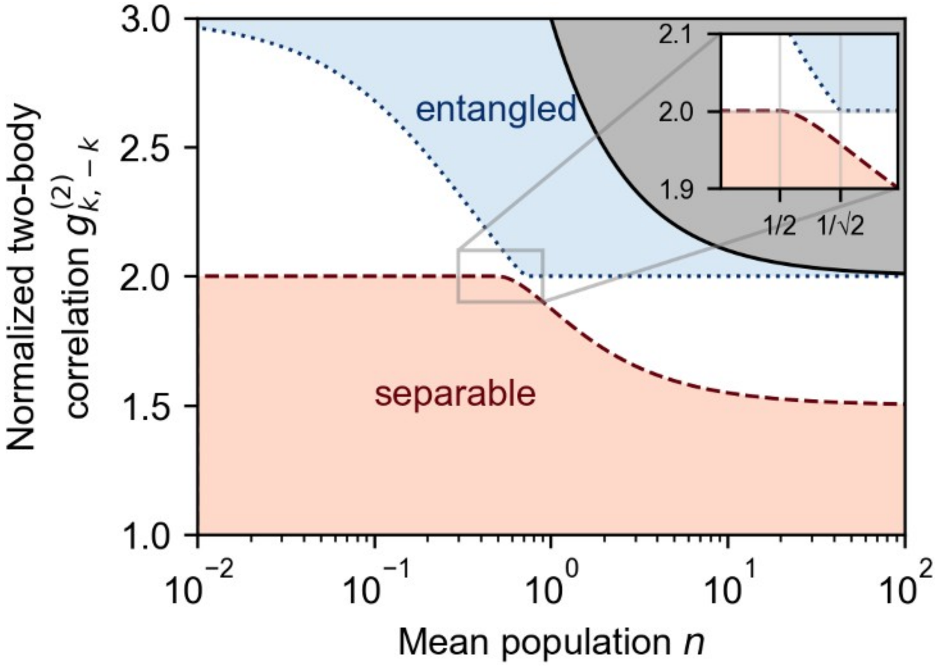
Ex: for zero mean Gaussian states,

$$G^{(2)}_{-k,k} = \langle \hat{a}^\dagger_k \hat{a}^\dagger_{-k} \hat{a}_k \hat{a}_{-k} \rangle = n_k n_{-k} + |\langle \hat{a}_k \hat{a}_{-k} \rangle|^2 + |\langle \hat{a}_k \hat{a}^\dagger_{-k} \rangle|^2$$

Populations

Entanglement is guaranteed when these quantities are sufficiently greater than the populations

If we assume that each mode exhibits a thermal probability distribution, $G^{(2)}$ is an entanglement witness[1]:

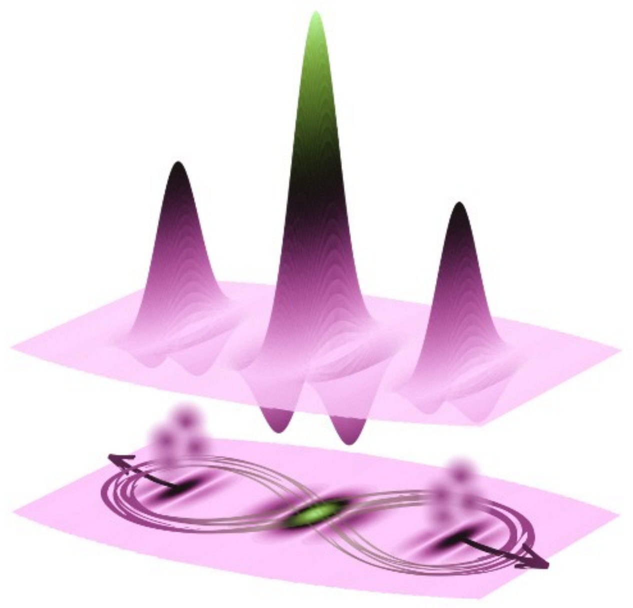


If we also measure the four-body correlation function, entanglement can be certified and quantified through logarithm negativity [1].

