



Cosmology and quantum simulation

Non separability of phonon pairs in a time-modulated BEC linked to inflationary scenarii

Victor Gondret, Charlie Leprince, Quentin Marolleau, Marc Cheneau,
Denis Boiron, Chris Westbrook, Amaury Micheli and Scott Robertson

Equipe Optique Atomique Quantique



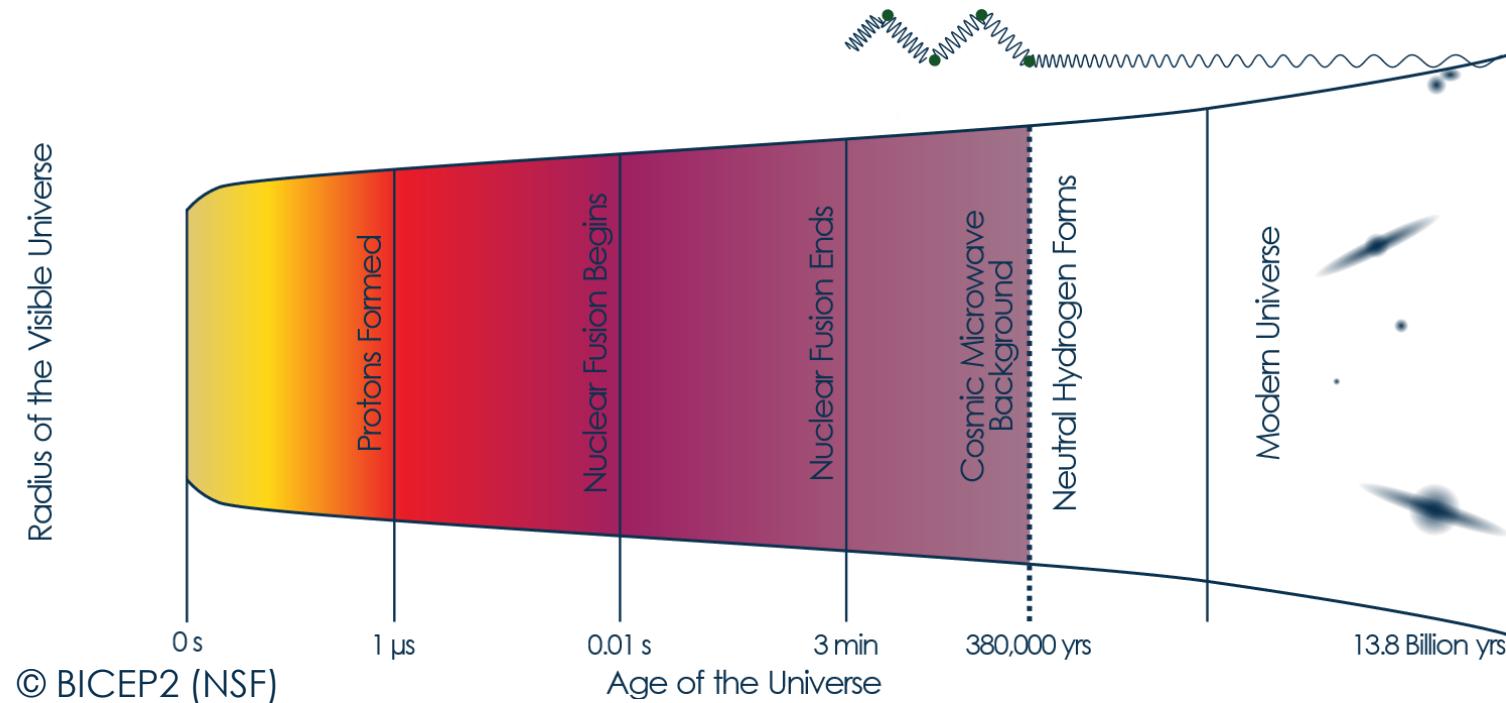
Outline

- ① Inflation in cosmology
- ② Analog gravity and Bose-Einstein Condensates
- ③ Work in progress @LCF

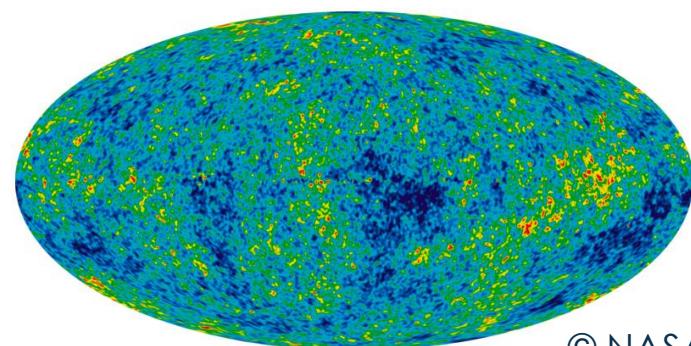
① Inflation in cosmology

The Standard Model

Radius of the Visible Universe



Cosmic microwave background at 2.7254 (6) K.



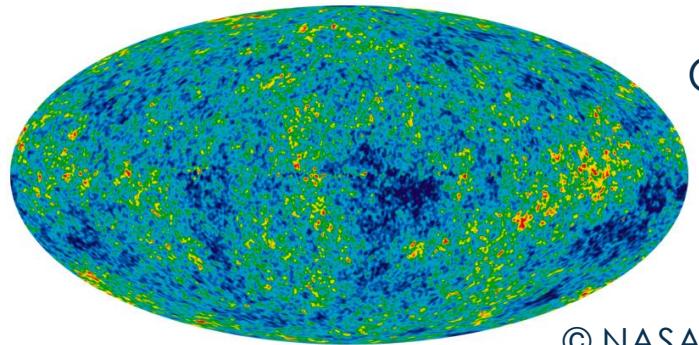
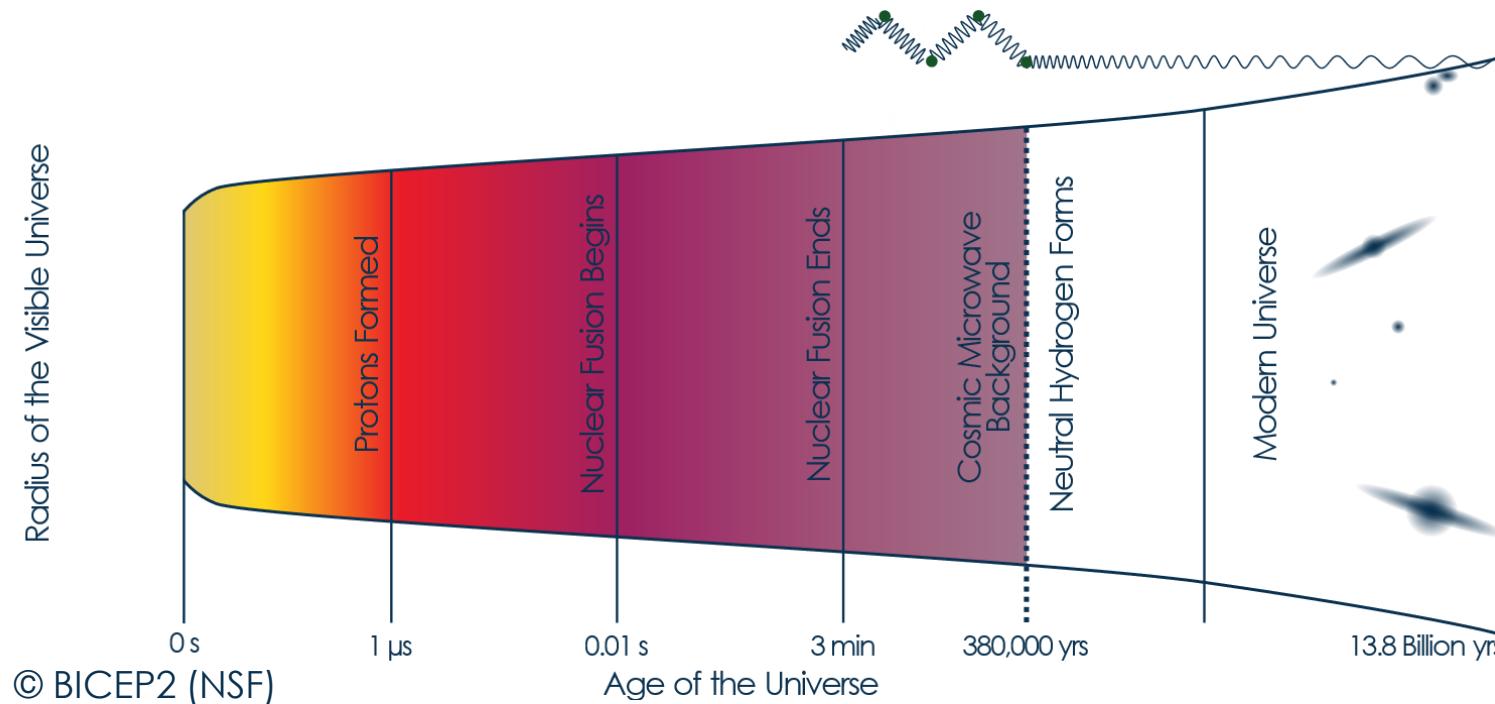
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Successes of the Standard Model

- 👍 Successful account for nucleosynthesis and the relative abundance of light elements,
- 👍 Predicts the fact that universe is expanding (Hubble law)
- 👍 Right prediction of the Cosmic Background Radiation temperature

The Standard Model

Radius of the Visible Universe



Cosmic microwave background at 2.7254 (6) K.

Failures of the Standard Model

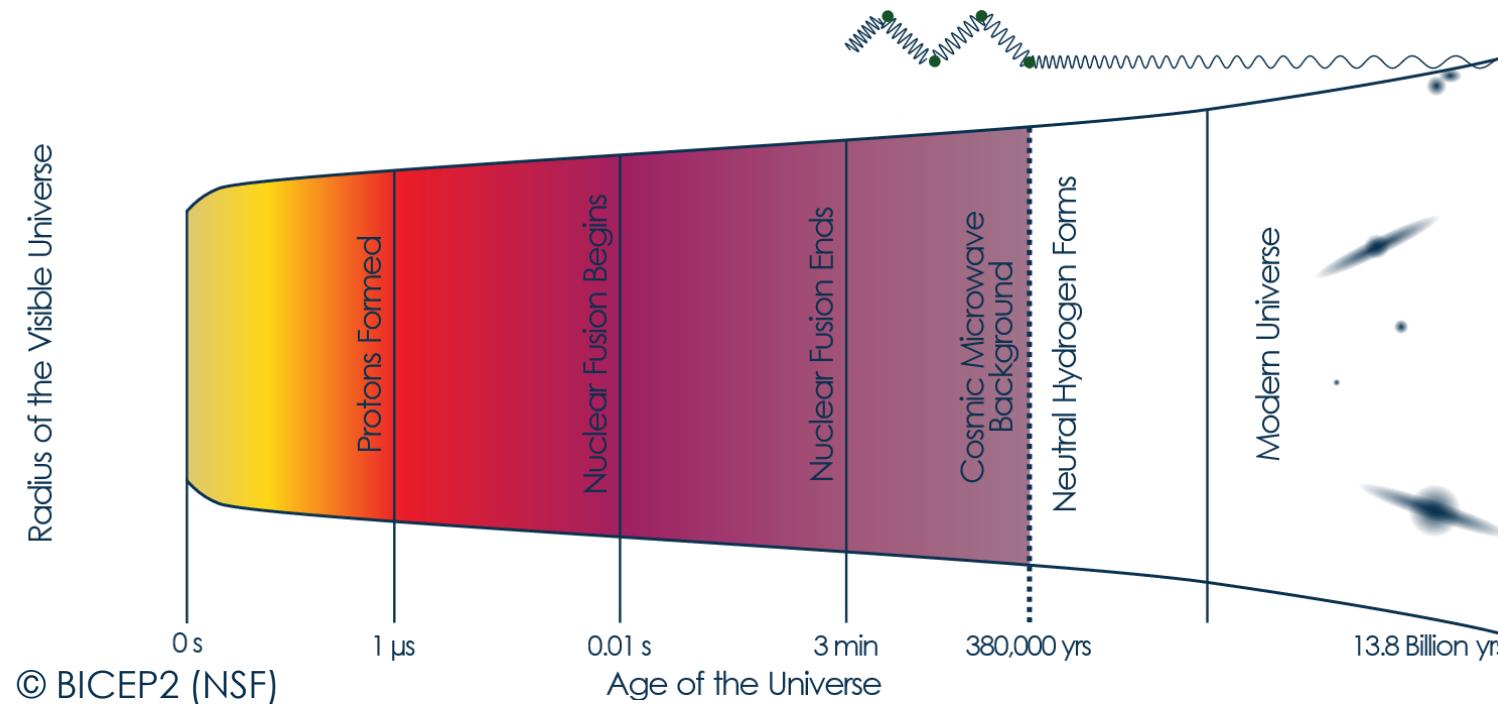
- 👎 The horizon problem
- 👎 The flatness problem
- 👎 ...

Horizon problem

Homogeneity of causally disconnected regions of space

The Standard Model

Radius of the Visible Universe



Failures of the Standard Model

- 👎 The horizon problem
- 👎 The flatness problem
- 👎 ...

Flatness problem

$$\frac{\rho}{\rho_C} \underset{1,02 \pm 0,02 \text{ now}}{\approx} 1 \quad \text{REALLY equal to one}$$

Energy density

Critical energy density

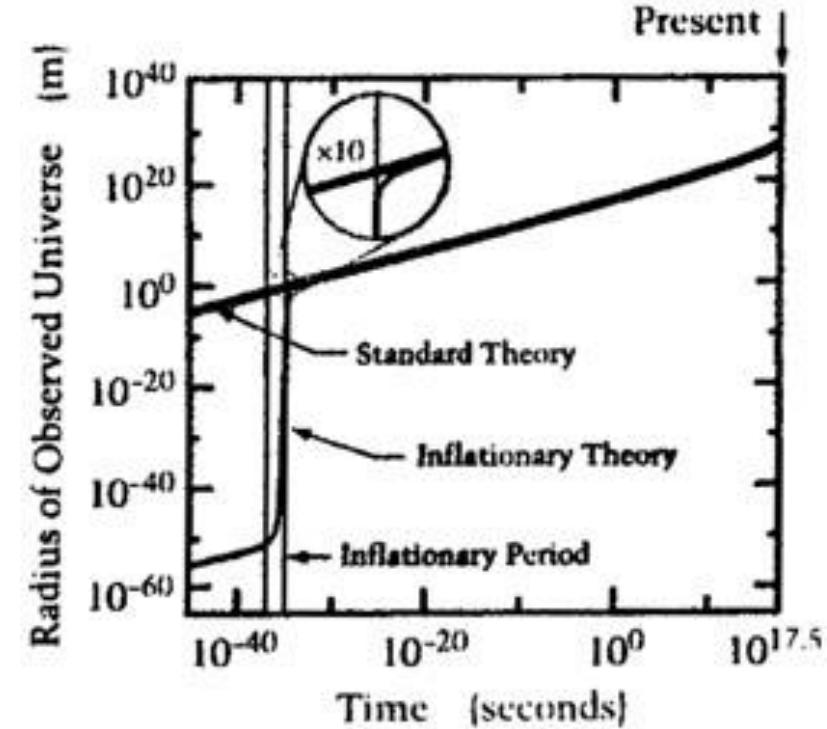
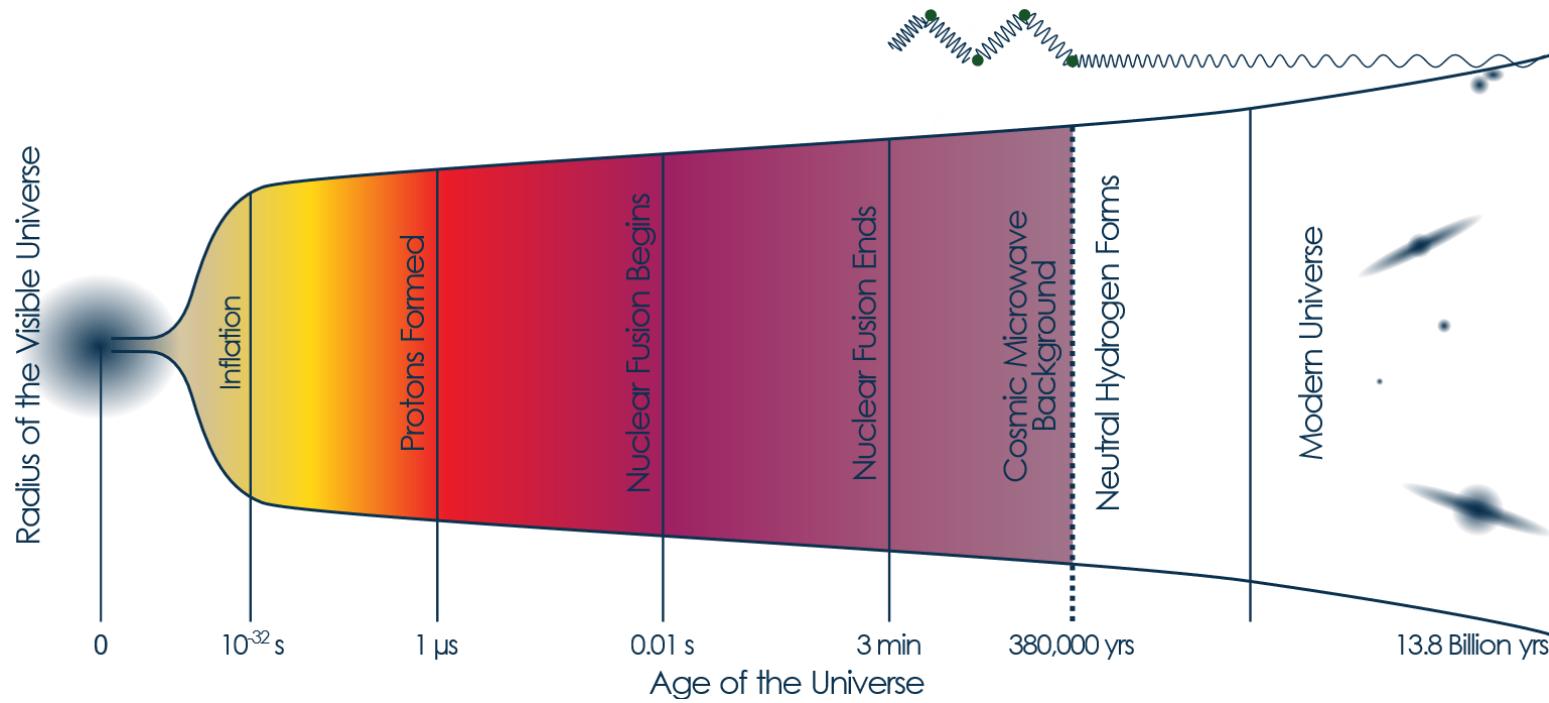
Horizon problem

Homogeneity of causally disconnected regions of space

A.H. Guth, Beem Line 27(3), 1997

S. Watson, An Exposition on Inflationary Cosmology, 2000

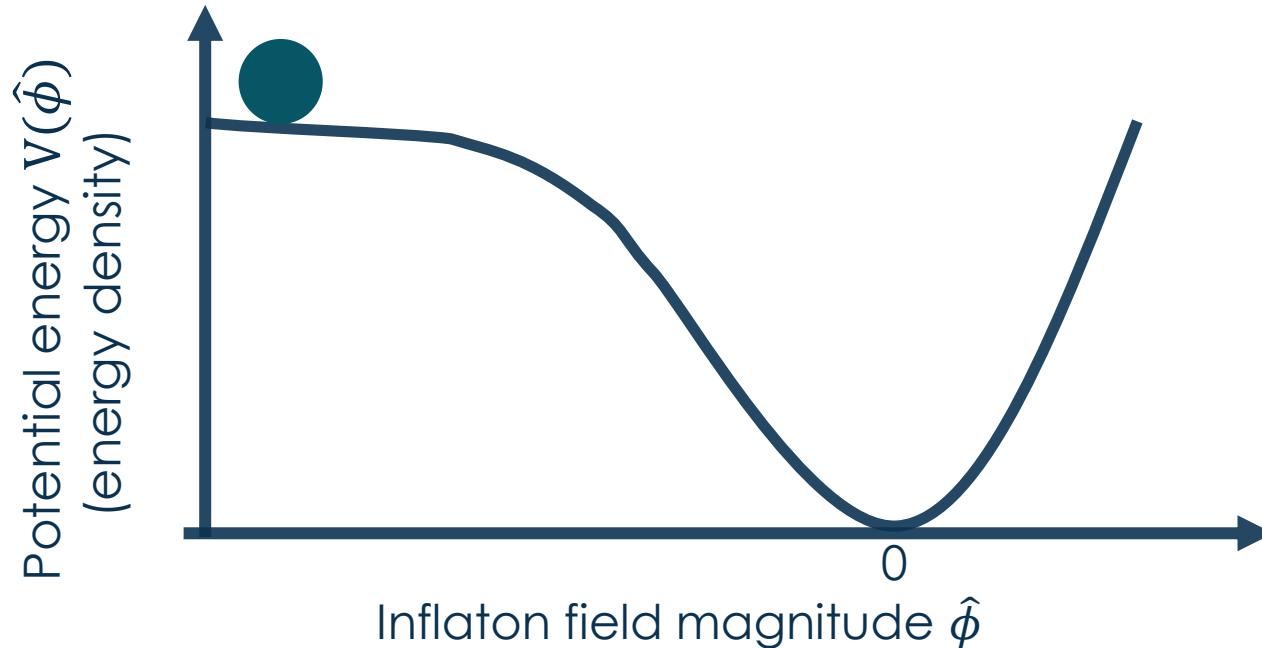
Theory of Inflation



- 👍 Horizon problem : regions are no more causally connected & the inflation washes out inhomogeneities
- 👍 Flatness problem : no matter initial conditions, inflation drives ρ/ρ_c to 1.

Guth, A.H., 1981. Inflationary universe: A possible solution to the horizon and flatness problems. Phys. Rev. D 23, 347–356.