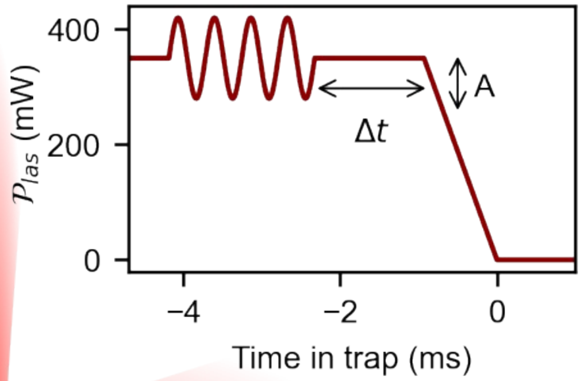
[1] Gondret *et al.*, Phys. Rev. Lett. **135**, 100201 (2025)

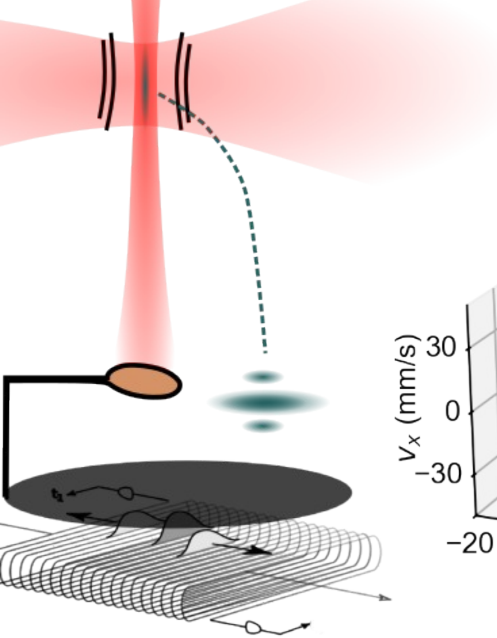
V. Observation of entanglement between collective excitations in a parametrically driven BEC

Gondret *et al.*, Observation of Entanglement in a Cold Atom Analog of Cosmological Preheating, Phys. Rev. Lett. **135**, 240603 (2025)



Thermal single mode probability distribution



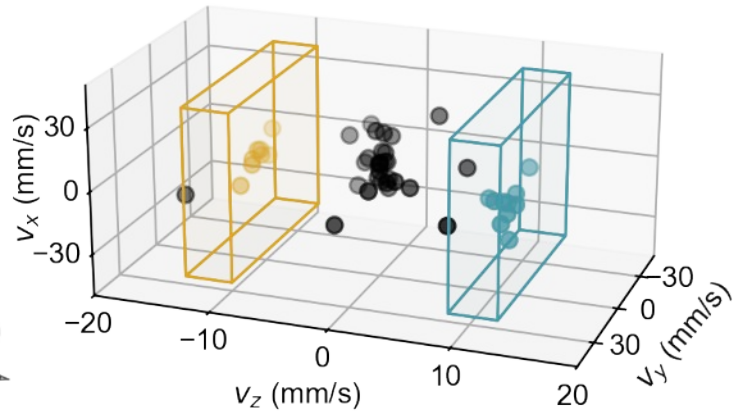


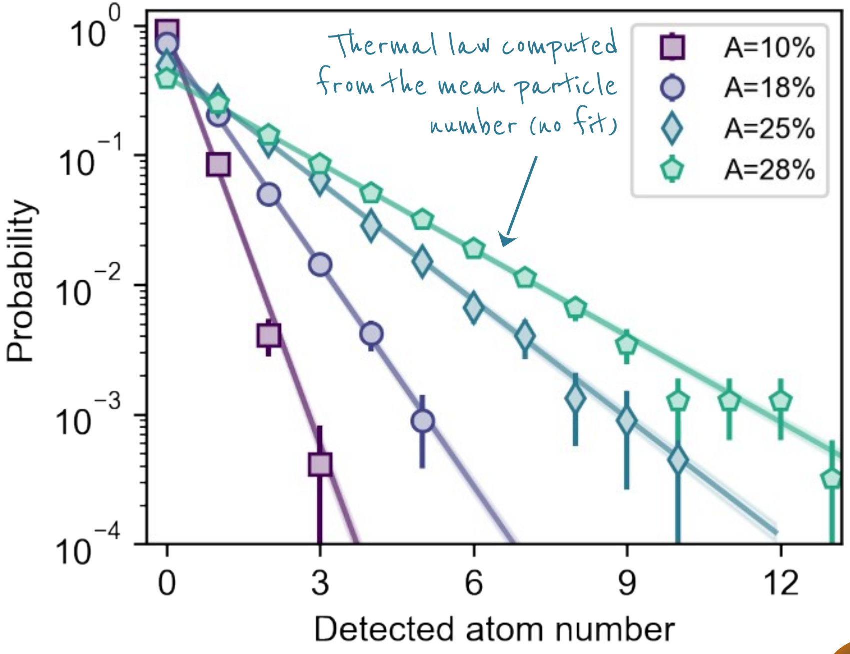
Protocol: We repeat the experiment varying the amplitude of the excitation A.