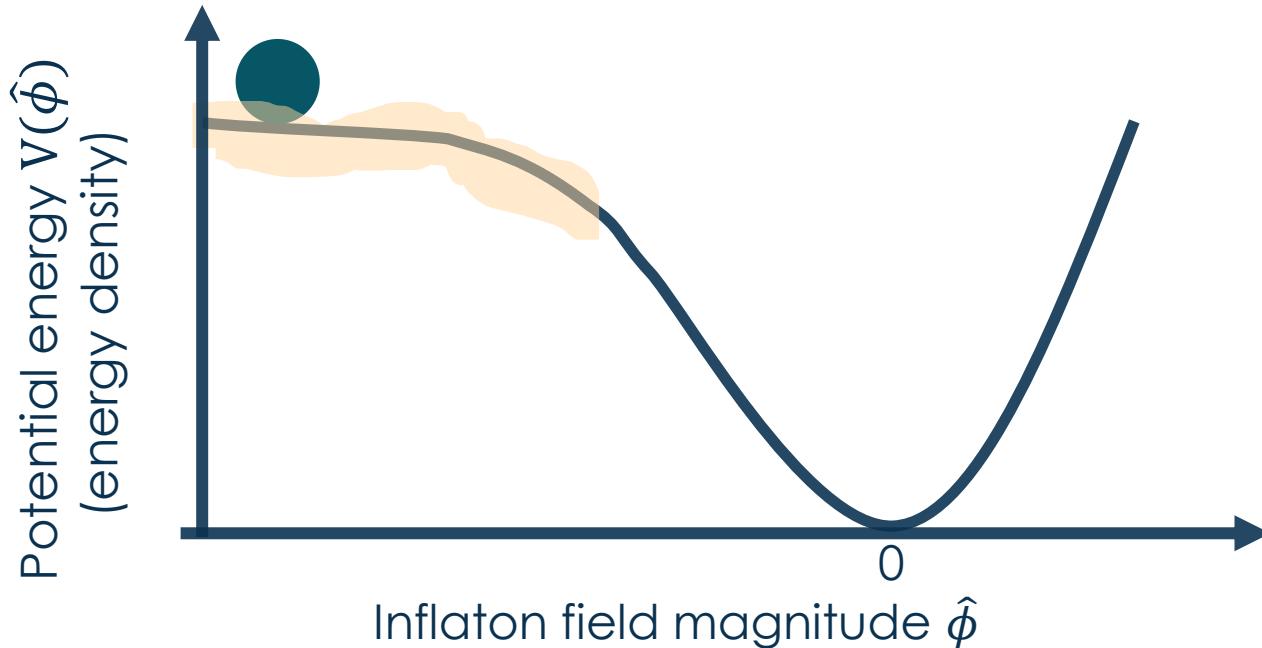
Theory of Inflation



Add a new field $\hat{\phi}$ called *inflaton*

A. H. Guth, Phys. Rev. D23, 347 (1981)

Theory of Inflation



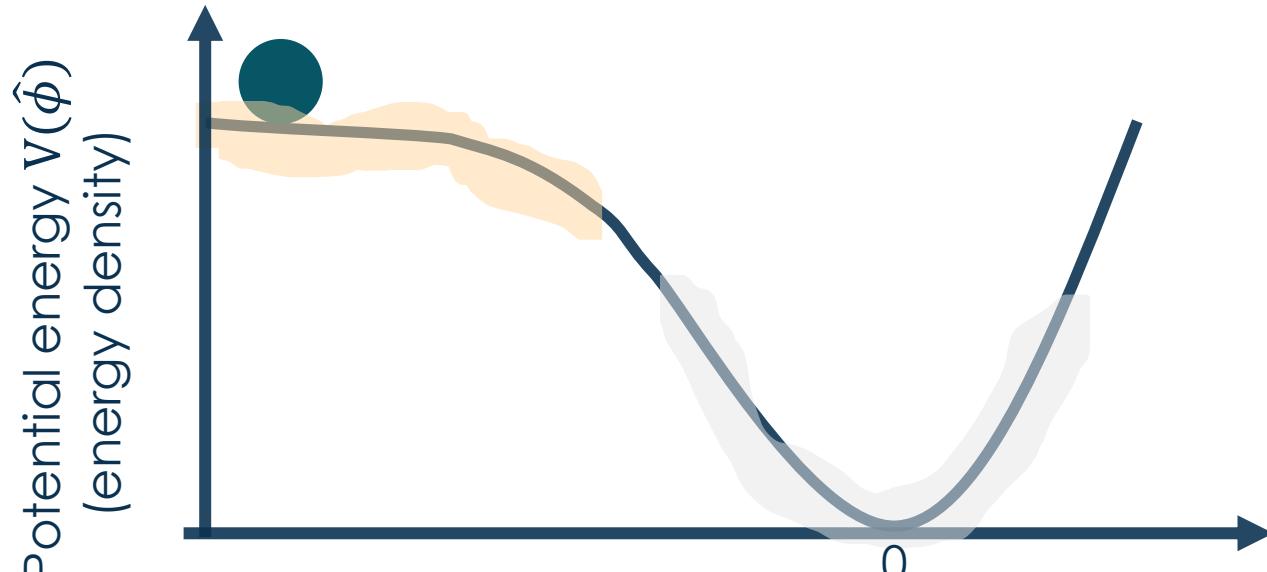
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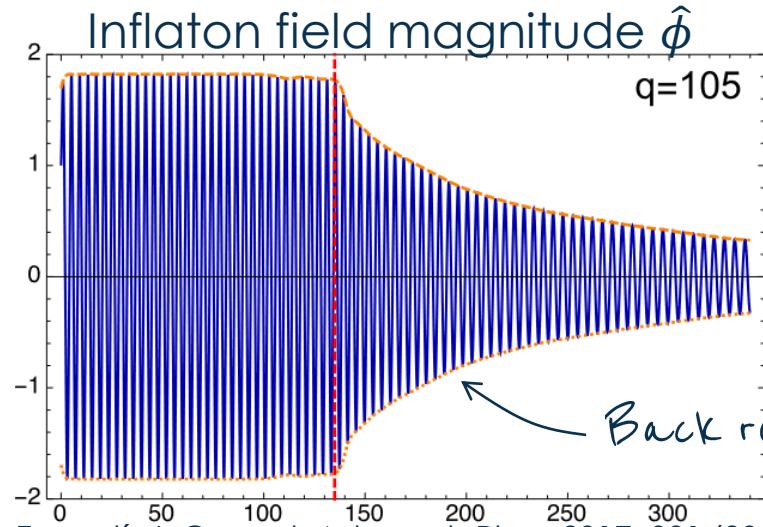
The inflaton *slowly rolls* from its initial state.
Its almost constant potential energy drives
the inflation.

A. Linde, Phys. Lett. 129B, 177 (1983).

Theory of Inflation



Decay of the inflaton field into particles



Add a new field $\hat{\phi}$ called *inflaton*

A. H. Guth, Phys. Rev. D23, 347 (1981)

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A. Linde, Phys. Lett. 129B, 177 (1983).

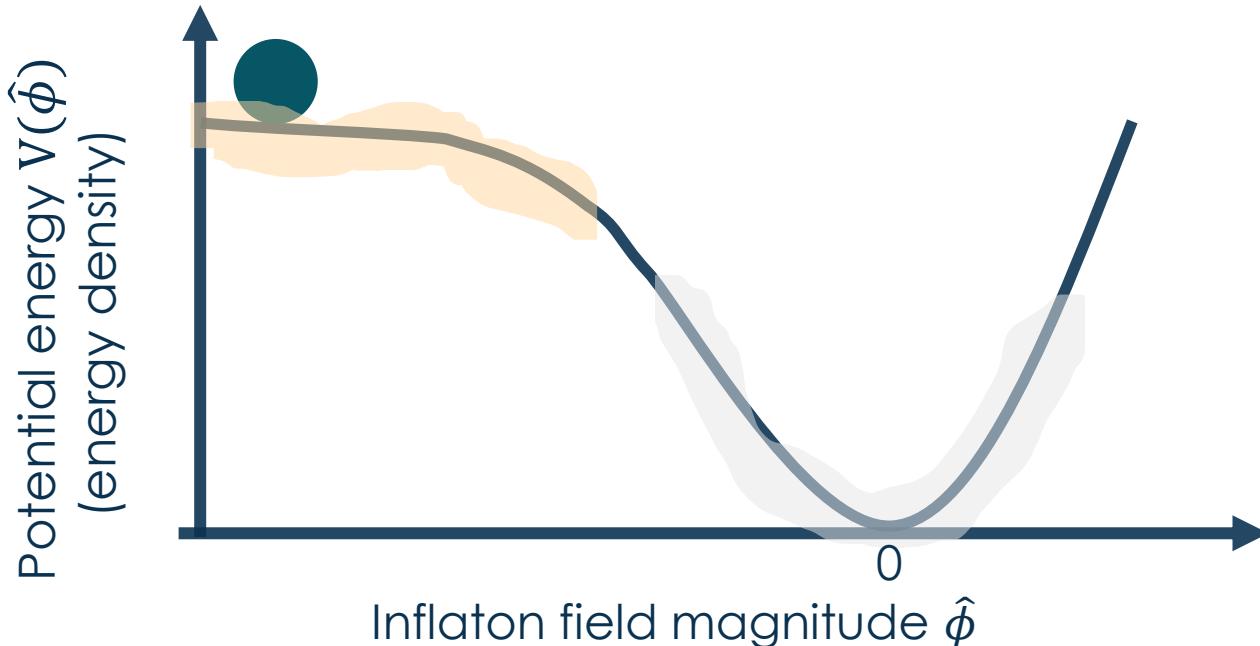
The inflaton starts to oscillate around its minimum and, coupled to matter fields, it creates particles through parametric resonance.

L. Kofman, A. Linde & A. Starobinsky, Phys. Rev. D 56, (1997).

A.H. Guth, Beem Line 27(3), 1997

S. Watson, An Exposition on Inflationary Cosmology, 2000

Theory of Inflation



$$V(\phi) = \lambda(\phi^2 - M^2)^2$$

Higgs potential

$$V(\phi) = \frac{1}{2}m^2\phi^2$$

Massive scalar field

$$V(\phi) = \lambda\phi^4$$

Self interacting scalar field

$$V(\phi) = 2H_i^2 \left(3 - \frac{1}{s}\right) e^{-\phi/\sqrt{s}}$$

String theory

Add a new field $\hat{\phi}$ called *inflaton*

A. H. Guth, Phys. Rev. D23, 347 (1981)

The inflaton slowly rolls from its initial state. Its almost constant potential energy drives the inflation.

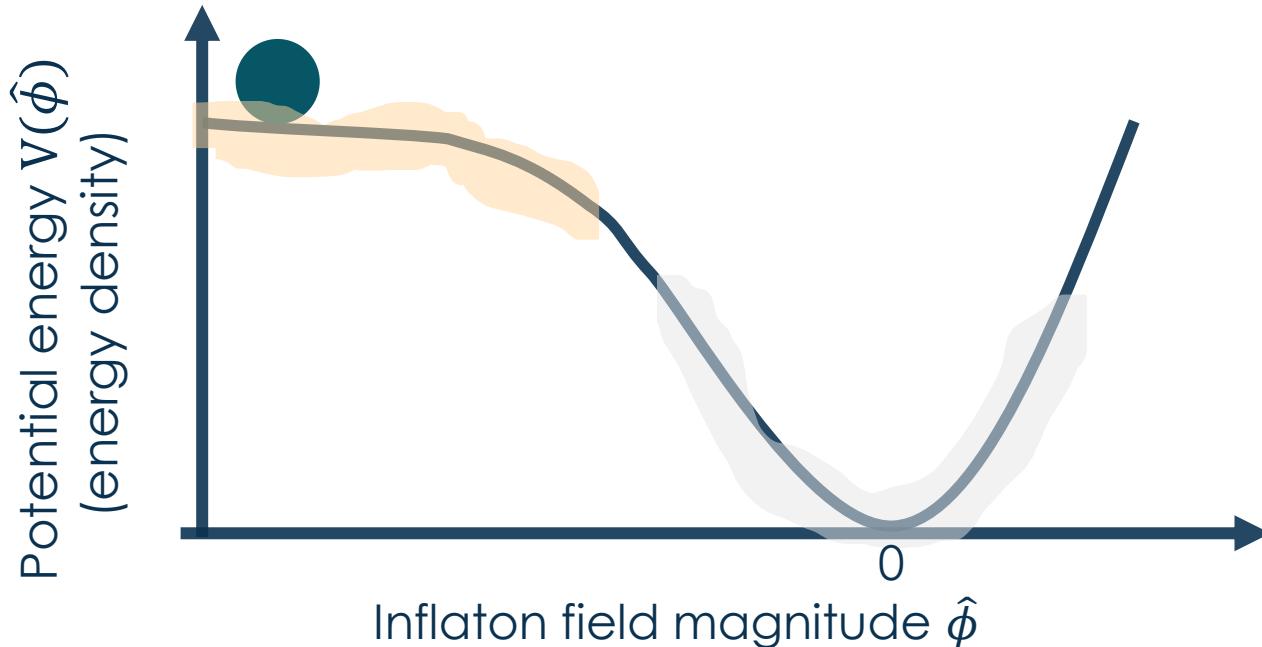
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A.H. Guth, Beem Line 27(3), 1997
S. Watson, An Exposition on Inflationary Cosmology, 2000

Theory of Inflation



Particles are created in pairs with opposite momenta in a two modes squeezed state. Interactions lead to decoherence and thermalization

D. Campo & R. Parentani, Phys. Rev. D **74**, 025001 (2006).

Add a new field $\hat{\phi}$ called *inflaton*

A. H. Guth, Phys. Rev. D23, 347 (1981)

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A.H. Guth, Beem Line 27(3), 1997
S. Watson, An Exposition on Inflationary Cosmology, 2000

② Analog gravity and Bose-Einstein Condensates

Analog gravity

PHYSICAL REVIEW LETTERS

VOLUME 46

25 MAY 1981

NUMBER 21

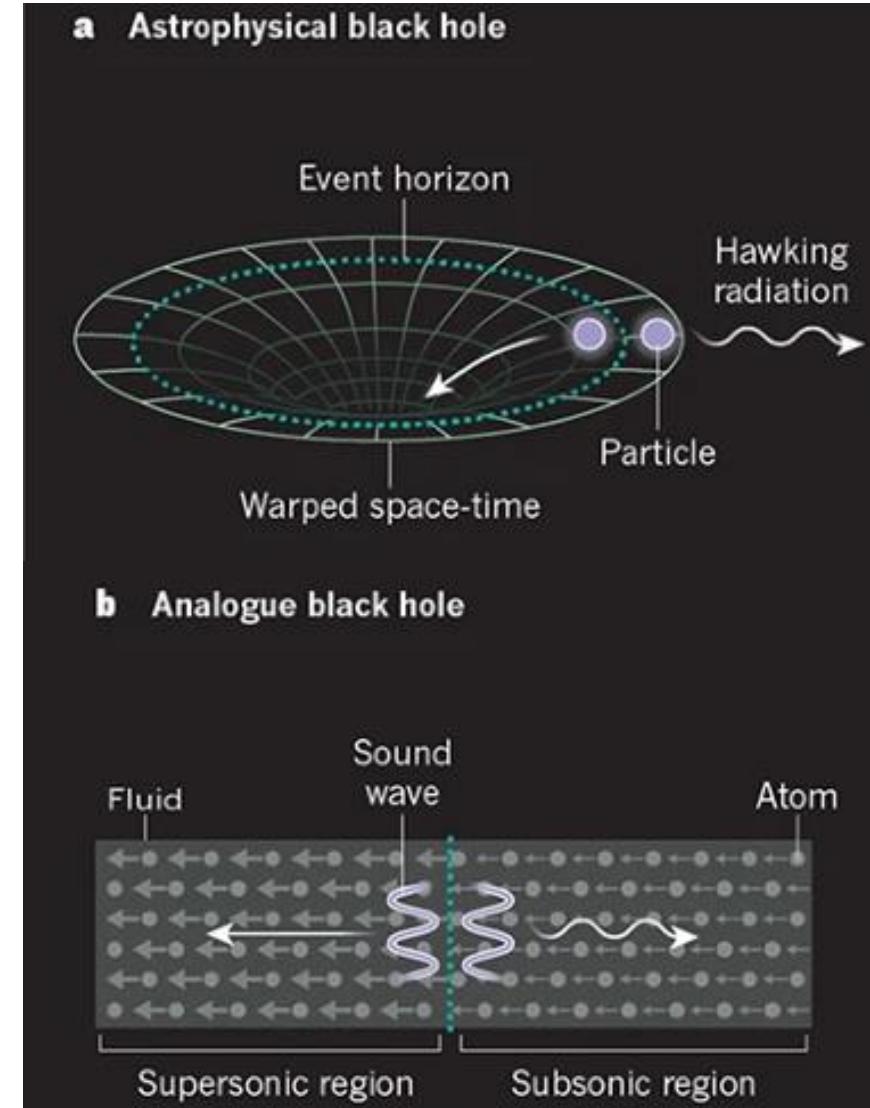
Experimental Black-Hole Evaporation?