Few thousand repetitions per point, 3 weeks 24/7 in total

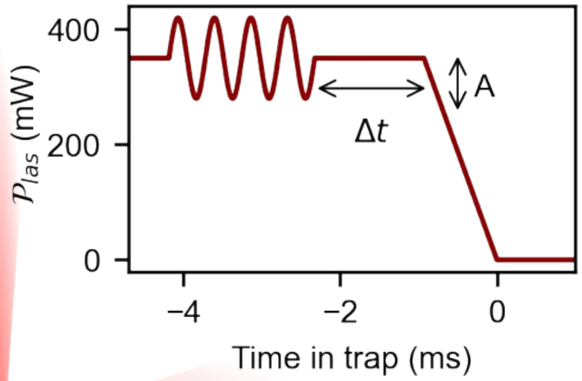
Measurement: probability distribution of each mode

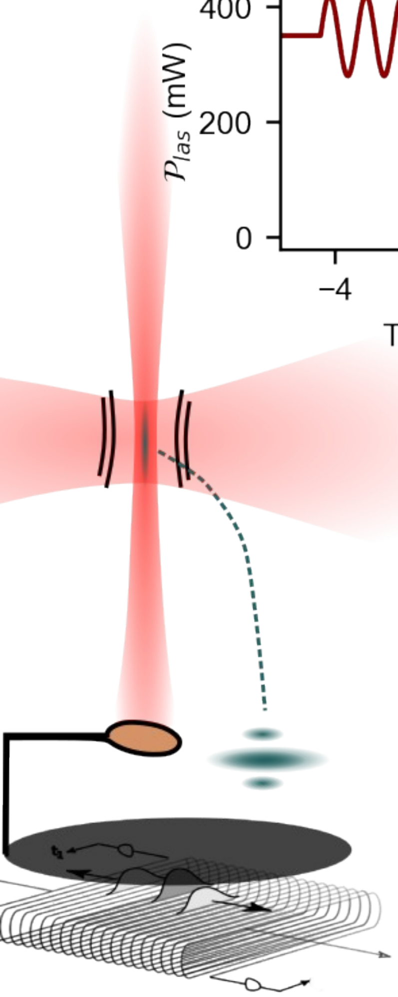
Result: thermal single mode probability distribution (i.e. bosonic bunching)





Thermal single mode probability distribution





Protocol: We repeat the experiment varying the amplitude of the excitation A.

Measurement: cross normalized two-body correlation function $g^{(2)}$

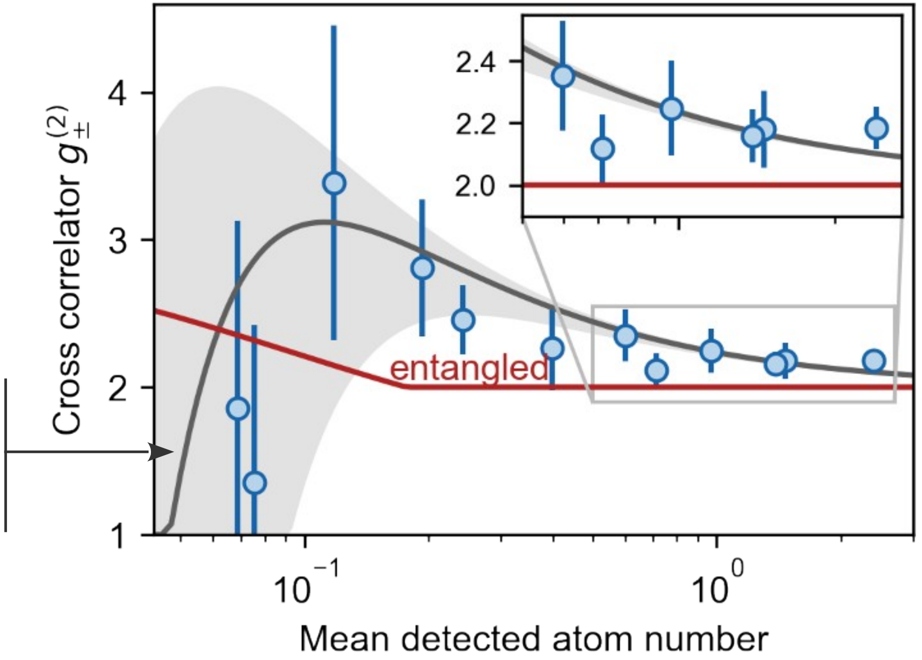
Result: assuming the two-mode state is Gaussian, it is entangled for sufficiently large excitation amplitude.