W. G. Unruh

Department of Physics, University of British Columbia, Vancouver, British Columbia V6T 2A6, Canada

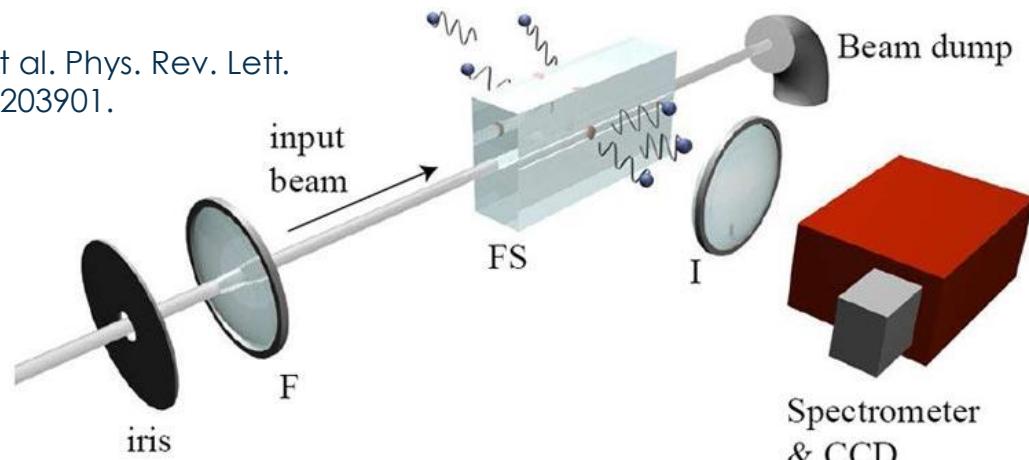
(Received 8 December 1980)

It is shown that the same arguments which lead to black-hole evaporation also predict that a thermal spectrum of sound waves should be given out from the sonic horizon in transsonic fluid flow.

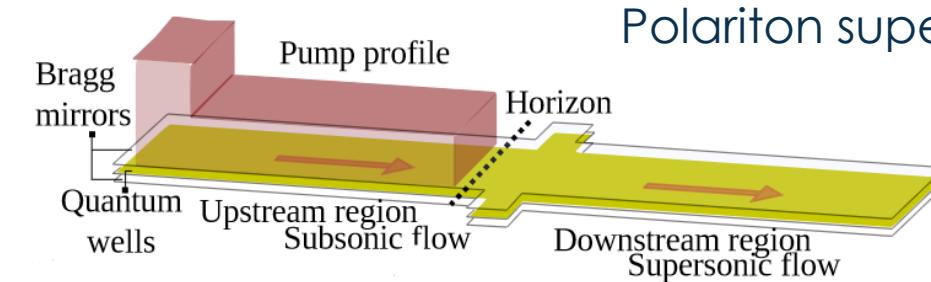


Analog gravity setups

F. Belgiorno, et al. Phys. Rev. Lett.
105.20 (2010): 203901.



Ultrashort laser pulse filaments



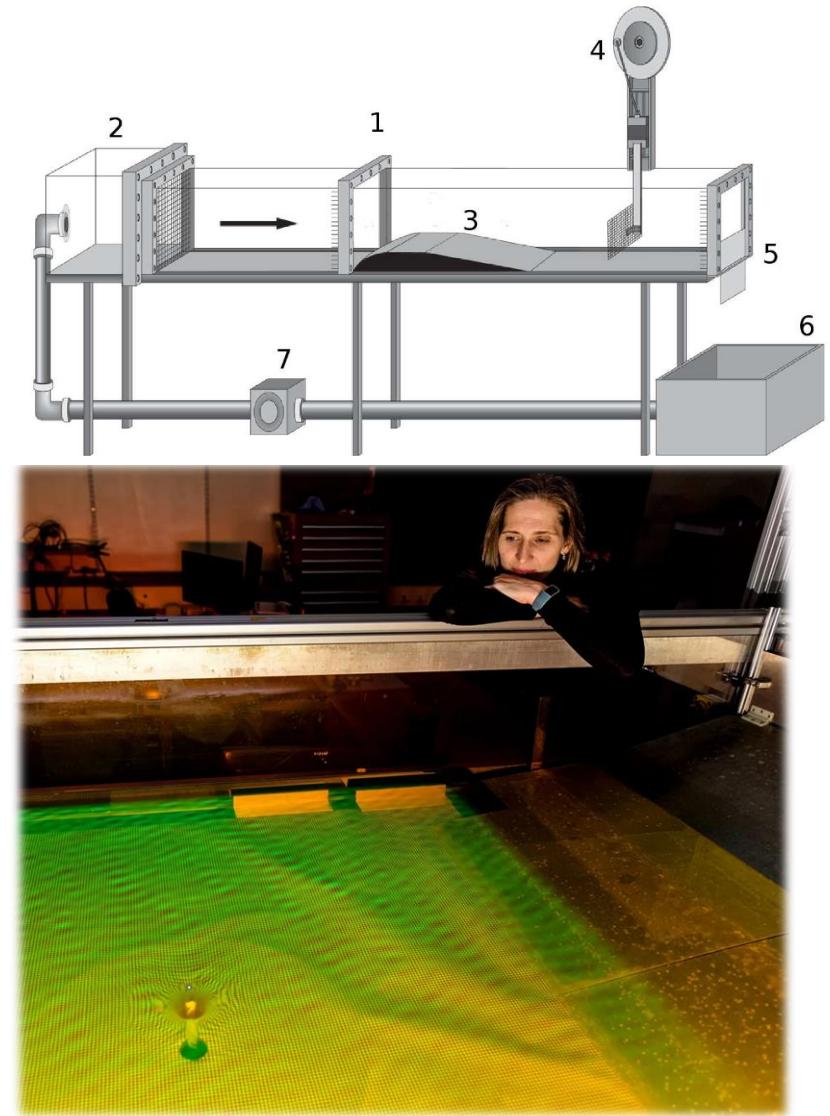
Polariton superfluid

M. Jacquet et al., Eur. Phys. J. D **76**, 152 (2022).

(...)

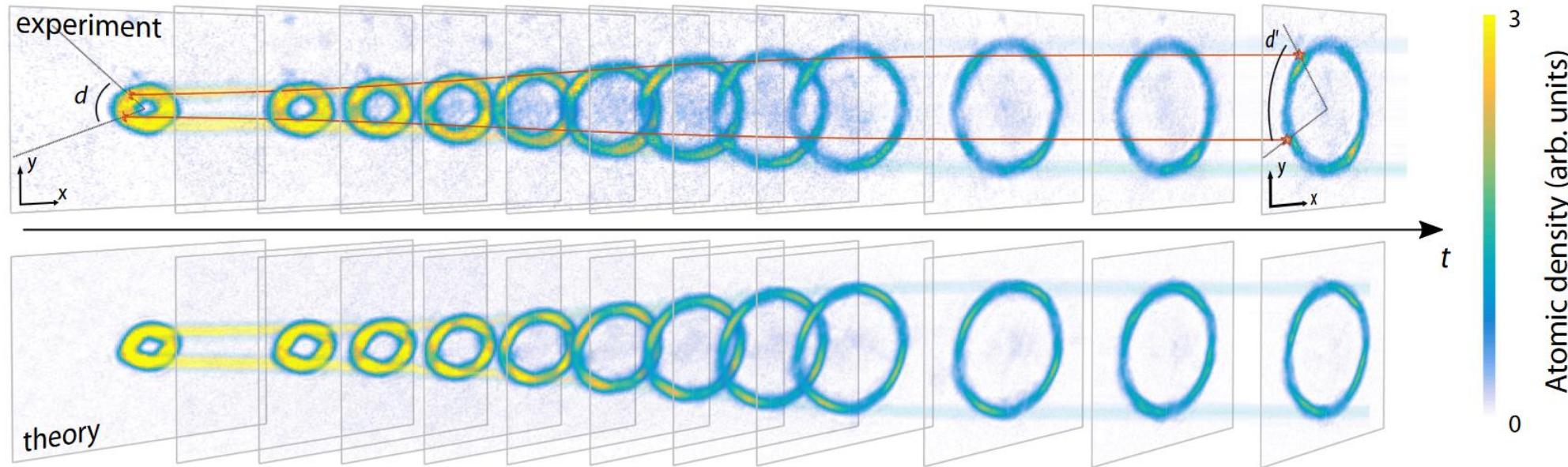
→ M. J. Jacquet, S. Weinfurter, and F. König, *The next Generation of Analogue Gravity Experiments*, Phil. Trans. R. Soc. A. **378**, 20190239 (2020).

Water tank



S. Weinfurter et al., Phys. Rev. Lett.
106, 021302 (2011).

BEC for analog inflation

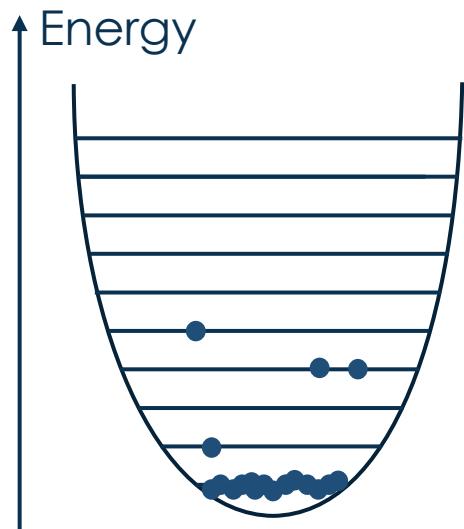


Experiment at Maryland, USA : a ring BEC is expanded at supersonic speed [1].

- [1] S. Eckel et al., A Rapidly Expanding Bose-Einstein Condensate: An Expanding Universe in the Lab, *Phys. Rev. X* **8**, 021021 (2018).
- [2] A. Chatrchyan et al., Analog Cosmological Reheating in an Ultracold Bose Gas, *Phys. Rev. A* **104**, 023302 (2021).
- [3] C. Viermann et al., Quantum Field Simulator for Dynamics in Curved Spacetime, *Nature* **611**, 260 (2022).

BEC for analog inflation

BEC is a macroscopic state at $k=0$

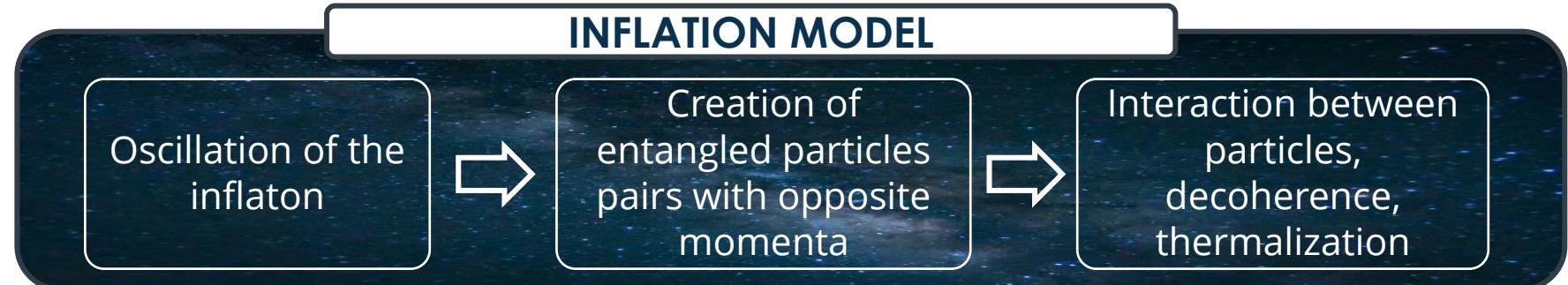
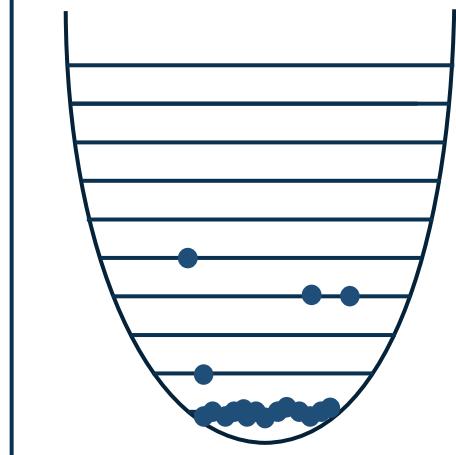


S. Robertson, F. Michel, and R. Parentani,
Phys. Rev. D **95**, 065020 (2017).
X. Busch and R. Parentani, Phys. Rev. D **88**,
045023 (2013).

BEC for analog inflation

BEC is a macroscopic state at $k=0$

Energy



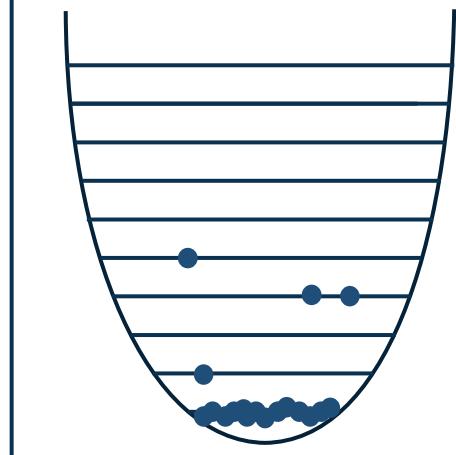
Preheating

Reheating

BEC for analog inflation

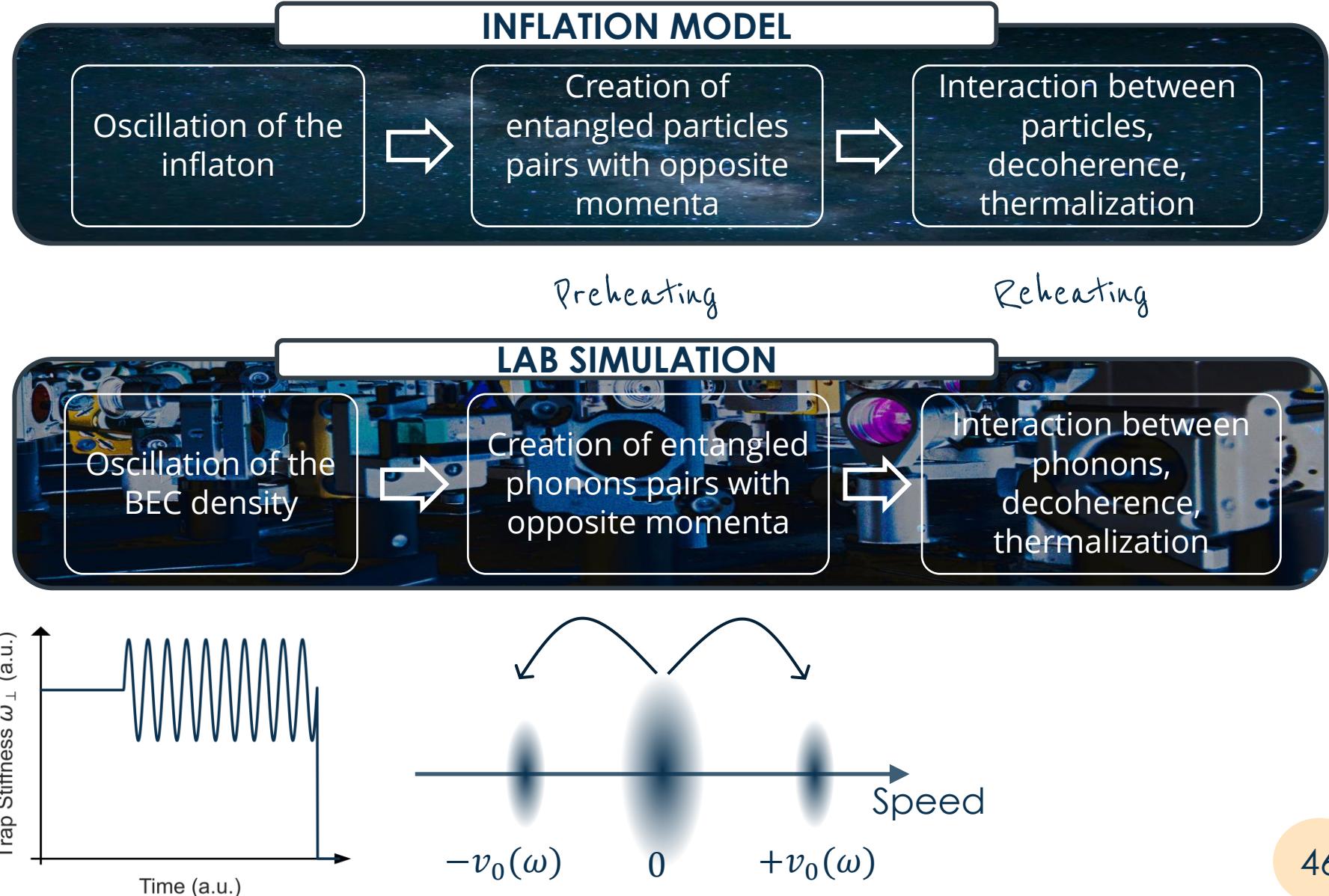
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Energy



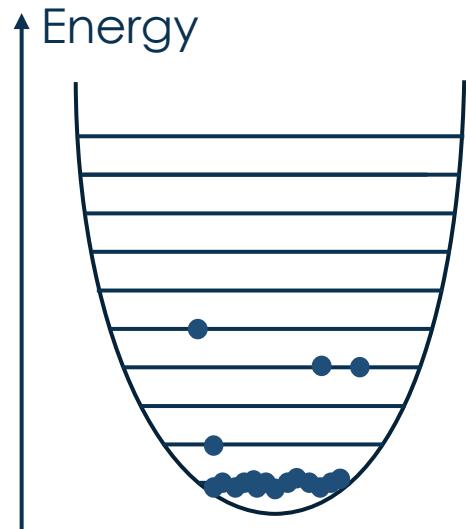
ω : excitation frequency

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045023 (2013).



BEC for analog inflation

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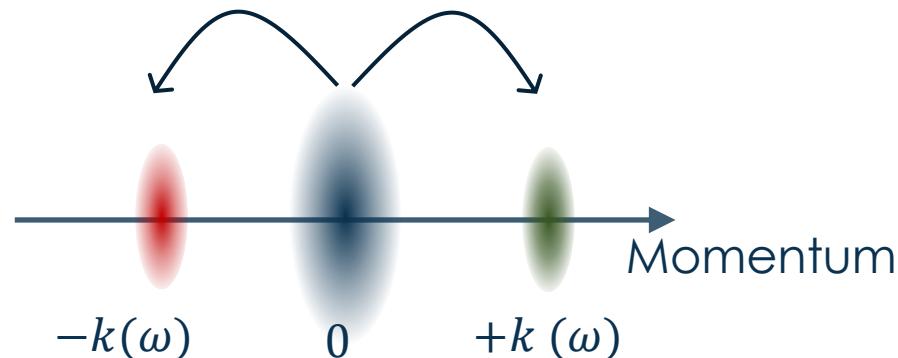
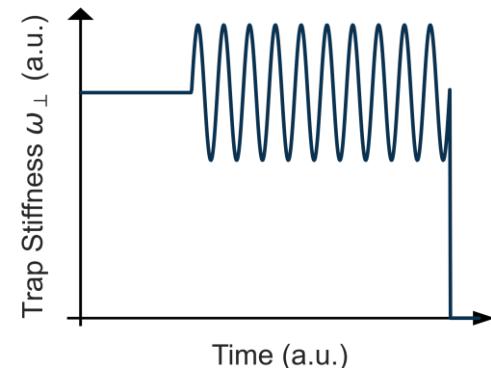
Second order correlation function :

Probing entanglement

$$g^{(2)}(\mathbf{k}_1, \mathbf{k}_2) = \frac{\langle n_{\mathbf{k}1} n_{\mathbf{k}2} \rangle}{\langle n_{\mathbf{k}1} \rangle \langle n_{\mathbf{k}2} \rangle}$$

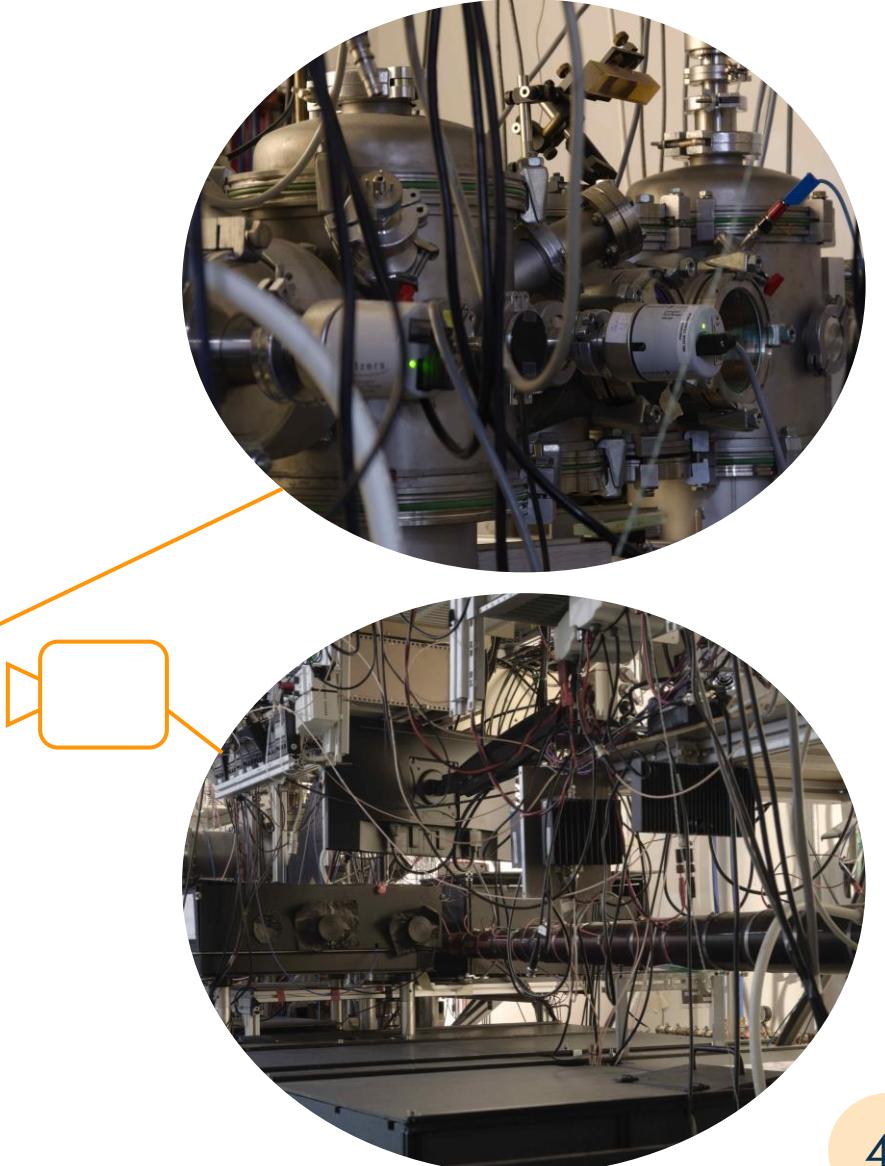
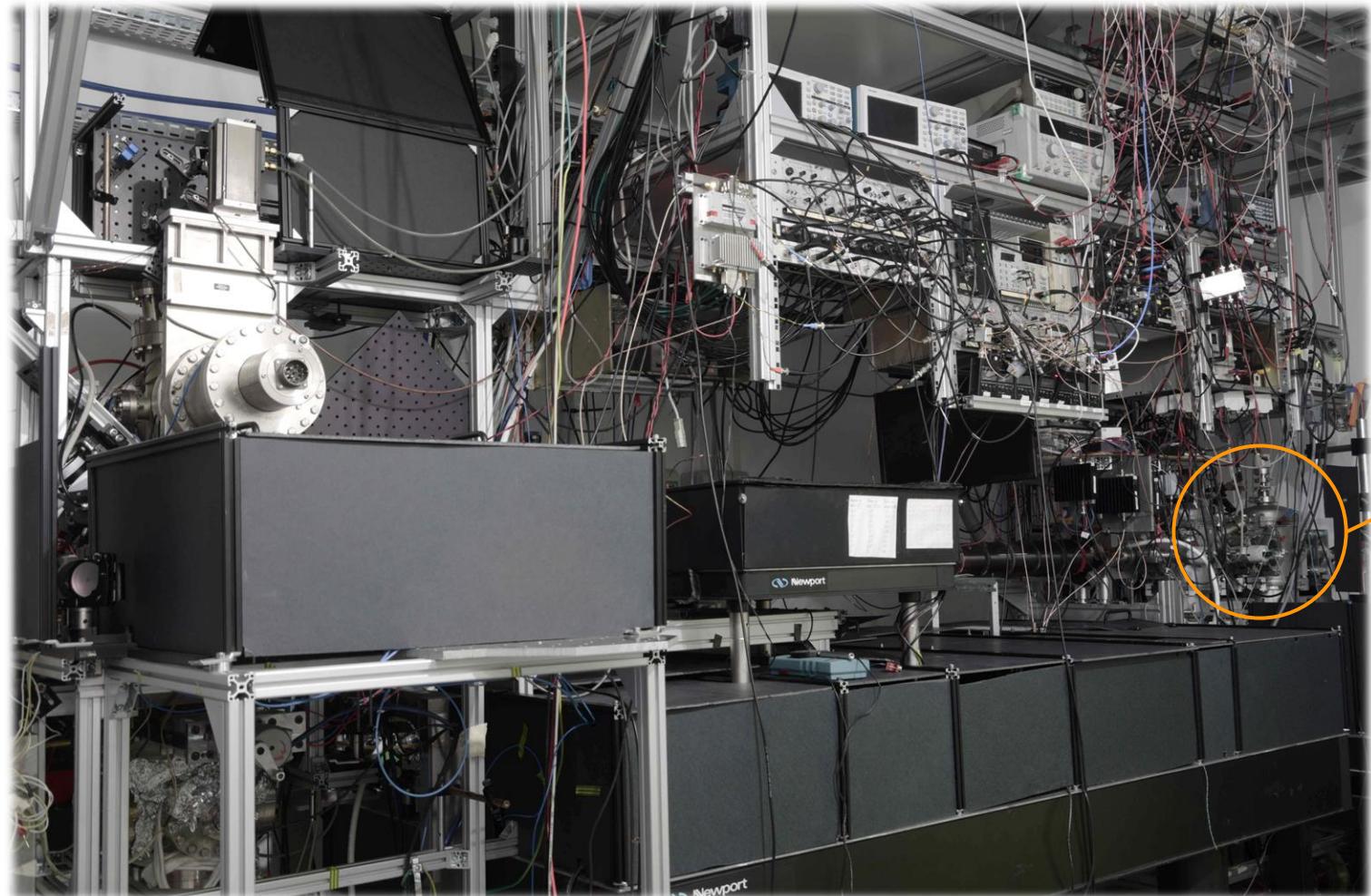
⟨ ... ⟩ : average over experimental realizations