Model: two-mode squeezed thermal state without free parameter.

- 25(5)% detector efficiency,
- 25(5) nK temperature. Fluctuation of

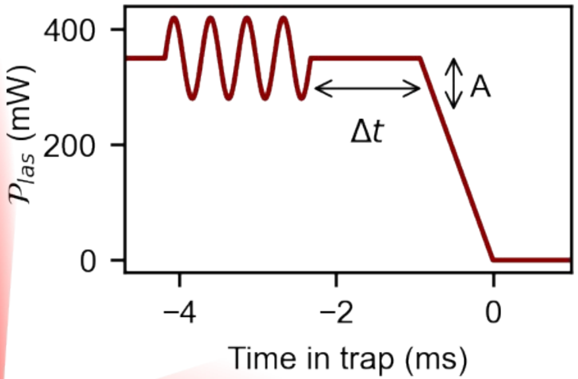
$$0.5 + 0.18(8)$$

Vacuum ↗ ↖ Thermal



Gondret *et al.*, Phys. Rev. Lett. **135**, 240603 (2025)

Continuing the driving



Protocol: vary the excitation duration Δt

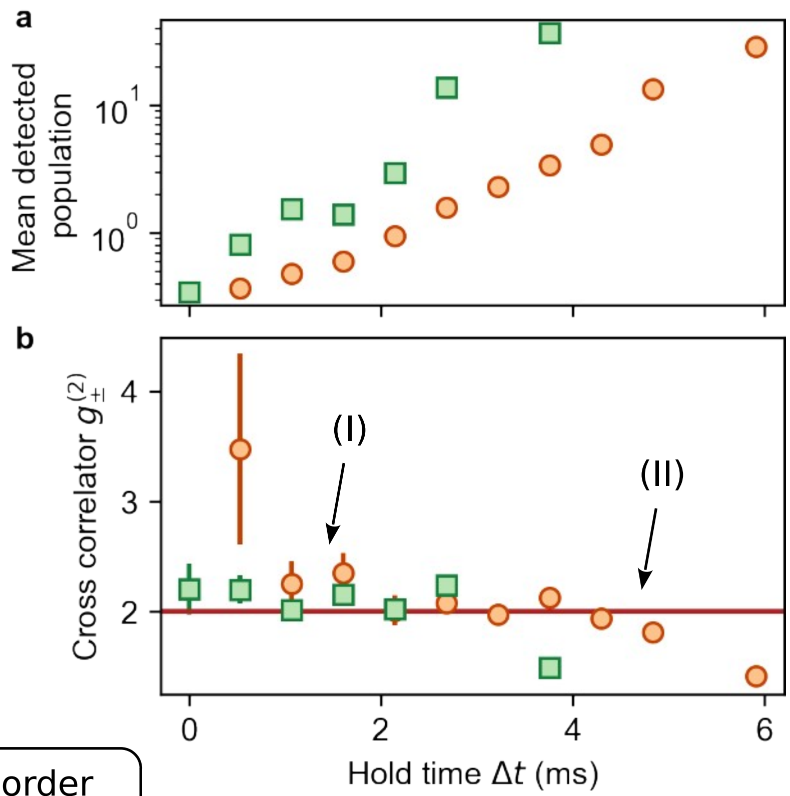
Expected in the two-mode squeezing model.

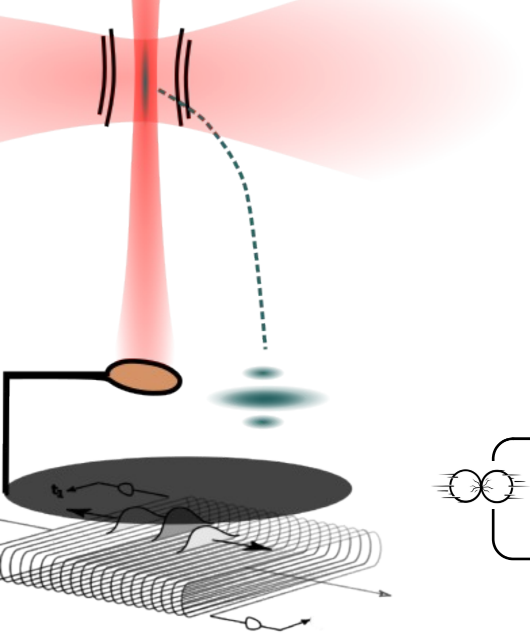
Results:

(I) $g^{(2)} \to 2$ as population grows

(II) At later time, $g^{(2)}$ drops below 2

Not expected





Onset of a late-time regime where higher order quasi-particle interactions become relevant.