- $g^{(2)}(k, -k) > 2$ for entangled (quasi) particles with momentum k
- $\hat{a}_k \hat{a}_{-k}^\dagger = 0$

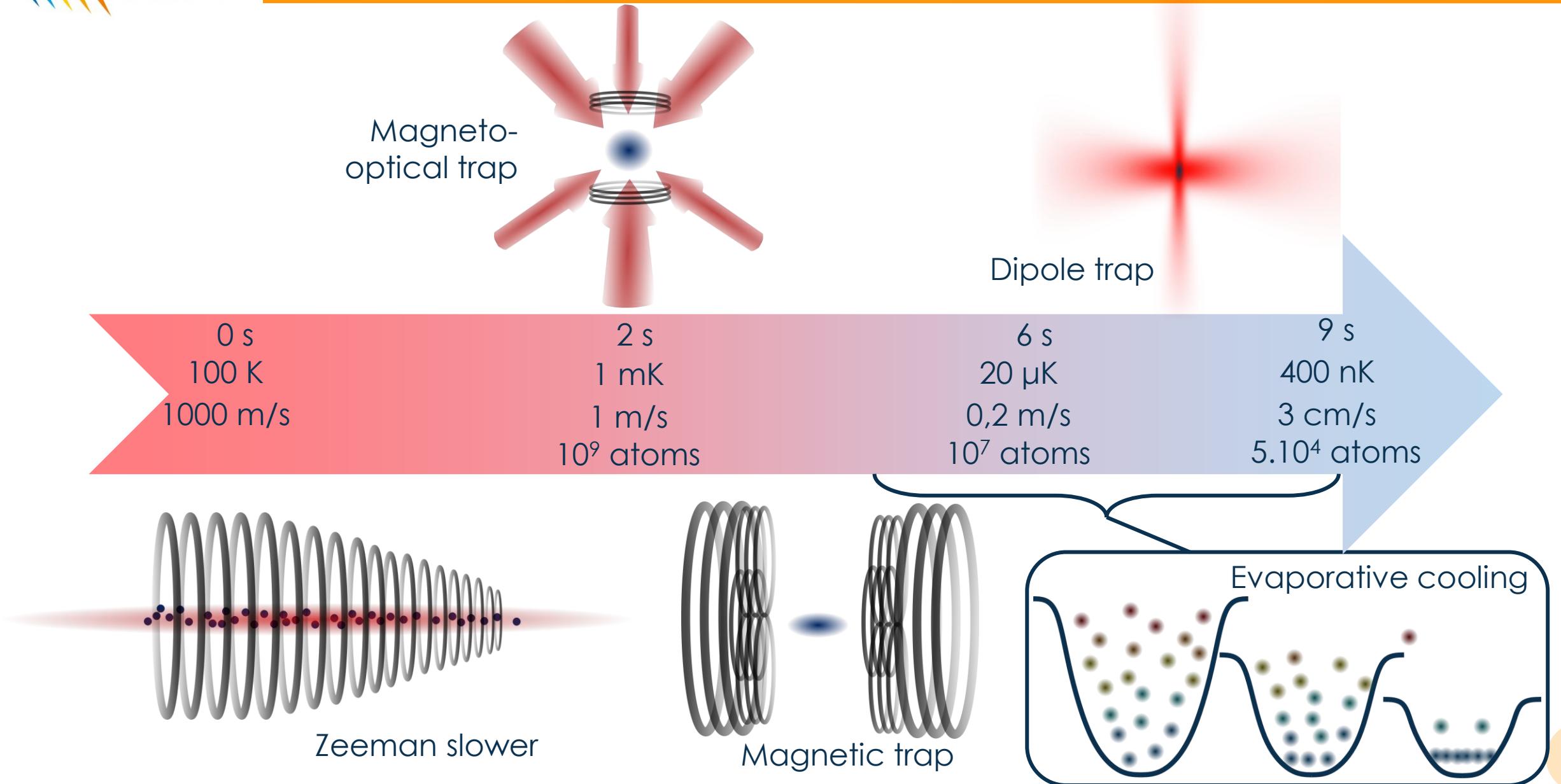


③ Work in progress @LCF

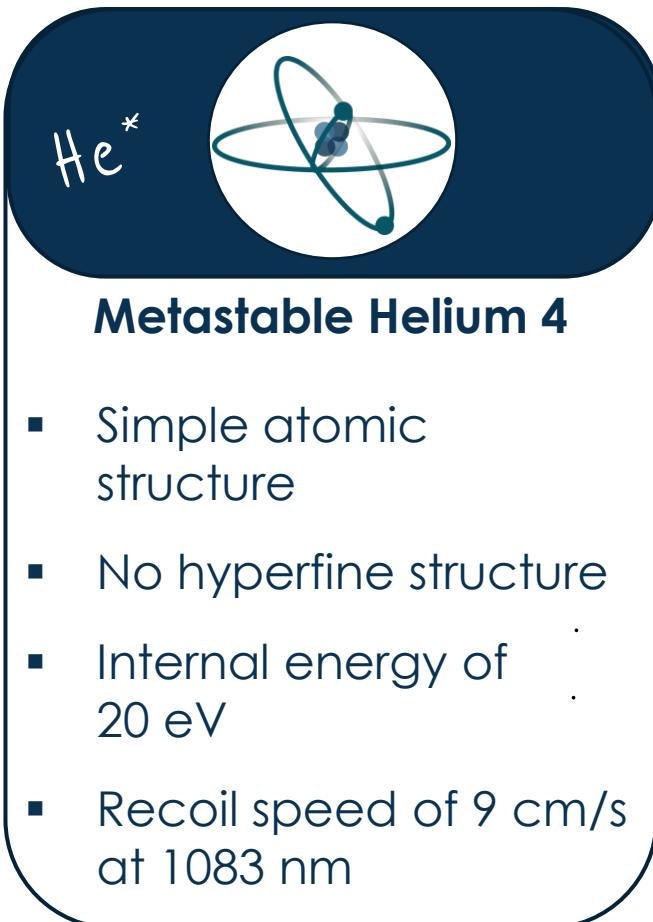
Experimental setup



An experimental cycle

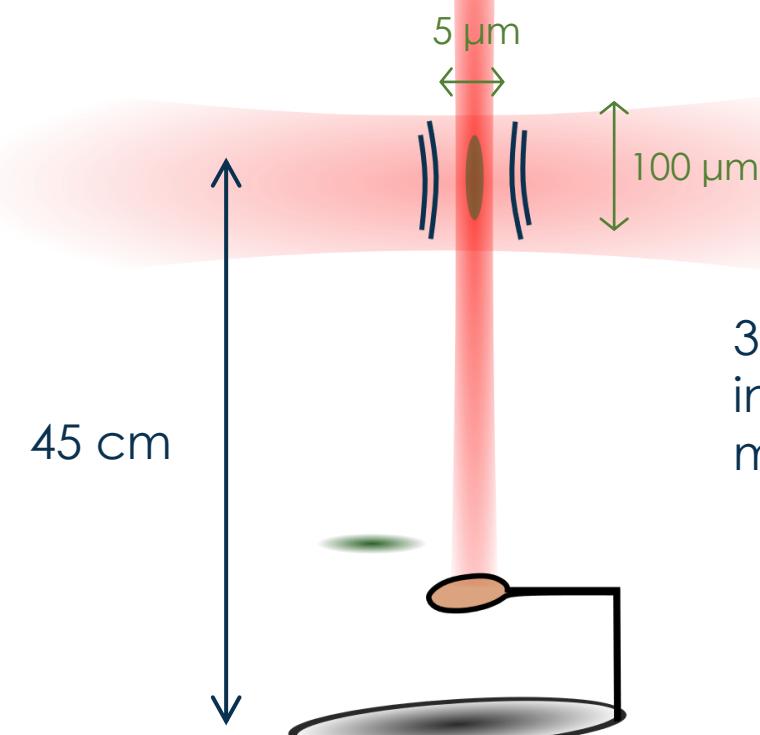


Detection of unique metastable helium atoms

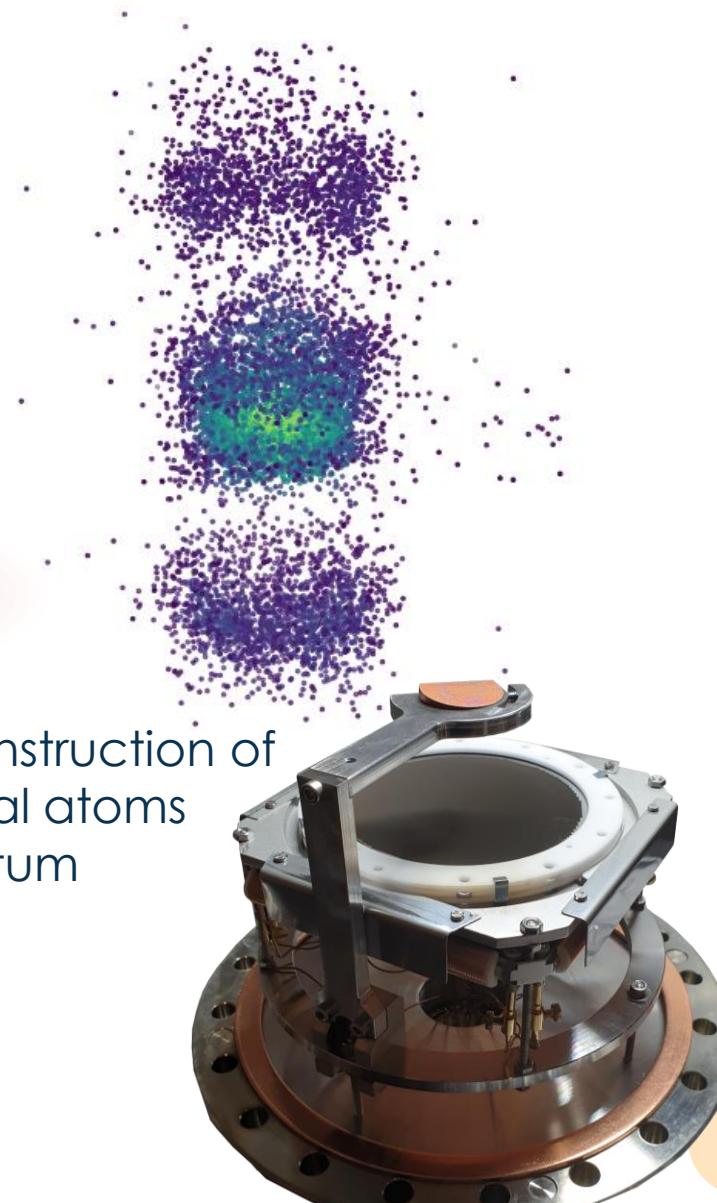


Cigar shape BEC in a crossed dipole trap

- $\omega_{\parallel} = 70 \text{ Hz}$
- $\omega_{\perp} = 1,3 \text{ kHz}$



3D reconstruction of individual atoms momentum



Detection of unique metastable helium atoms

He^*

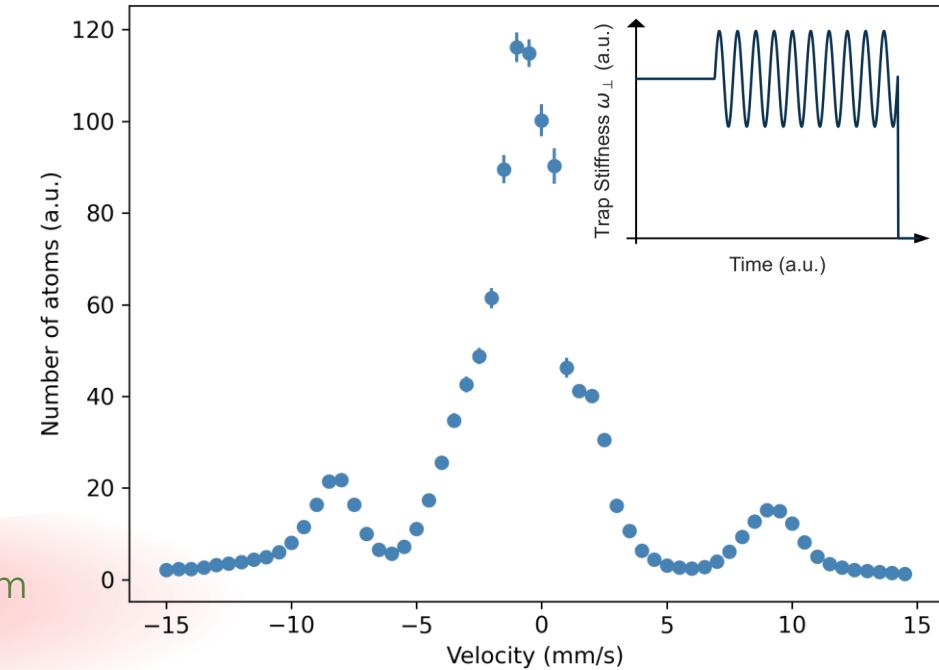
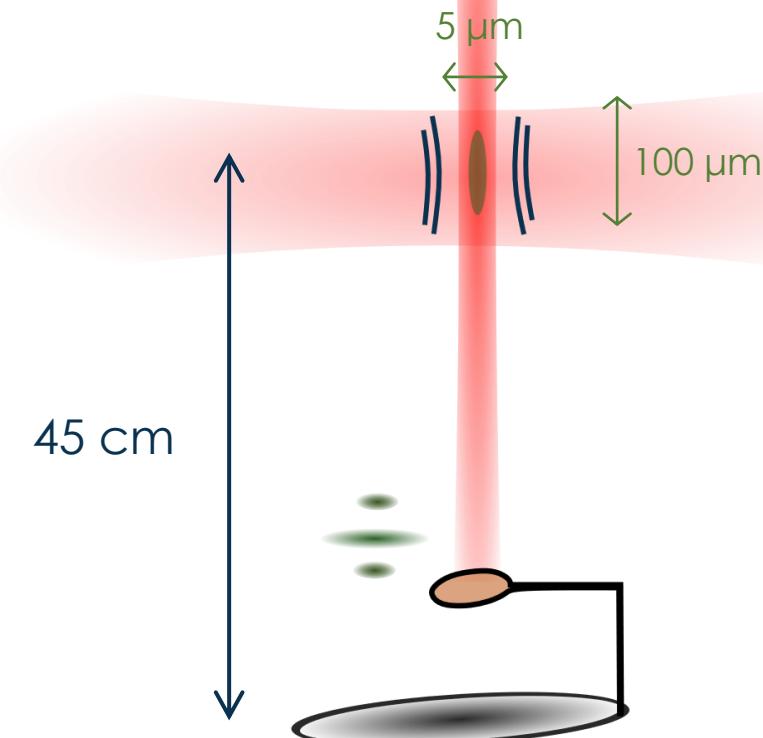


Metastable Helium 4

- Simple atomic structure
- No hyperfine structure
- Internal energy of 20 eV
- Recoil speed of 9 cm/s at 1083 nm

Cigar shape BEC in a crossed dipole trap

- $\omega_{\parallel} = 70 \text{ Hz}$
- $\omega_{\perp} = 1.3 \text{ kHz}$



Probing correlations

Second order correlation function :

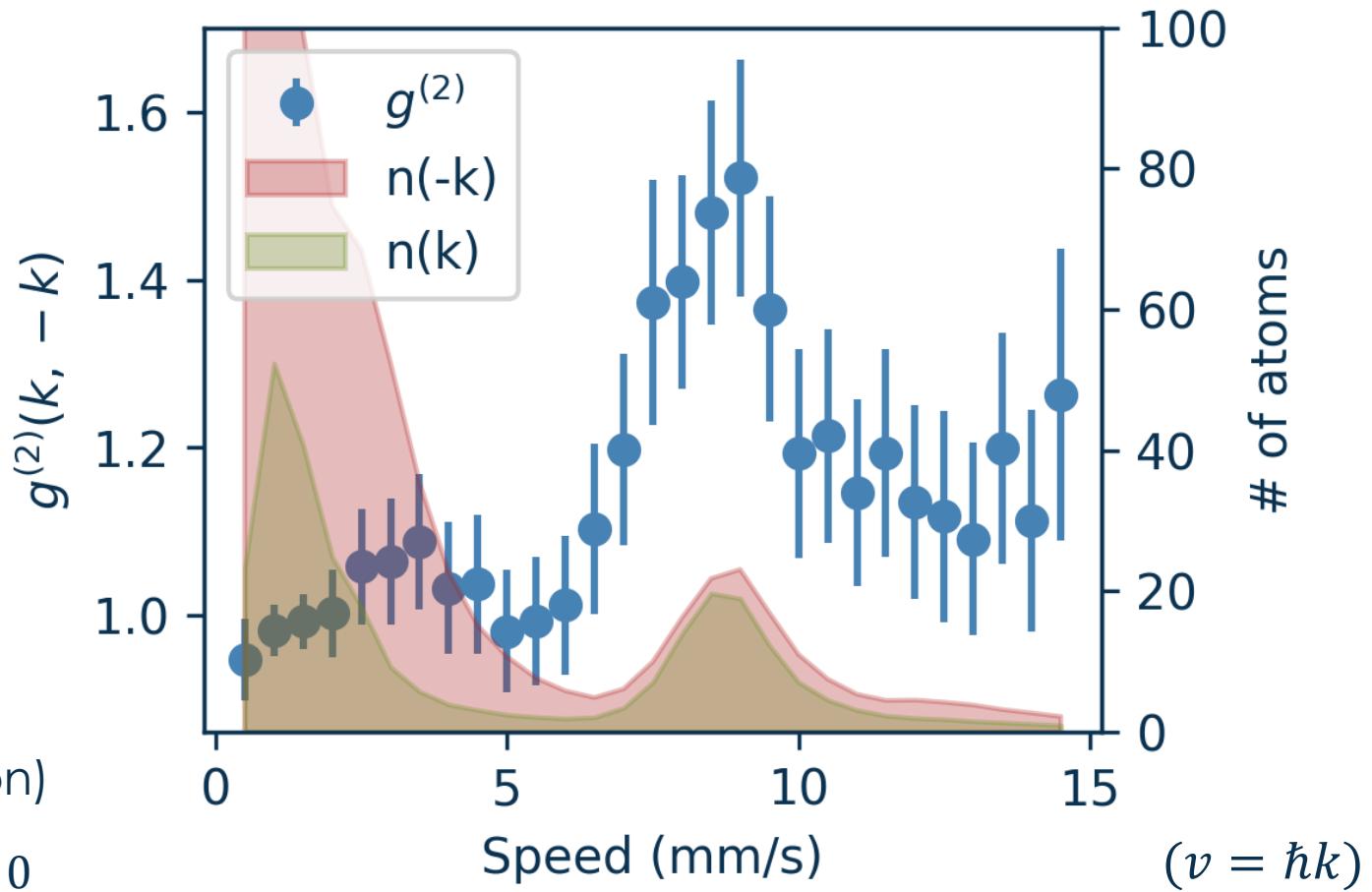
$$g^{(2)}(-k, k) = \frac{\langle n_{-k} n_k \rangle}{\langle n_{-k} \rangle \langle n_k \rangle}$$

$\langle \dots \rangle$: average over experimental realizations

- 👍 Clear correlations between opposite momenta particles,
- 👎 But $g^{(2)}$ still under 2

Perspectives

- 💡 Decrease the temperature
- 💡 Decrease the # of oscillations (\downarrow population)
- 💡 Check $g^{(2)} > 2 \leftrightarrow$ entanglement : $\hat{a}_k \hat{a}_{-k}^\dagger = 0$
- 💡 Study the thermalization



Thank you for your time !