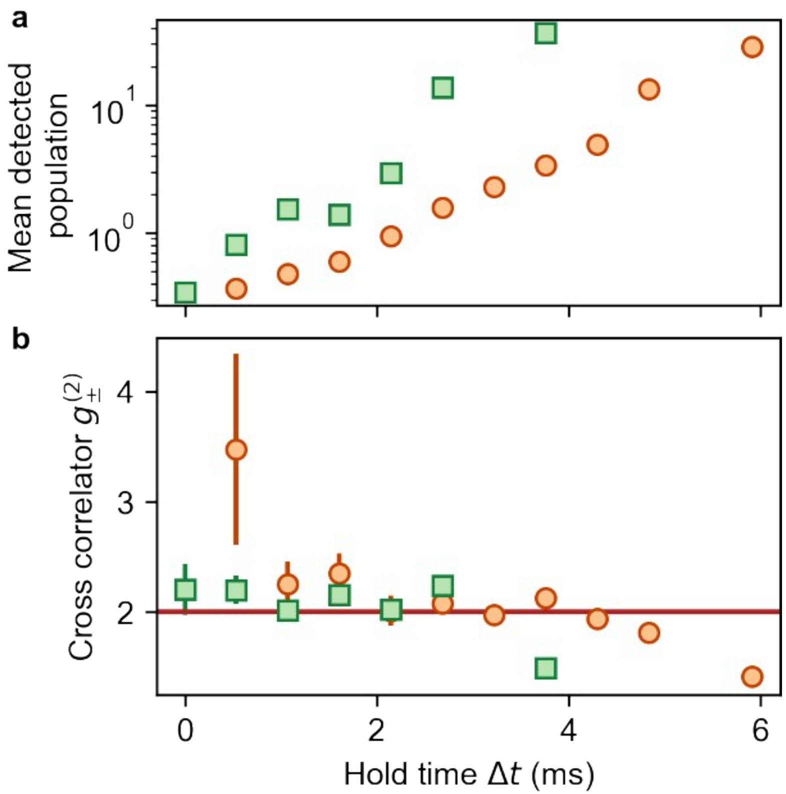
Gondret *et al.*, Phys. Rev. Lett. **135**, 240603 (2025)

Outlooks

- Onset of a late-time regime where higher order quasi-particle interactions become relevant.

Towards the study of the much-less understood interaction-dominated regime:

- decoherence of the resonant modes,
  Robertson *et al.*, Phys. Rev. D **98**, 056003 (2018)
- loss of Gaussianity,
  Schweigler *et al.*, Nat. Phys. **17**, 559 (2021)
  Bureik *et al.*, Nat. Phys. **21**, 57 (2025)
- appearance of higher order peaks,
  Gregory *et al.*, arXiv:2410.08842 (2025)
- back-reaction of the quasiparticles on the BEC...
  Butera and I. Carusotto, Phys. Rev. Lett. **130**, 241501 (2023)



Gondret *et al.*, Phys. Rev. Lett. **135**, 240603 (2025)

Summary and Outlooks

Take home message

- Analogy between quasiparticles in a fluid and particles in curved space time,
- Two-mode entanglement of Gaussian state can be assessed from particle number correlation,
- Observation of vacuum amplification through entanglement between quasiparticles in a BEC,
- Towards a non-Gaussian regime

From Gaussian matter-waves manipulated with light...

...to non-Gaussian light upon interacting with matter !

- Detection of non-classicality of non-Gaussian states,
- Generation of bright non-classical light,

¿Estados Gaussianos son fome verdad?

Take home message

- Analogy between quasiparticles in a fluid and particles in curved space time,
- Two-mode entanglement of Gaussian state can be assessed from particle number correlation,
- Observation of vacuum amplification through entanglement between quasiparticles in a BEC.

Thank you for your attention!

Refs

- arXiv 2411.09284
- arXiv 2503.09555
- arXiv 2506.22024
- arXiv 2508.01654

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with Clothilde Lamirault, Rui Dias, Léa Camier, Amaury Micheli, Charlie Leprince, Quentin Marolleau, Scott Robertson, Denis Boiron, Christoph I. Westbrook