On the experimental side



On the theory side



Some lecture

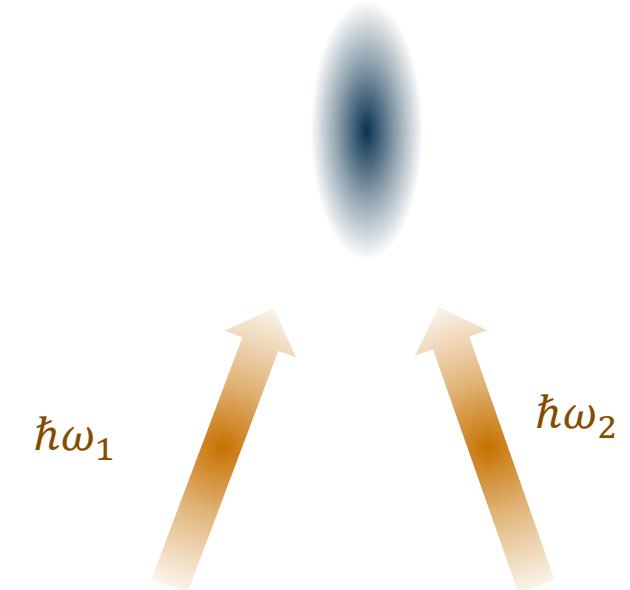
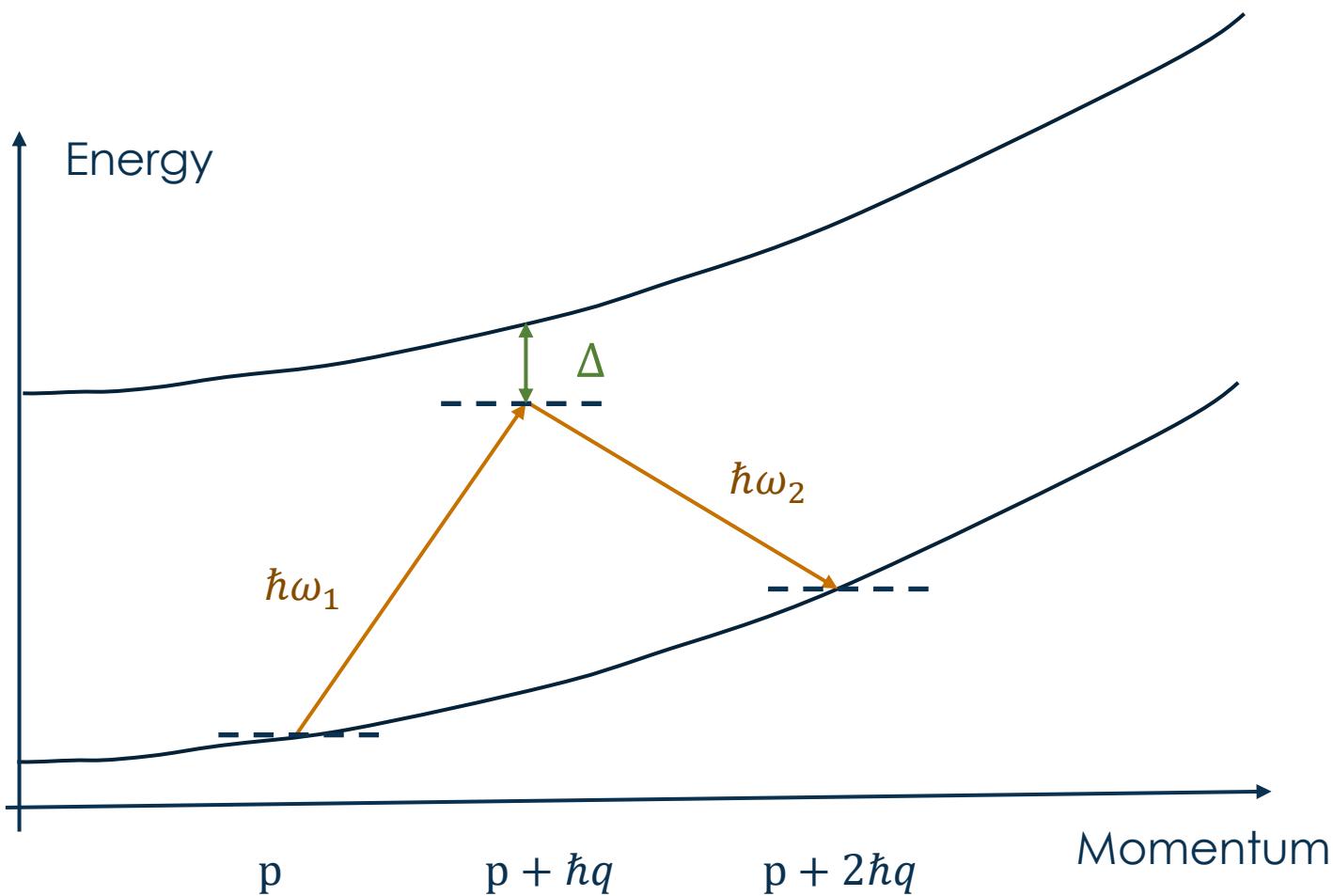
- S. Watson, *An Exposition on Inflationary Cosmology*, 2000
- J.-C. Jaskula *et al.*, Phys. Rev. Lett. **109**, 220401 (2012).
- S. Robertson, F. Michel, and R. Parentani, Phys. Rev. D **95**, 065020 (2017).
- A. Micheli and S. Robertson, Phys. Rev. B **106**, 214528 (2022).



Fundings **anr**®



Supplement : Bragg diffraction for atomic mirrors

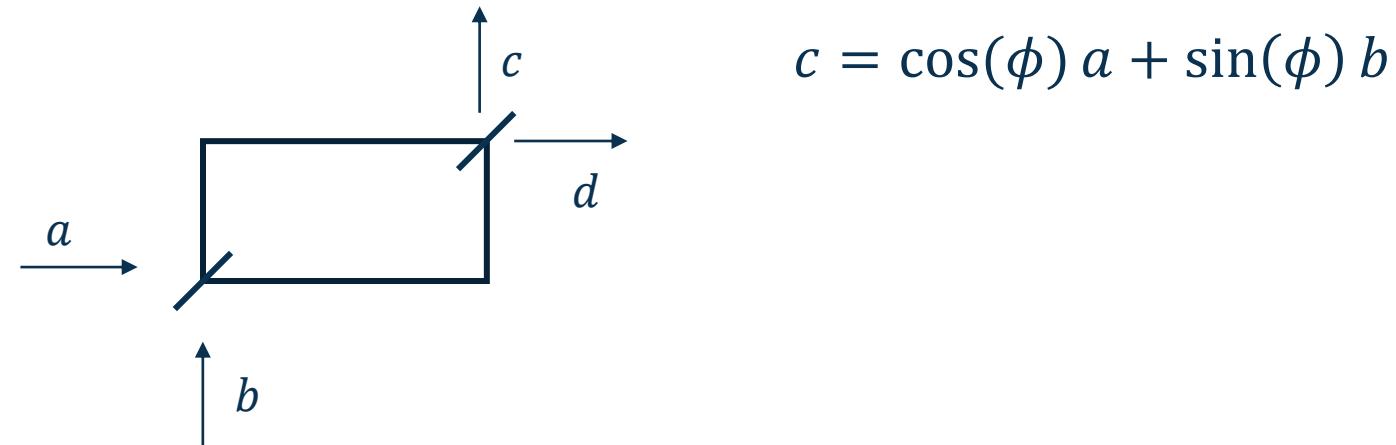


$$q = \frac{2\pi}{\lambda} \sin \theta$$

Supplement : checking the non-separability criteria

$$\langle \hat{n}_k \hat{n}_{-k} \rangle = \langle \hat{b}_k^\dagger \hat{b}_{-k}^\dagger \hat{b}_k \hat{b}_{-k} \rangle = n_k n_{-k} + |\langle \hat{b}_k \hat{b}_{-k} \rangle|^2 + \underbrace{|\langle \hat{b}_k^\dagger \hat{b}_{-k} \rangle|^2}_{0????}$$

if the state is separable : $\leq \overbrace{n_k n_{-k}}$

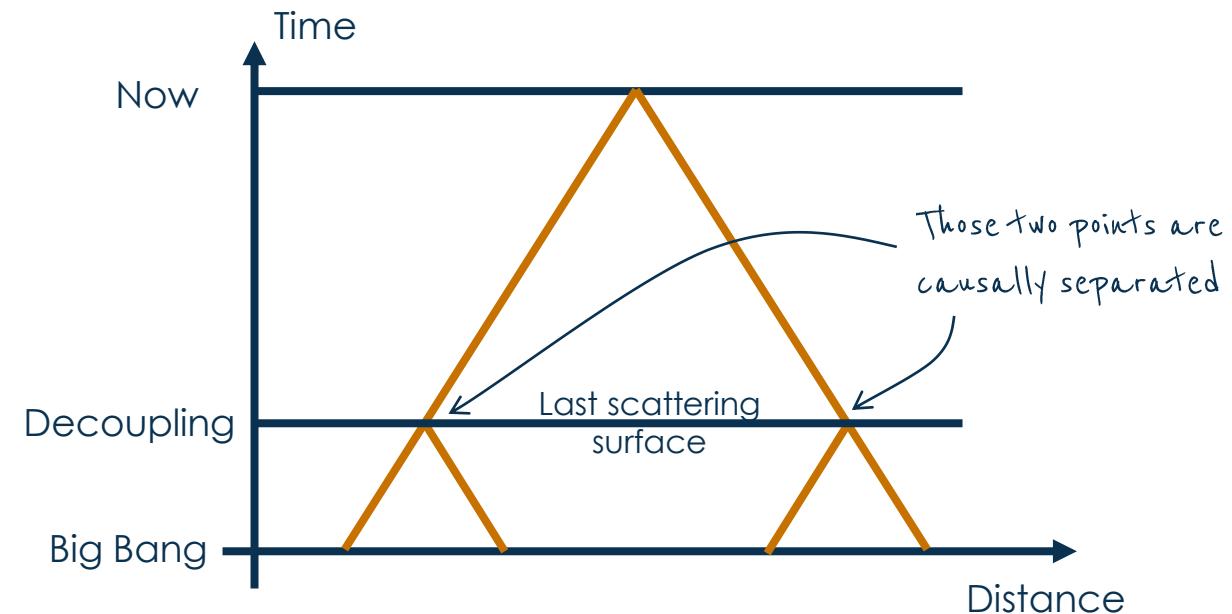
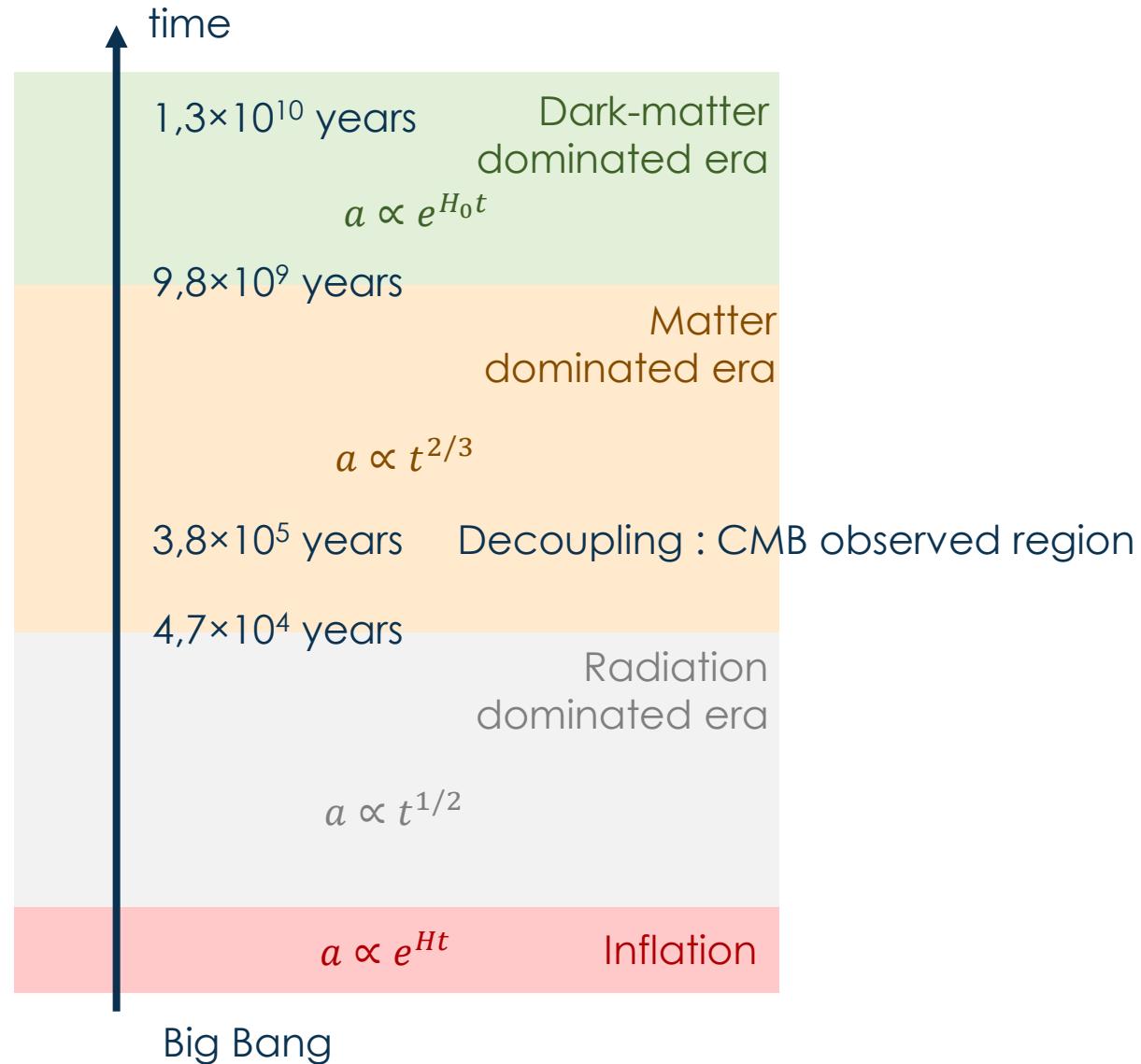


$$c = \cos(\phi) a + \sin(\phi) b$$

$$n_c = \langle c^\dagger c \rangle = \cos(\phi)^2 a^\dagger a + \sin(\phi)^2 b^\dagger b + \cos \phi \sin \phi (a^\dagger b + b^\dagger a)$$

→ Check that this term is zero.

Supplement : inflation scenario



Supplement : excitation spectrum of quasi-particles

